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# Intrabeam Scattering in the Recycler Ring: A Fresh Look

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## **Abstract**

In this report, we estimate the longitudinal and transverse growth rates for the Recycler Ring for a range of relative momentum fraction and a coasting beam of  $2 \times 10^{12}$  antiprotons. The coupling between the horizontal and vertical degrees of freedom is ignored.

# Intrabeam Scattering: Recycler Ring

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This is coulomb scattering between beam particles inside a given beam - bunch or otherwise. There are two approaches in estimating the growth rates:

- **Bjorken-Mtingwa Formalism:** S-matrix formalism and relativistic quantum mechanical treatment - the diffusion equation follows from Fermi's Golden Rule and conservation laws
- **Piwinsky:** Relativistic classical coulomb scattering. Following a treatment of Bane (EPAC-2002) for high energy extension/approximation.

## Assumptions:

- Assumes no coupling between the vertical and horizontal planes.
- No vertical dispersion.
- The lattice version 22 is used - latest.
- Coasting beam - no bunch structure.

# Intrabeam Scattering: Bjorken-Mtingwa Formalism

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$$\frac{1}{T_p} = \frac{1}{\sigma_p} \frac{d\sigma_p}{dt}$$

$$\frac{1}{T_x} = \frac{1}{\epsilon_x^{1/2}} \frac{d\epsilon_x^{1/2}}{dt}$$

$$\frac{1}{T_y} = \frac{1}{\epsilon_y^{1/2}} \frac{d\epsilon_y^{1/2}}{dt}$$

$$\frac{1}{T_i} = 4\pi A(\log) \left\langle \int_0^\infty \frac{d\lambda \lambda^{1/2}}{|(L + \lambda I)|^{1/2}} \left\{ \text{Tr}(L^i) \text{Tr}\left(\frac{1}{L + \lambda I}\right) - 3 \text{Tr}(L^i) \left(\frac{1}{L + \lambda I}\right) \right\} \right\rangle$$

$$A = \frac{r_0^2 c N}{16 \sqrt{2} \pi^3 C \beta_r^3 \gamma^4 \epsilon_x \epsilon_y \sigma_p}$$

$$L = L^p + L^x + L^y$$

$$L^p = \begin{Bmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{Bmatrix}$$

$$L^x = \begin{Bmatrix} 1 & -\gamma\phi_x & 0 \\ -\gamma\phi_x & \gamma^2 \mathcal{H}_x / \beta_x & 0 \\ 0 & 0 & 0 \end{Bmatrix}$$

$$L^y = \begin{Bmatrix} 0 & 0 & 0 \\ 0 & -\gamma^2 \mathcal{H}_y / \beta_y & \gamma\phi_y \\ 0 & -\gamma\phi_y & 1 \end{Bmatrix}$$

$$\mathcal{H} = \frac{\eta^2 + (\beta\eta' - \beta'\eta/2)^2}{\beta}$$

$$\phi = \eta' - \frac{1}{2}\beta'\eta/\beta$$

where the average is taken all over the ring and:

$r_0$  - classical proton radius

$N$  - the number of particles in the bunch

$\beta_r$  - Rleativistic  $v/c$

(log) - Coulomb log factor given by  $\ln(\frac{2}{\theta_{min}})$

$\epsilon_i$  - emittance with  $i = x, y$

$\beta_i, \eta$  - lattice beta and dispersion function with  $i = x, y$

$C$  - RR circumference

$\sigma_p$  - the relative energy spread

One can obtain a useful approximation for high  $\gamma$  ( $\approx$ : 10 or higher)

$$\frac{\beta_x}{\epsilon_x}, \frac{\beta_y}{\epsilon_y} \ll \frac{\gamma^2 \eta^2}{\epsilon_x \beta_x}, \frac{\beta_x \gamma^2 \phi^2}{\epsilon_x}, \frac{\gamma^2}{\sigma_p^2}$$

$$\frac{1}{\tau_p} = \pi^2 \alpha^2 M N (\log) \frac{\gamma}{\Gamma \sigma_p^2} \int_0^\infty d\lambda \frac{\sqrt{\lambda} [2a\lambda + b]}{[\lambda^3 + a\lambda^2 + b\lambda + c]^{3/2}}$$

$$a = \frac{\gamma^2 \eta^2}{\epsilon_x \beta_x} + \frac{\gamma^2 \phi^2 \beta_x}{\epsilon_x} + \frac{\gamma^2}{\sigma_p^2}$$

