

Implications of Higgsless Models of EWSB

Note Title

1/6/2004

OUTLINE

- 1) HIGGSLESS EWSB
- 2) PHENOMENOLOGICALLY VIABLE MODEL
- 3) MORE UNITARITY ISSUES
- 4) COLLIDER PHENOMENOLOGY

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IBL.

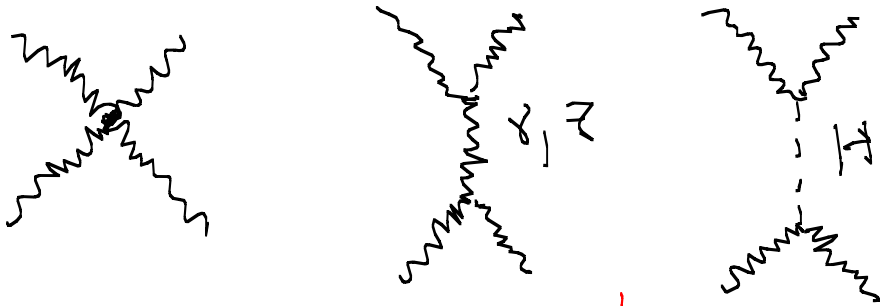
J. TIRRO

hep-ph/0312193

WHAT IS THE HIGGS FOR?

- GIVES MASS TO THE W, Z
- UNITARIZES $W_L W_L \rightarrow W_L W_L$ SCATTERING AMPLITUDE

$$A = a s^2 + b s + c + \mathcal{O}\left(\frac{1}{s}\right)$$



TERMS GROWING
LIKE s^2 CANCEL

TERMS GROWING LIKE
 s CANCEL

GAUGE FIELDS ON AN INTERVAL

TAKE $M^4 \times I$
 \uparrow
 INTERVAL

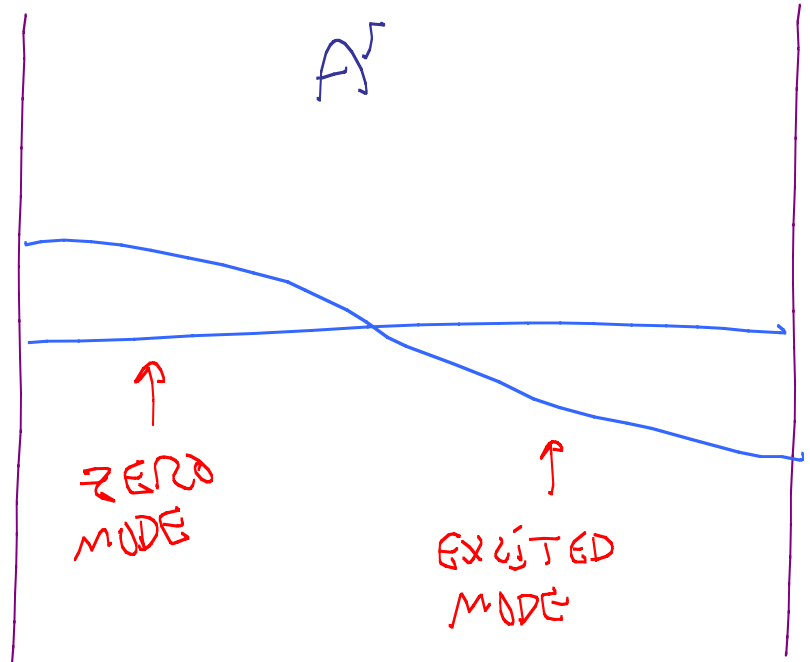
PUT A GAUGE FIELD
 IN THE BULK

$$A^\mu = A^{(0)\mu} + \sum_{n=1}^{\infty} A_n^{(\mu)} \cos\left(\frac{n\pi y}{r_c}\right) + \sum_{n=1}^{\infty} A_n^{(\mu)} \sin\left(\frac{n\pi y}{r_c}\right)$$

$$A^5 = A^{(0)5} + \sum_{n=1}^{\infty} A_n^{(5)} \cos\left(\frac{n\pi y}{r_c}\right) + \sum_{n=1}^{\infty} A_n^{(5)} \sin\left(\frac{n\pi y}{r_c}\right)$$

NORMAL ORBIFOLD
 BOUNDARY CONDITIONS

$$\partial_5 A^\mu = 0, \quad A^5 = 0$$



NOTE MOMENTUM

$$p^2 = p_\mu p^\mu + \underbrace{p_5 p^5}_{= \partial_5 \partial^5} \leftarrow \text{LOOKS LIKE A MASS TERM}$$

MASS IS CONNECTED
 TO CURVATURE IN 5th \rightarrow

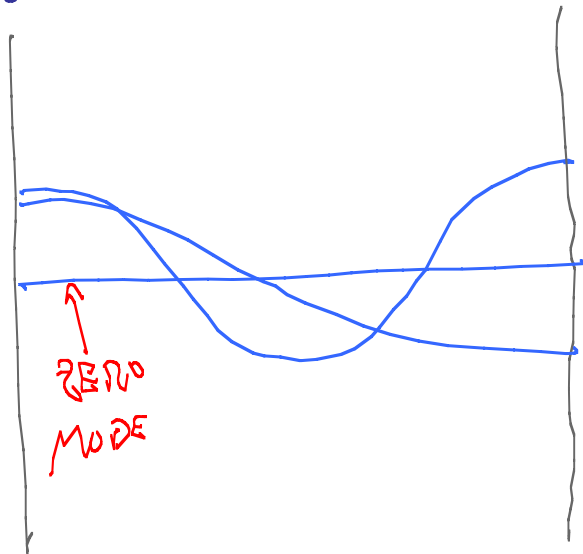
DIFFERENT BOUNDARY CONDITIONS

CSAKI, GROJEAN,
MURAYAMA, PILLO & TERNING
hep-ph/0305237

GENERATING MASSES

$$\partial_5 A^\mu = 0$$

$$\partial_5 A^\mu = 0$$



NORMAL ORBIFOLD

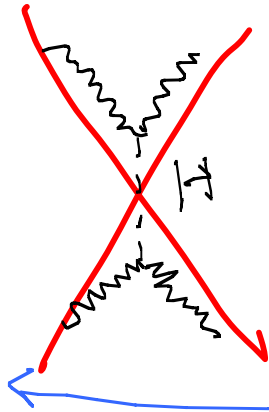
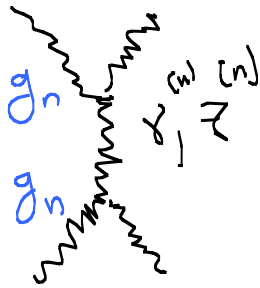
$$\partial_5 A^\mu = 0$$