$$b = \left[ \left( \frac{\beta_x}{\epsilon_x} + \frac{\beta_y}{\epsilon_y} \right) \left( \frac{\gamma^2 \eta^2}{\epsilon_x \beta_x} + \frac{\gamma^2}{\sigma_p^2} \right) + \frac{\beta_x \beta_y \gamma^2 \phi^2}{\epsilon_x \epsilon_y} \right]$$

$$c = \frac{\beta_x \beta_y}{\epsilon_x \epsilon_y} \left( \frac{\gamma^2 \eta^2}{\epsilon_x \beta_x} + \frac{\gamma^2}{\sigma_p^2} \right)$$

$$\frac{1}{\tau_x} = \frac{\sigma_p^2}{\tau_p} \left\langle \left[ \frac{\eta^2}{\epsilon_x \beta_x} + \frac{\beta_x \phi^2}{\epsilon_a} \right] \right\rangle$$

$$\Gamma_{unbunched} = \frac{(2\pi)^{5/2}}{\sqrt{2}} \beta^3 \gamma^3 M^3 \epsilon_x \epsilon_y \sigma_p C$$

# Intrabeam Scattering - Piwinsky Formalism

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$$\frac{1}{T_p} = A \langle \left[ \frac{\sigma_h^2}{\sigma_p^2} f(a_p, b_p, q) \right] \rangle$$

$$\frac{1}{T_x} = A \langle \left[ f\left(\frac{1}{a_p}, \frac{b_p}{a_p}, \frac{q}{a_p}\right) + \frac{\eta_x^2 \sigma_h^2}{\beta_x \epsilon_x} f(a_p, b_p, q) \right] \rangle$$

$$\frac{1}{T_y} = A \langle \left[ f\left(\frac{1}{b_p}, \frac{a_p}{b_p}, \frac{q}{b_p}\right) + \frac{\eta_y^2 \sigma_h^2}{\beta_y \epsilon_y} f(a_p, b_p, q) \right] \rangle$$

$$\frac{1}{\sigma_h^2} = \frac{1}{\sigma_p^2} + \frac{\eta_x^2}{\beta_x \epsilon_x} + \frac{\eta_y^2}{\beta_y \epsilon_y}$$

$$a_p = \frac{\sigma_h}{\gamma} \sqrt{\frac{\beta_x}{\epsilon_x}}$$

$$b_p = \frac{\sigma_h}{\gamma} \sqrt{\frac{\beta_y}{\epsilon_y}}$$

$$q = \sigma_h \beta \sqrt{\frac{2d}{r_0}}$$

$$f(a_p, b_p, q) = 8\pi \int_0^1 du \frac{1 - 3u^2}{PQ} \left\{ 2 \ln \left[ \frac{q}{2} \left( \frac{1}{P} + \frac{1}{Q} \right) \right] - 0.577 \right\}$$

$$P^2 = a_p^2 + (1 - a_p^2)u^2$$

$$Q^2 = b_p^2 + (1 - b_p^2)u^2$$

The above formalism has to be modified for rings with  $\langle \mathcal{H} \rangle \neq \langle \eta^2/\beta \rangle$ . Therefore by replacing  $\mathcal{H}$  with  $\eta^2/\beta$ , we obtain the [modified Piwinsky Formalism](#). Also with proper choice of the maximum impact parameter  $d$ , the coulomb logarithms can be made same. Thus we can write:

$$\frac{1}{T_p} \approx \frac{r_0^2 c N (\log) \sqrt{2\pi}}{8C \gamma^3 (\epsilon_x \epsilon_y)^{3/4} \sigma_p^3} \langle \sigma_H h(a, b) (\beta_x \beta_y)^{-1/4} \rangle$$

$$h(a, b) = \frac{2\sqrt{ab}}{\pi} \int_0^\infty \frac{du}{\sqrt{(a^2 + u^2)(b^2 + u^2)}}$$

$$\frac{1}{\sigma_H^2} = \frac{1}{\sigma_p^2} + \frac{\mathcal{H}_x}{\epsilon_x} + \frac{\mathcal{H}_y}{\epsilon_y}$$

# Computations Summary

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## Notes:

- $N = 2 \times 10^{12}$
- $\epsilon_x = 5 - 20 \pi$ -mm-mr (Normalized 95%)
- $\epsilon_y = \epsilon_x$  for simplicity
- Rates are computed as a function  $\sigma_p = \Delta/P$  for various values of  $\epsilon_x$ .
- Lattice averaging is done via breaking up the Ring into 1300 parts
- One has to be very careful with numerical integrations as quantities of relative magnitudes of E+24 in the denominator!