$$A^\mu = 0$$



GAUGE-BREAKING
BC'S

UNITARITY



SUM RULES

$$g_w^2 = \sum_{n=1}^{\infty} g_n^2$$

TERMS GROWING LIKE s^2 CANCEL

TERMS GROWING LIKE s CANCEL

LIKE

$$4M_w^2 g_w^2 = 3 \sum_{n=1}^{\infty} M_n^2 g_n^2$$

VALID AT

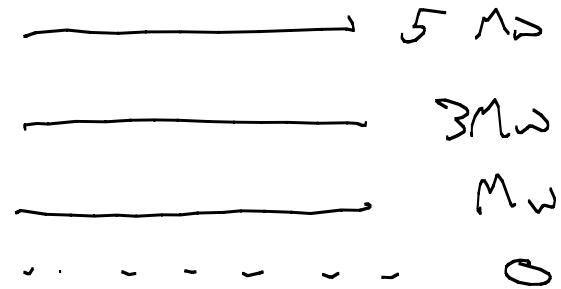
ASYMPTOTICALLY HIGH s

(POSSIBLY) VIABLE MODEL

- IN FLAT SPACE THE MASS SPECTRUM IS ROUGHLY $\frac{2n-1}{4R}$

TOO LIGHT!

← COMPACTIFICATION RADIUS



- WITHOUT A HIGGS DOUBLET, NO CUSTODIAL SU(2). g -PARAMETER WAY OFF

BOTH PROBLEMS SOLVED BY

$SU(2)_L \times SU(2)_R \times U(1)_{B-L}$
IN A RANDALL-SUNDIUM (WARPED) SCENARIO

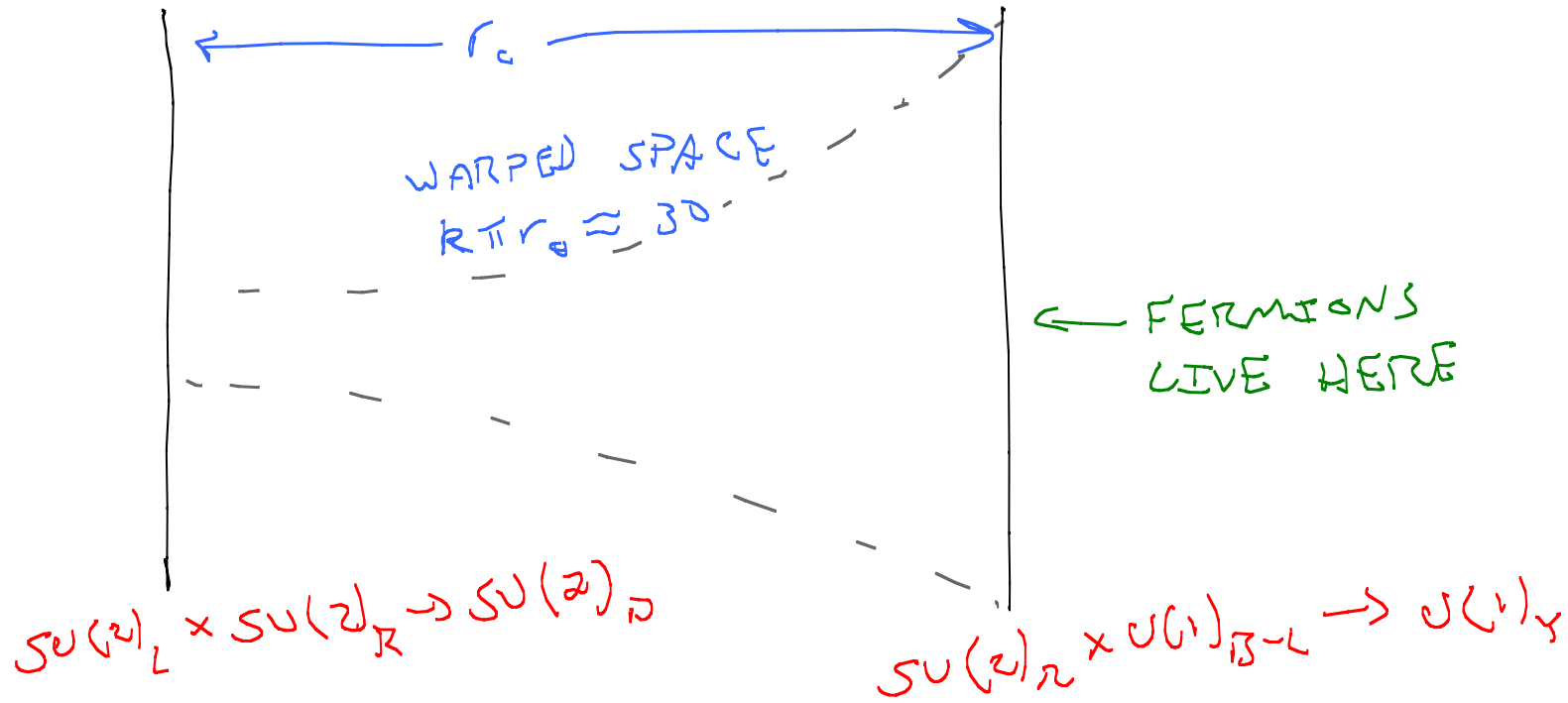
CSAKI, GROJEAN, PYLE,
+ FERNING hep-ph/0308088
NOMURA hep-ph/0309148

AGASHE, DELGADO, MAY
+ SUNDIUM hep-ph/0308006

BREAKING PATTERN

"TeV" BRANE

"PLANCK" BRANE



OVERALL $SU(2)_L \times SU(2)_R \times U(1)_{B-L} \rightarrow U(1)_Q$

MODEL PARAMETERS

g_L - $SU(2)_L$ COUPLING
 g_R - $SU(2)_R$ COUPLING
 g' - $U(1)$ COUPLING
 k, λ - SCALING BETWEEN
 PLANCK & TEV SCALES

$$k = \frac{g_R}{g_L} \quad \lambda = \frac{g'}{g_L}$$

\tilde{g}_L
 \tilde{g}_Y } BRANE LOCALIZED
 KINETIC TERMS
 ON PLANCK BRANE
 SEE NOMURA

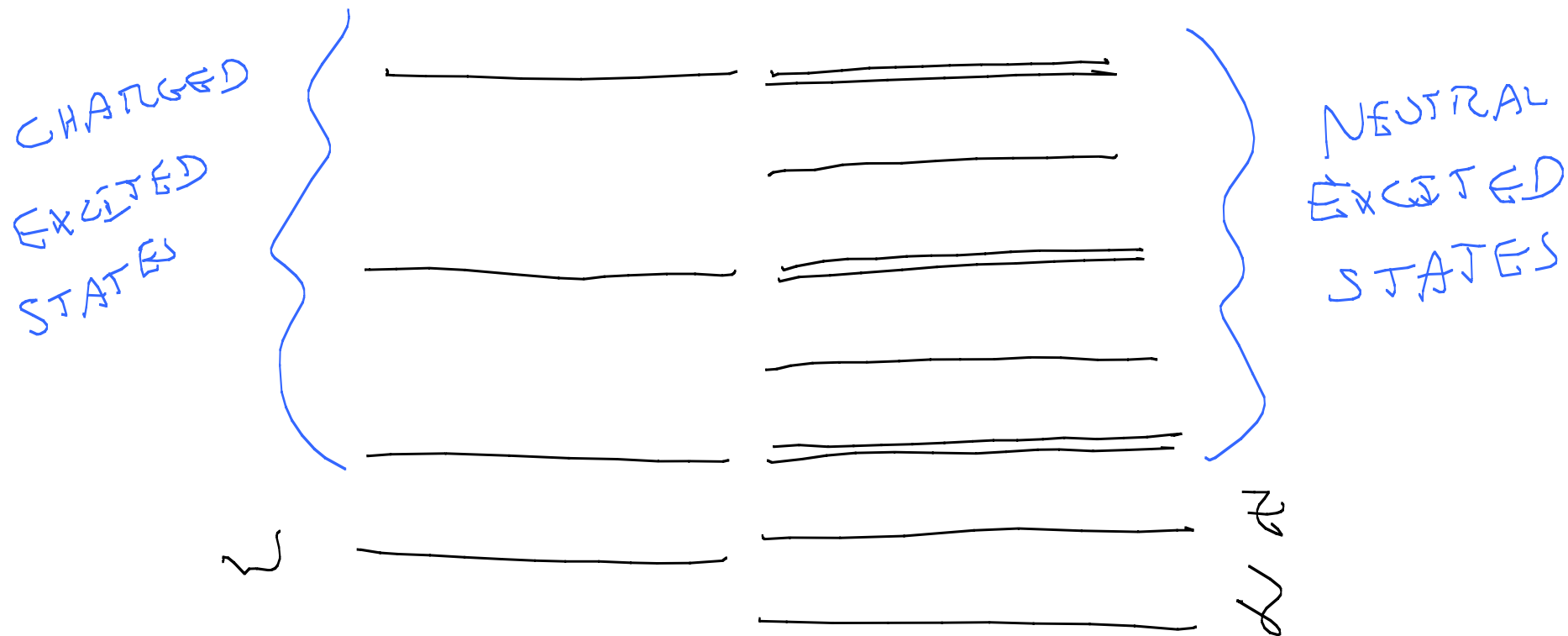
$$\delta_L = \frac{k g_L^2}{2 \tilde{g}_L^2}, \quad \delta_Y = \text{etc.}$$