# Recycler Ring IBS Results

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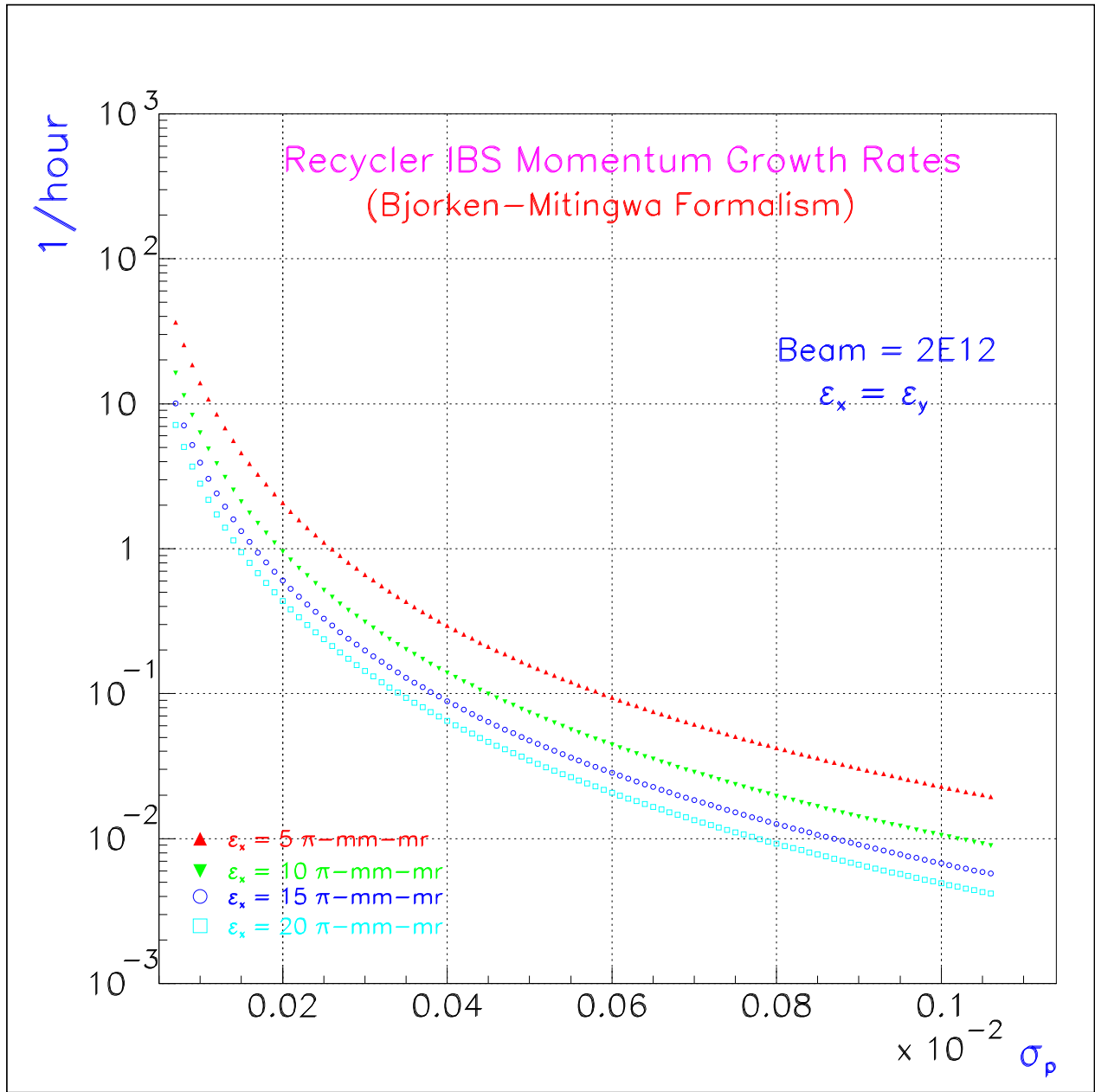
Quantity	$\epsilon_x$ [ $\pi$ -mm-mr]	$\sigma_p$	<b>BJM</b> Formalism	<b>PWKB</b> Formalism
Mom. Growth (1/hr)	5.0	0.00025	1.11	1.17
Trans. Growth (1/hr)	5.0	0.00025	0.06	0.06
$d(\epsilon_x)/dt$ ( $\pi$ -mm-mr/hr)	5.0	0.00025	0.30	0.32
Mom. Growth (1/hr)	5.0	0.00075	0.05	0.07
Trans. Growth (1/hr)	5.0	0.00075	0.02	0.03
$d(\epsilon_x)/dt$ ( $\pi$ -mm-mr/hr)	5.0	0.00075	0.12	0.17
Mom. Growth (1/hr)	10.0	0.00025	0.52	0.45
Trans. Growth (1/hr)	10.0	0.00025	0.01	0.01
$d(\epsilon_x)/dt$ ( $\pi$ -mm-mr/hr)	10.0	0.00025	0.14	0.12
Mom. Growth (1/hr)	10.0	0.00075	0.02	0.03
Trans. Growth (1/hr)	10.0	0.00075	0.01	0.01
$d(\epsilon_x)/dt$ ( $\pi$ -mm-mr/hr)	10.0	0.00075	0.06	0.08