$$\delta_Y = -2 \frac{k^2 \lambda^2}{\lambda^2 + k^2} \delta_L$$

k, λ → FIXED BY M_W
 δ_L → " " G_F
 λ → " " M_Z

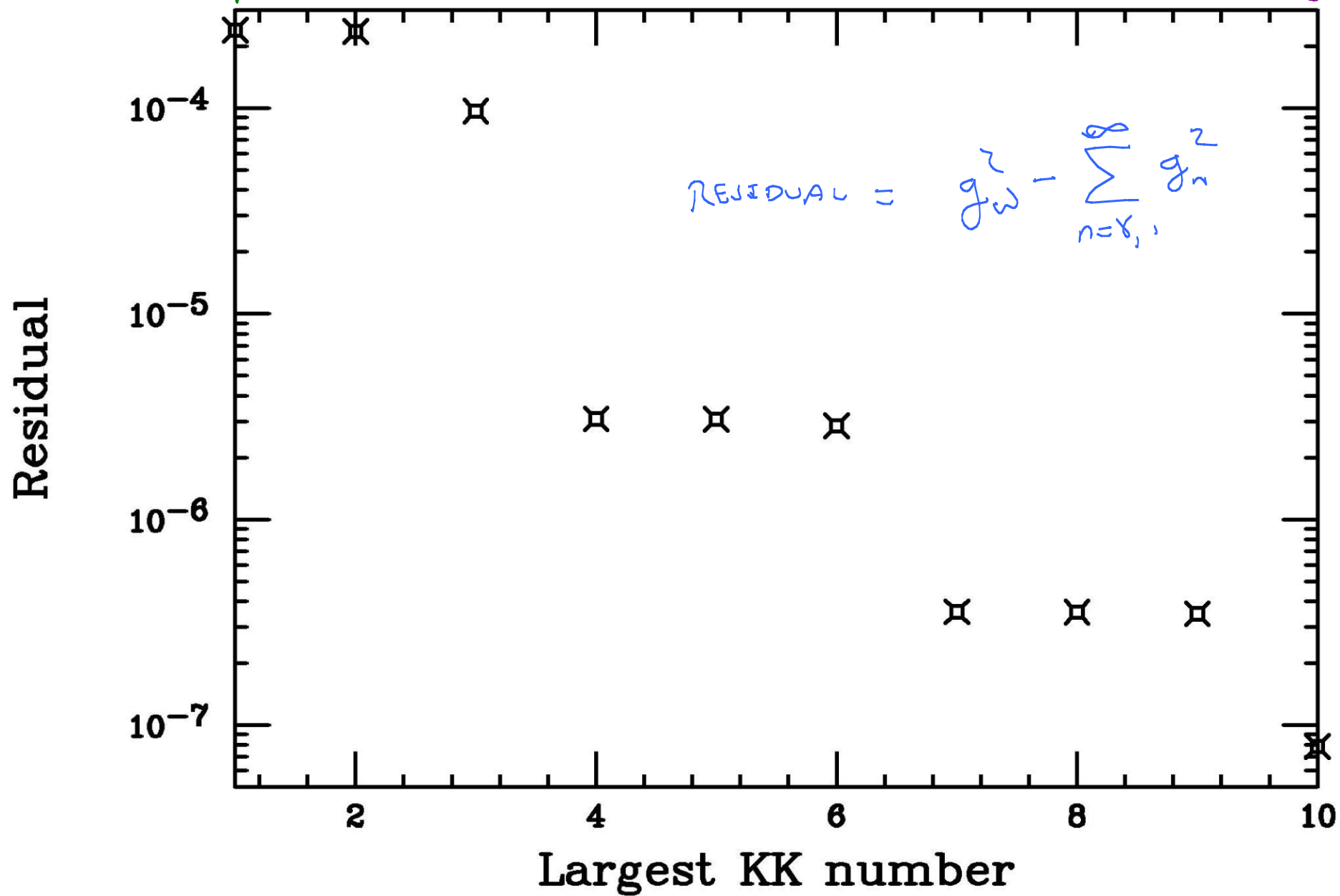
k LEFT AS FREE PARAMETER

SPECTRUM



ASYMPTOTIC UNITARITY \rightarrow SUM RULES

$k=2$



PRECISION ELECTROWEAK DATA

EXAMPLE CALCULATION

$\sin^2 \theta_{05}$ ← DEFINED FROM $\frac{M_W^2}{M_Z^2}$
EXACT IN OUR SCHEME

CAN DEFINE

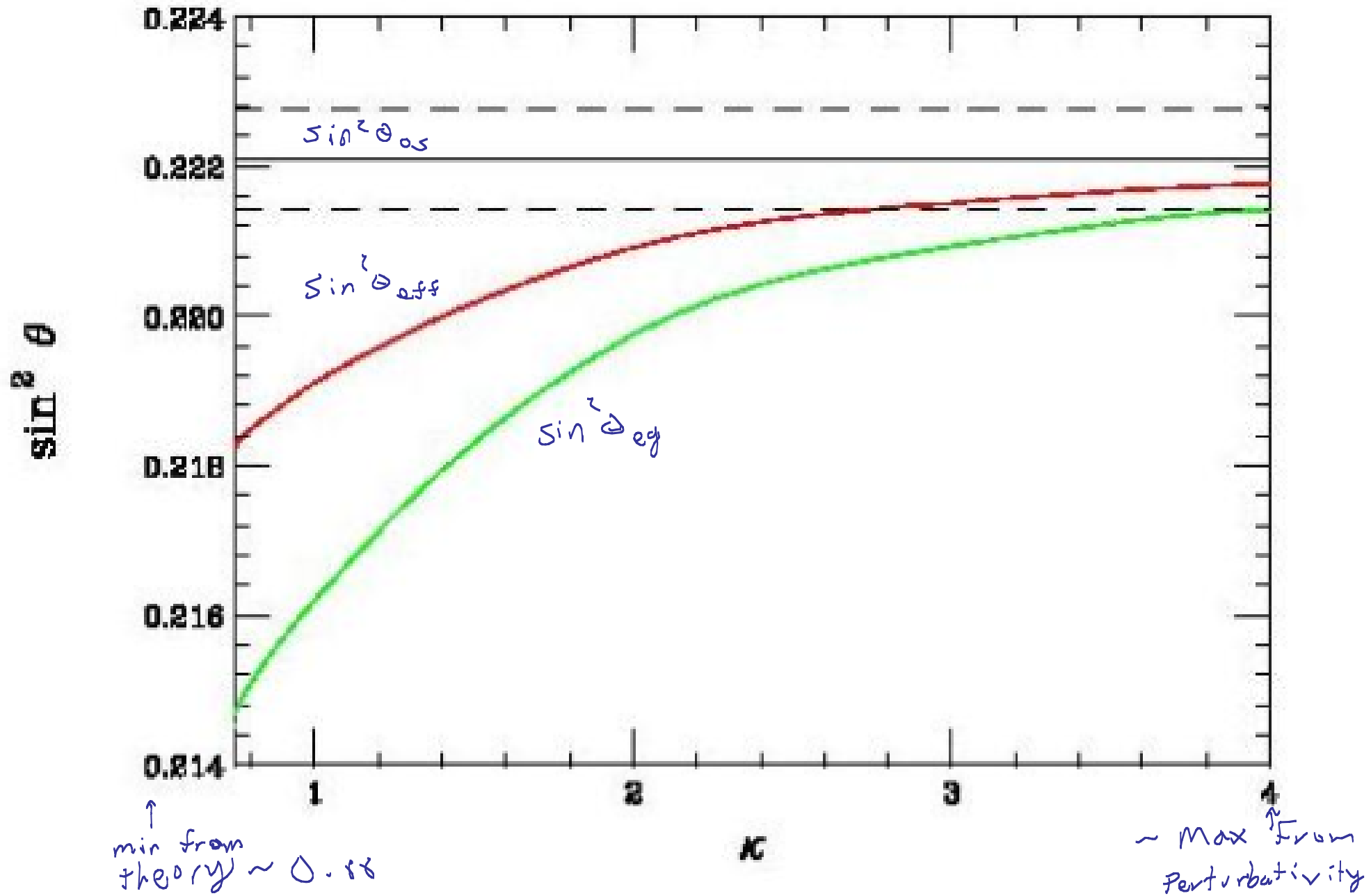
$$\sin^2 \theta_{eq} = \frac{e^2}{g^2}$$

$\sin^2 \theta_{eff}$ ← FROM Z -POLE

MEASURE
DEVIATIONS
OF THESE

ALSO BARBIERI,
POMARILLO + RATTARZI
hep-ph/0310285

BURDMAN + NOMURA
hep-ph/0312247



UNITARITY ISSUES

RECALL \rightarrow SUM RULES ARE VALID AT ASYMPTOTICALLY
HIGH s .

WHAT IF UNITARITY BREAKS BEFORE THIS REGIME?