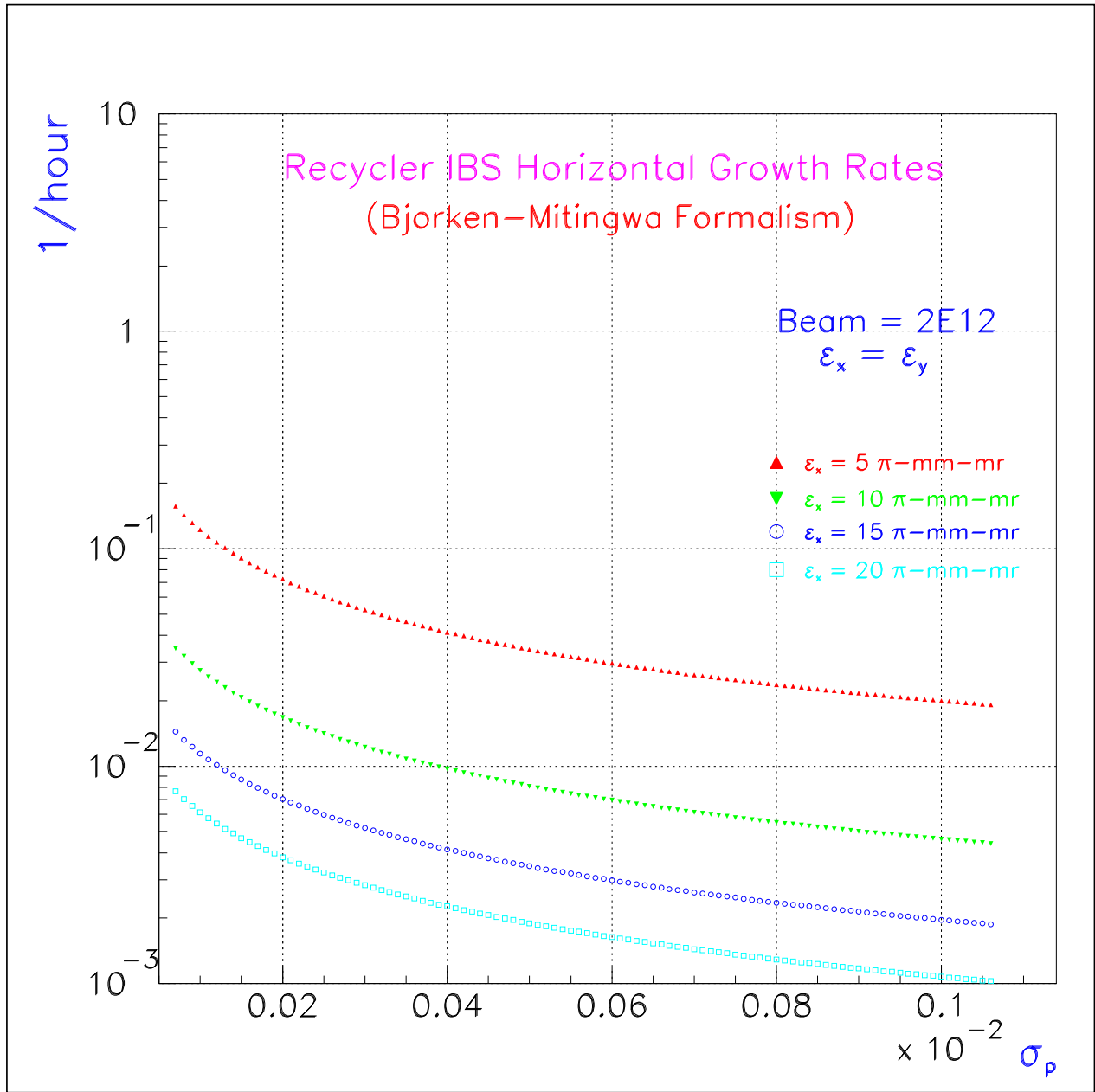
Beam = 2.0E+12

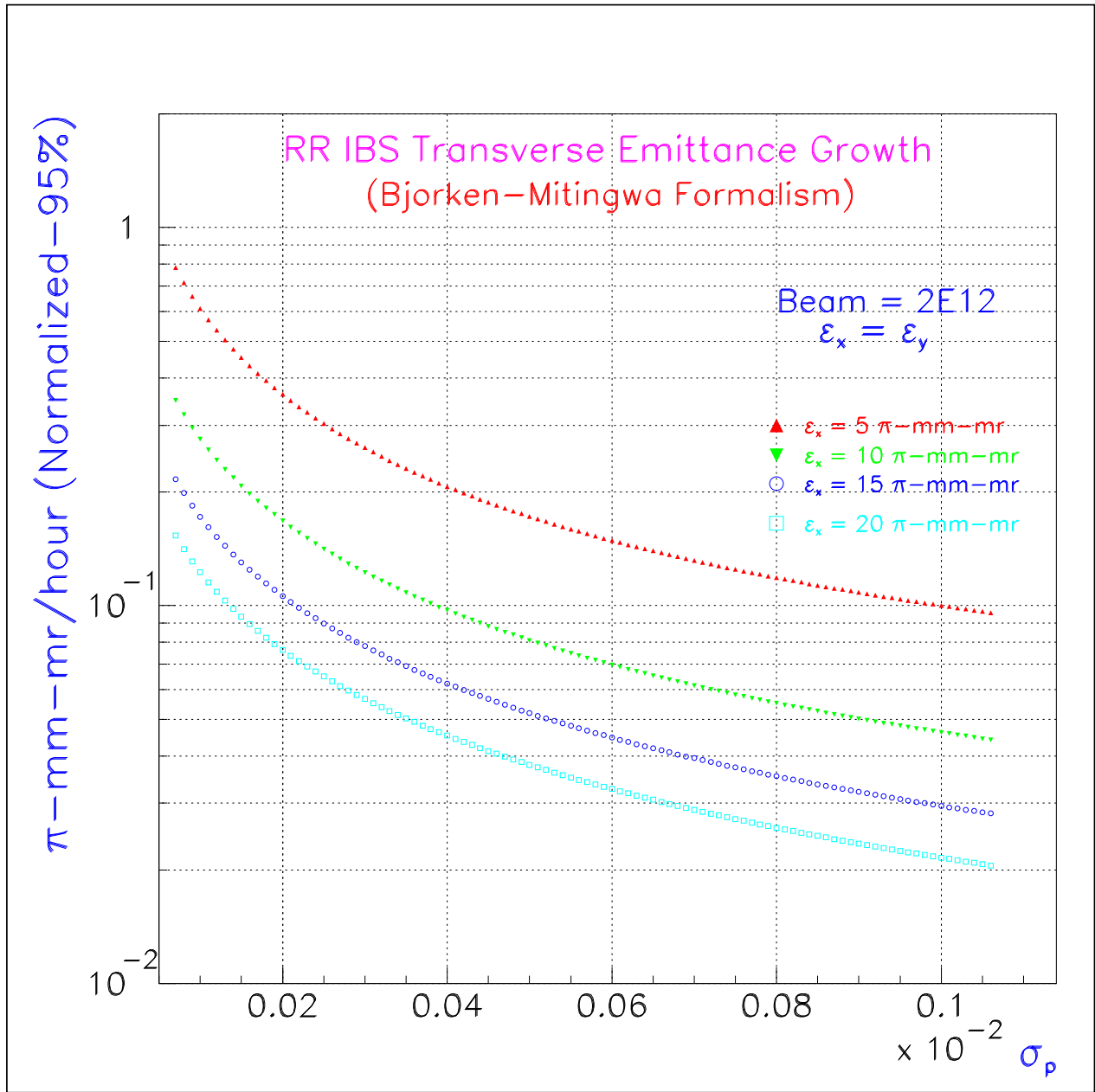
$\epsilon_x = \epsilon_y$

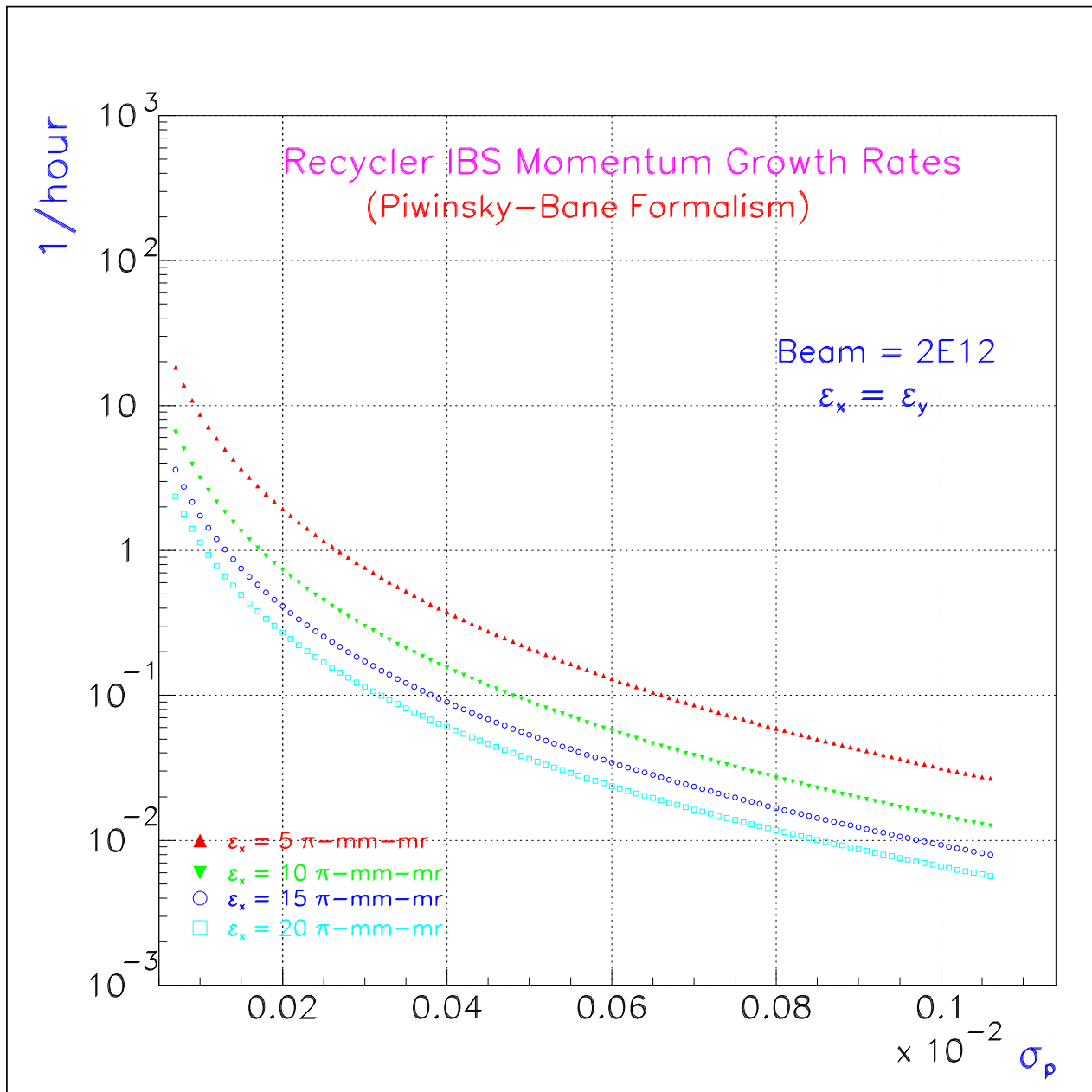
clog = 20.0

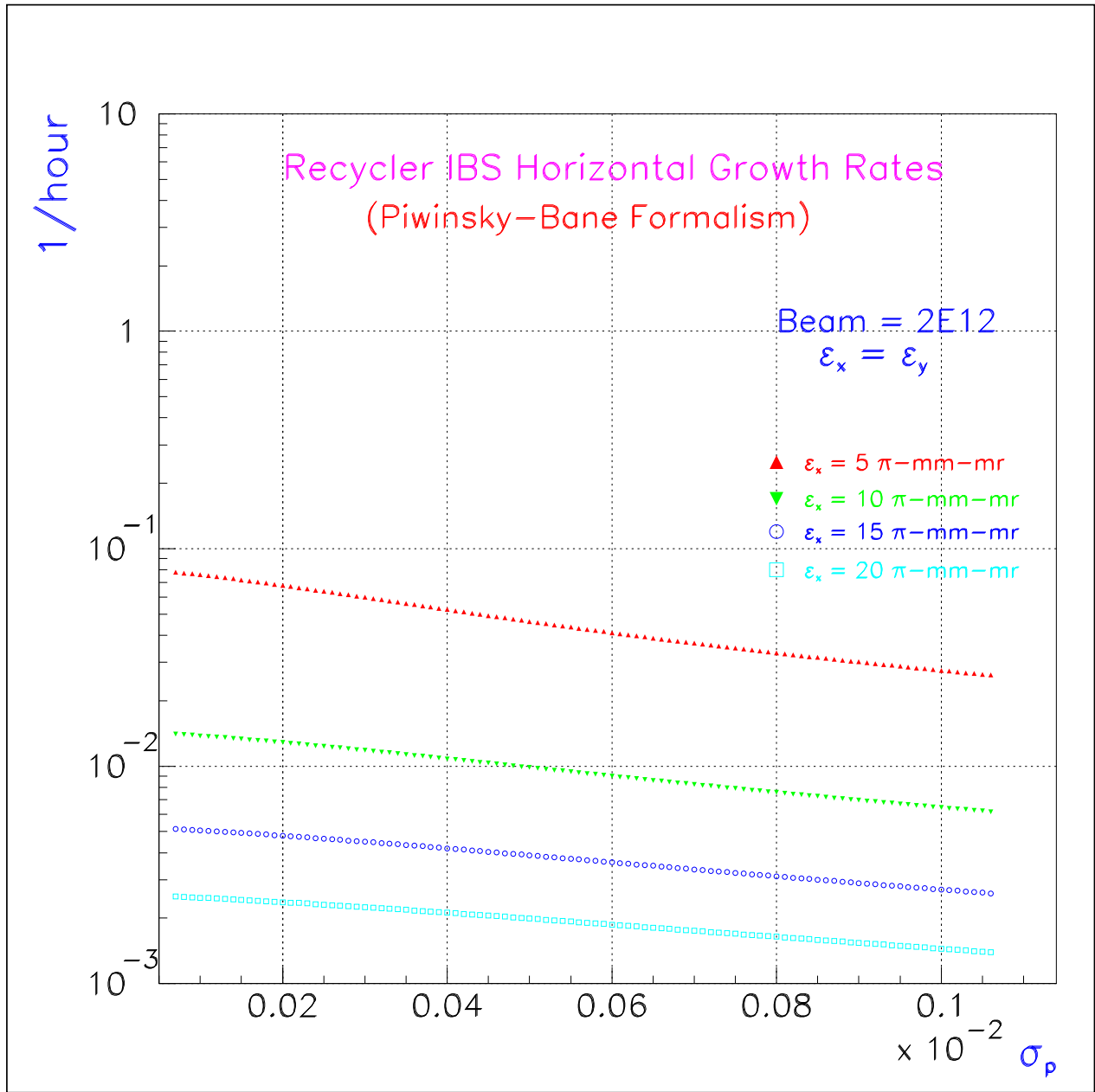
Coasting Beam

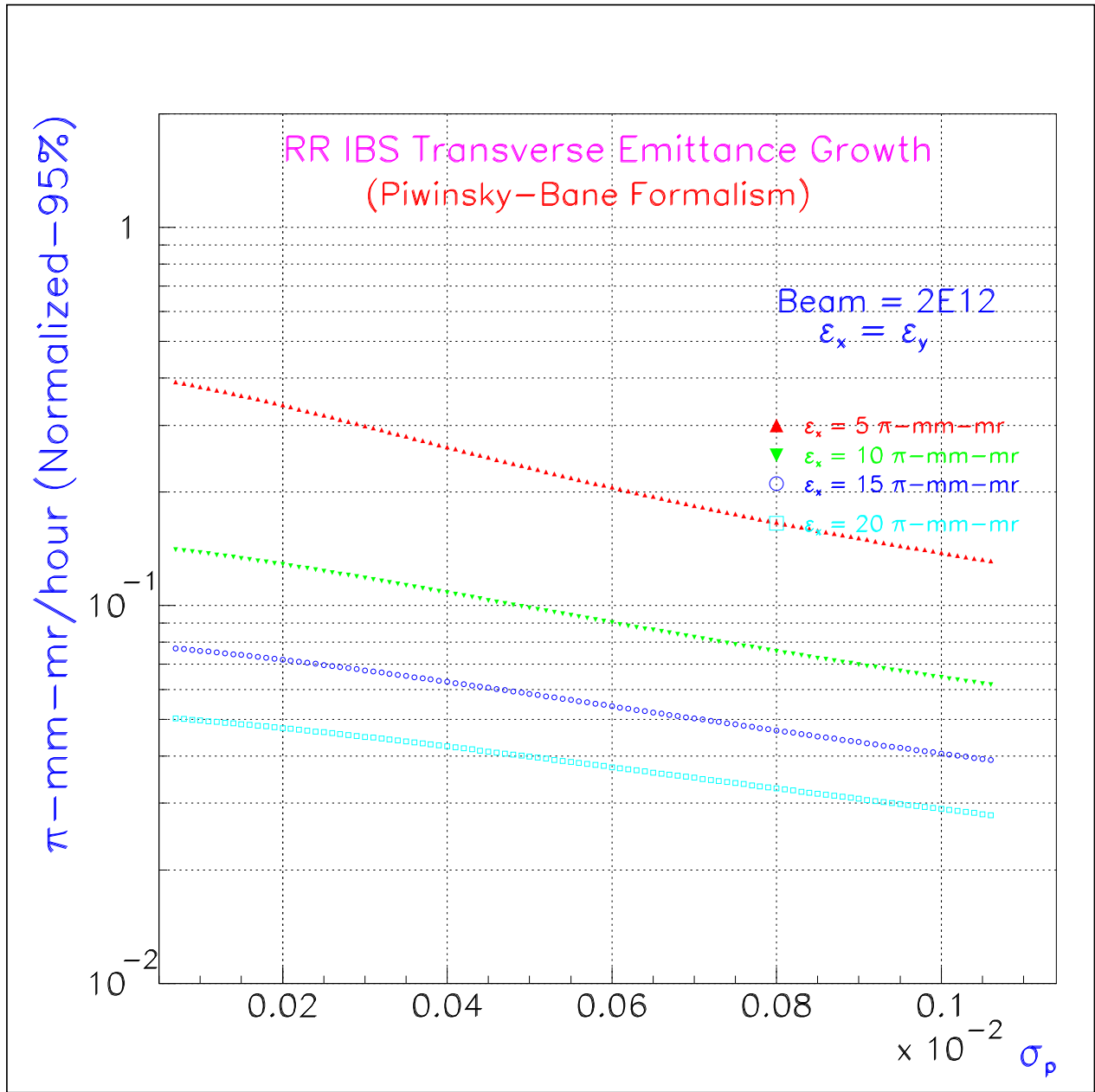












# Summary and Outlook

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We completed a fresh look at the intrabeam scattering rates for the Recycler Ring using Bjorken-Mtingwa formalism and that of Piwinsky. With in the high energy approximation, both results are consistent with each other.

- The transverse emittance growth rate is about  $0.12 \pi\text{-mm-mr}/\text{hour}$  for a beam of  $2 \times 10^{12}$  particles,  $\epsilon_{tN95\%} = 10 \pi\text{-mm-mr}$ ,  $\sigma_p = 0.00025$ .
- The longitudinal growth rate is  $0.5/\text{hr}$  - may test the performance of the longitudinal cooling ??
- There are also similar calculations by Pat Colestock and Alexi Burov. The results are comparable.

Looking ahead:

- Look at the final equation of state for the Recycler Ring - all cooling and heating processes included both in transverse and longitudinal degrees of freedom
- Detailed treatment of Ion production and clearing mechanisms for the Recycler Ring.



## **Acknowledgement**

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## **References**

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- [2] A. Piwinski, Proc. 9th Int. Conf. on High Energy Accelerators, 1974, P.405.
- [3] K. L. F. Bane, Proceedings of EPAC 2002, Paris, France.