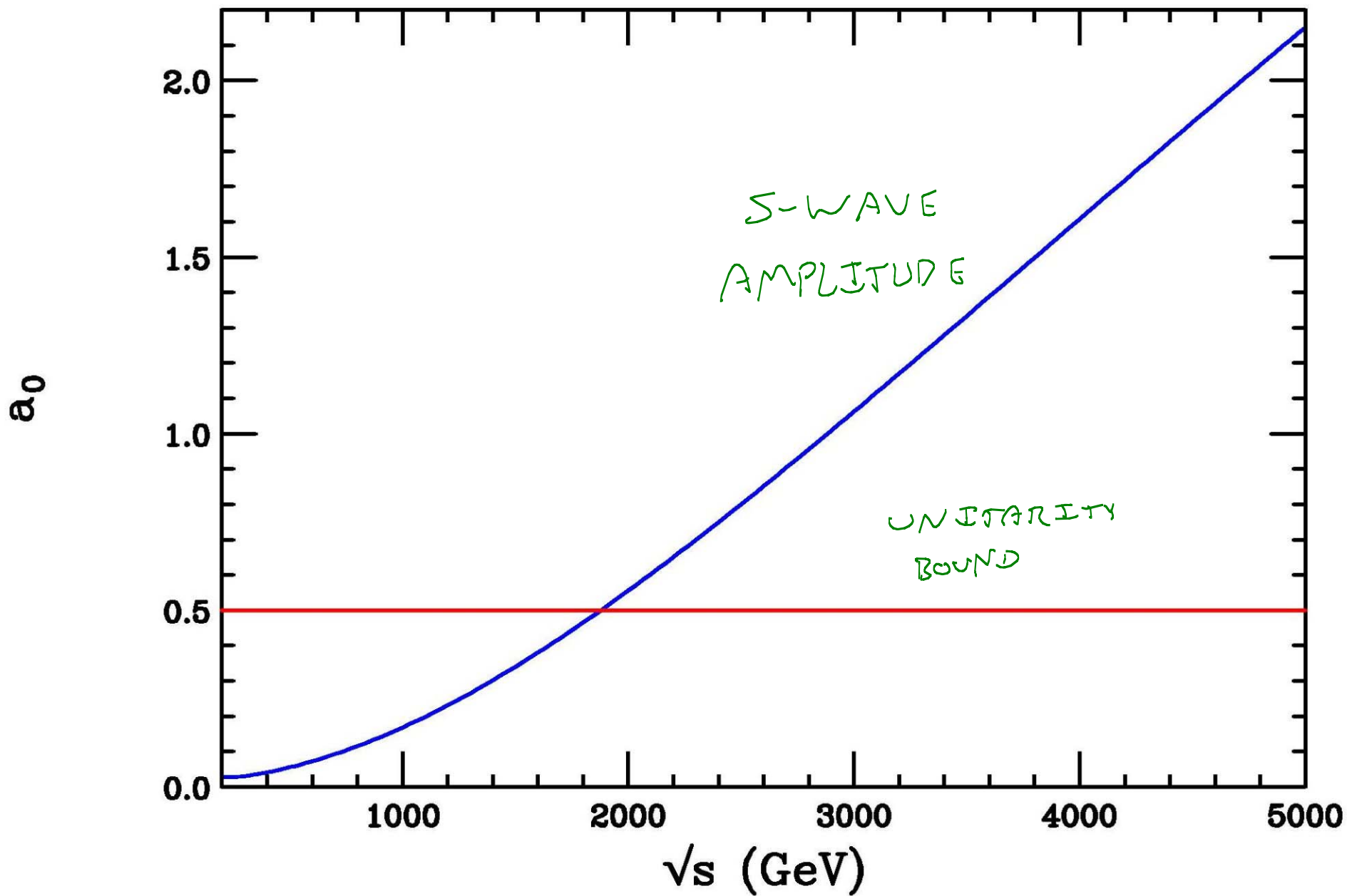
PARTIAL WAVE UNITARITY TEST

$$a_0 = \frac{1}{32\pi} \int_{-1}^1 d\cos\theta A(\omega, \omega \rightarrow \omega, \omega)$$

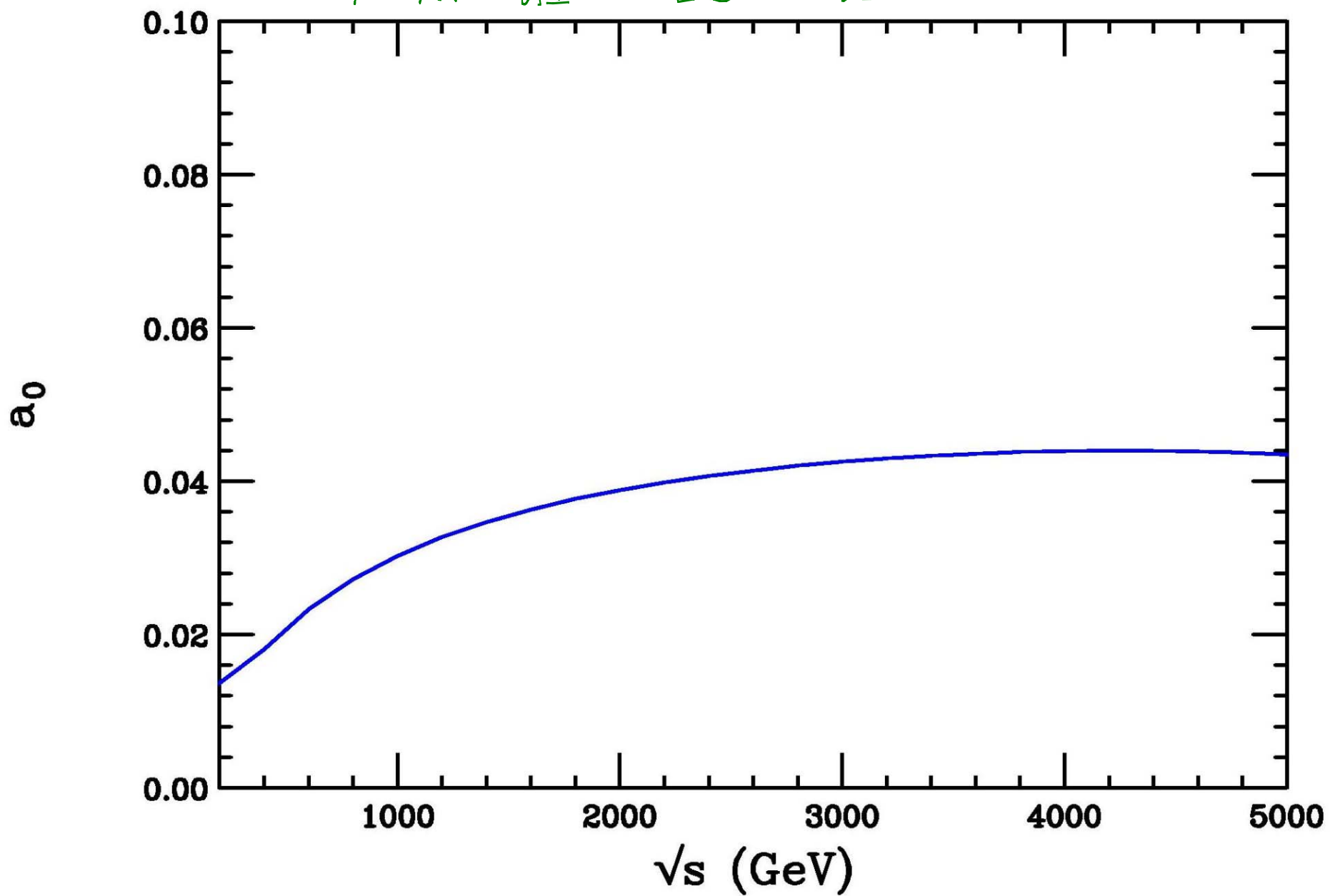
$$|\operatorname{Re}(a_0)| \leq \frac{1}{2}$$

WARPED HIGGSLESS MODEL

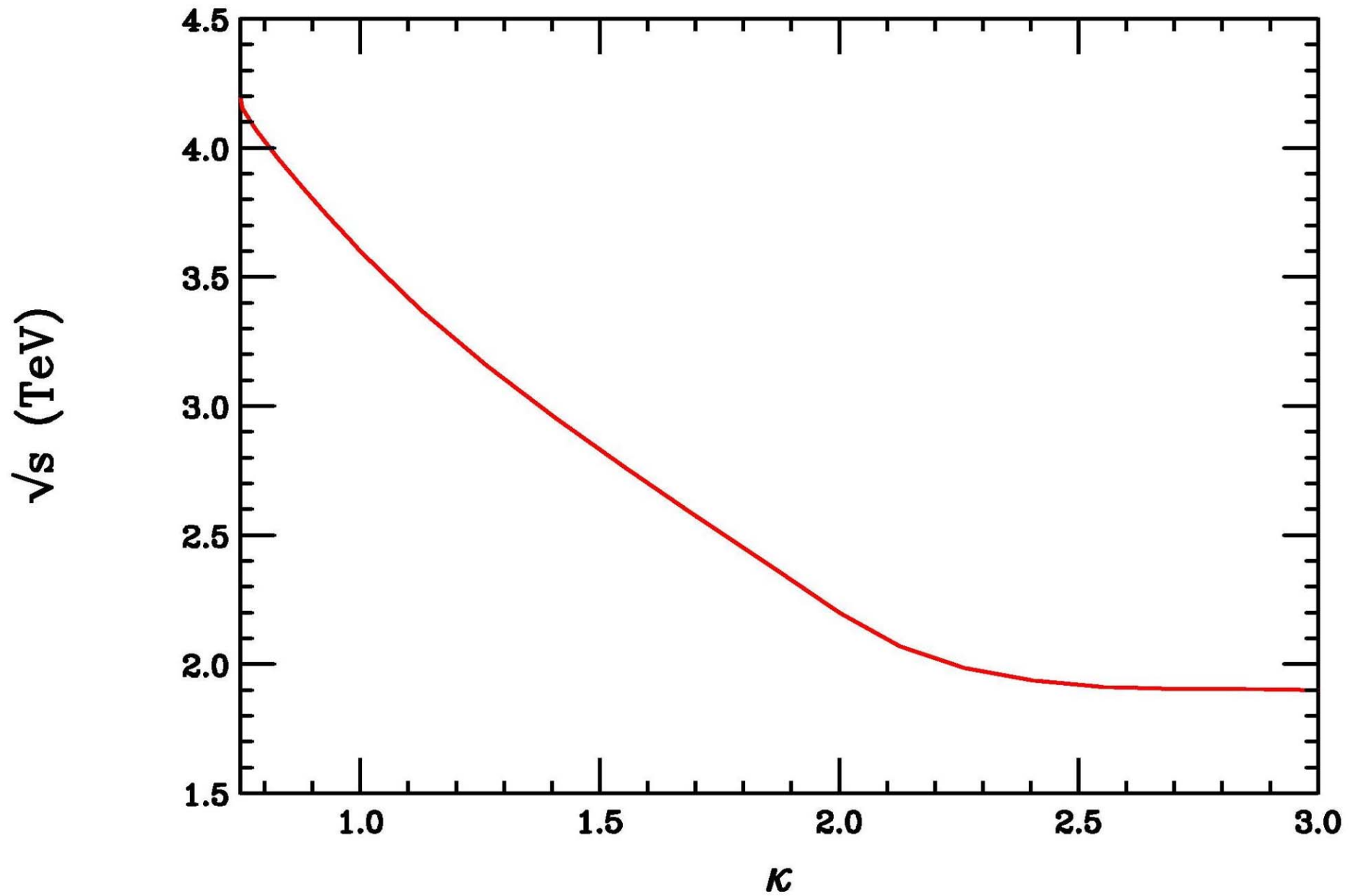
$k=3$



FLAT HIGGSLESS MODEL



ONSET OF UNITARITY VIOLATION



COLLIDER ISSUES

IMPORTANT SIGNATURES

- NO HIGGS SCALAR ← COULD BE A SCALAR
RADEON, FOR EXAMPLE

- RISING $w_L w_L$ SCATTERING ← STUDIED
CROSS SECTION IN GENERAL
TIM BARKLOW e^+e^-

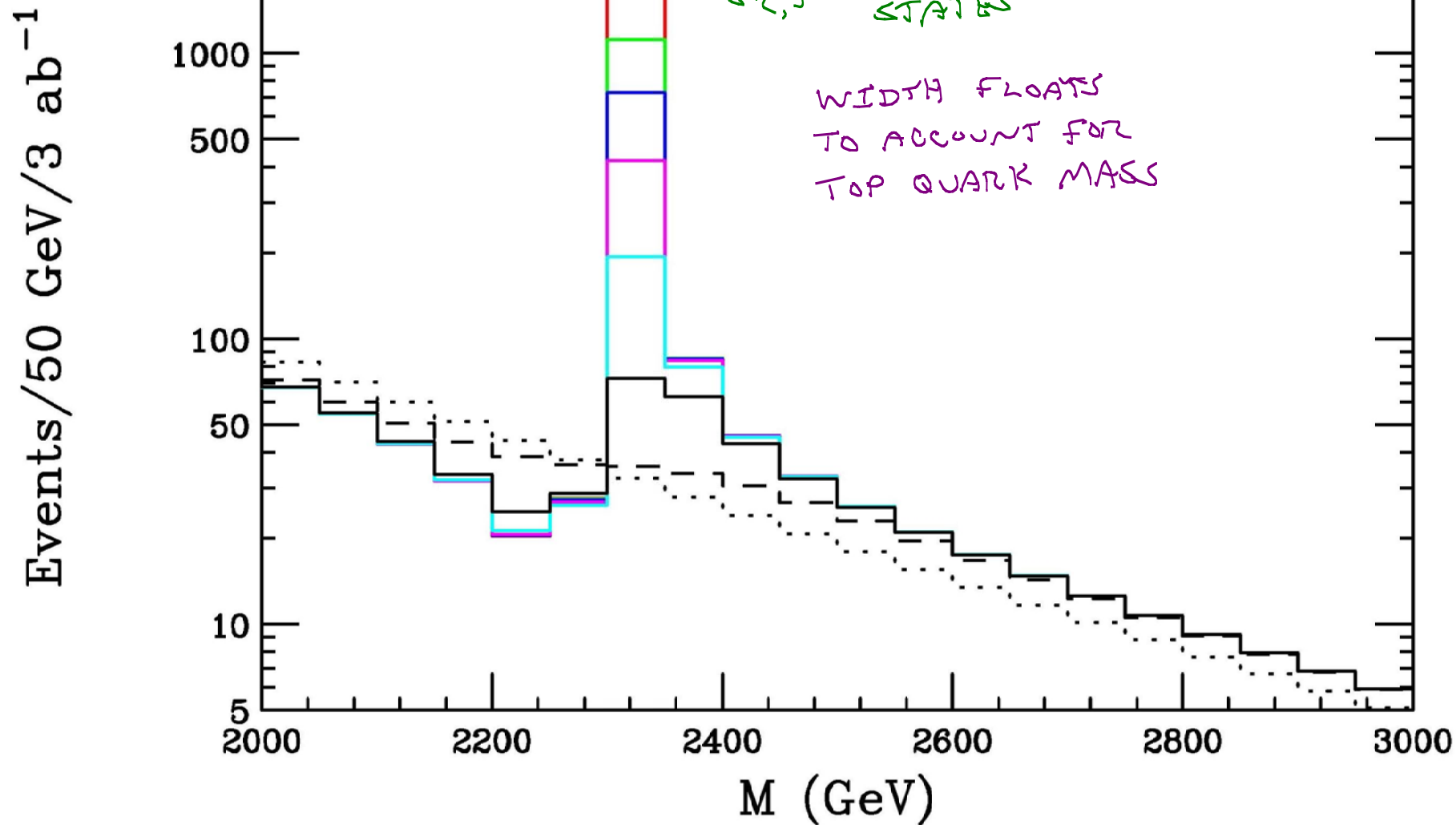
- DOUBLED Z' STATES ← FEATURE OF SEVERAL
EXTRA DIMENSIONAL
MODELS
RIZZO hep-ph/0305077

- GLUON RESONANCES

- GRAVITON RESONANCES ← SMALL, UNLIKE
GENERIC RS

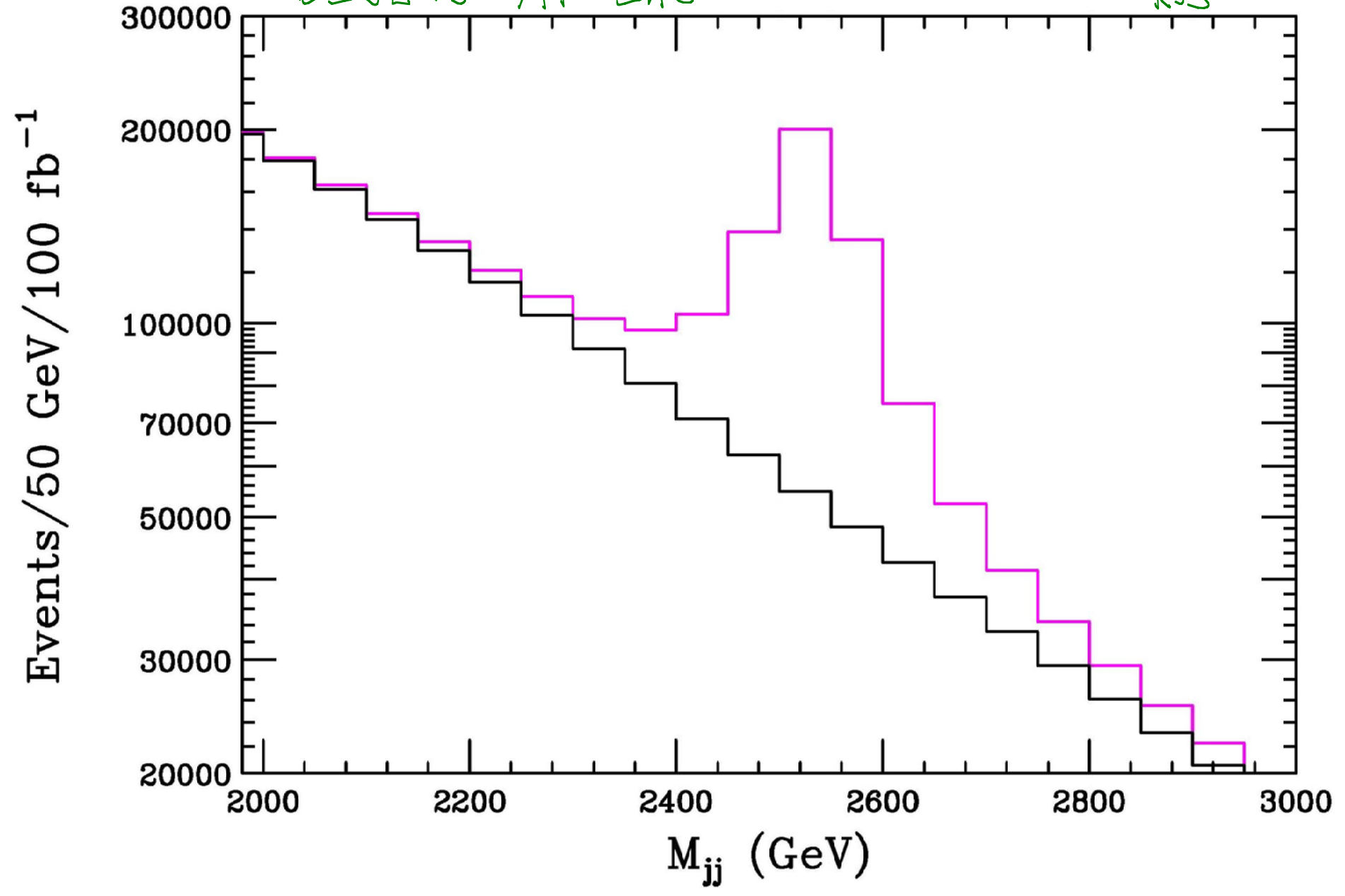
DRELL-YAN PRODUCTION AT LHC

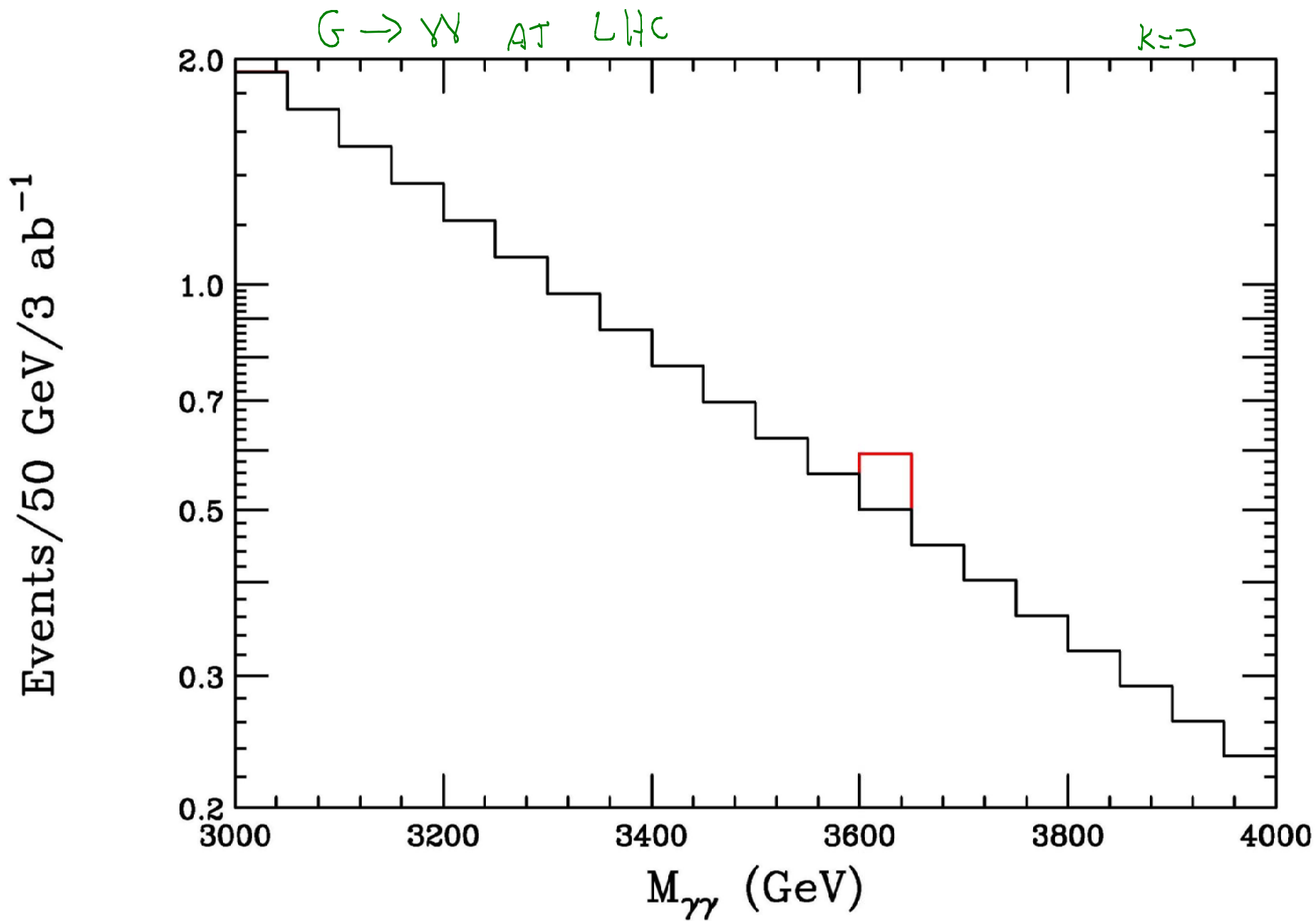
k=3



DIJETS AT LHC

K=3





LINEAR COLLIDER (TO BE DONE)

MOST INTERESTING NEW QUESTION

WHAT CAN THE LC LEARN ABOUT THE
COUPLINGS OF A DOUBLED KK STATE
WHILE COLLIDING BELOW THRESHOLD?

- MASS & (TWO STATES) KNOWN FROM LHC
↑
MAYBE

UNDER CERTAIN CONDITIONS, QUITE A LOT

- $\pi\pi \rightarrow \gamma\gamma$ | 0305077
9612304

CONCLUSIONS

- EXTRA DIMENSIONS MAY PROVIDE AN INTERESTING ALTERNATIVE TO THE HIGGS MECHANISM
- CURRENTLY NO VIABLE MODEL, BUT THE IDEA IS STILL YOUNG (< 1 yr)
- RICH PHENOMENOLOGY THAT NEEDS A LINEAR COLLIDER TO UNTANGLE.

