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MCNPX Model/Table Comparison

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MCNPX MODEL/TABLE COMPARISON

by

John S. Hendricks

ABSTRACT

MCNPX is a Monte Carlo N-Particle radiation transport code extending the capabilities of MCNP4C. As with MCNP, MCNPX uses nuclear data tables to transport neutrons, photons, and electrons. Unlike MCNP, MCNPX also uses (1) nuclear data tables to transport protons; (2) physics models to transport 30 additional particle types (deuterons, tritons, alphas, pions, muons, etc.); and (3) physics models to transport neutrons and protons when no tabular data are available or when the data are above the energy range (20 to 150 MeV) where the data tables end. MCNPX can mix and match data tables and physics models throughout a problem. For example, MCNPX can model neutron transport in a bismuth germinate (BGO) particle detector by using data tables for bismuth and oxygen and using physics models for germanium. Also, MCNPX can model neutron transport in UO₂, making the best use of physics models and data tables: below 20 MeV, data tables are used; above 150 MeV, physics models are used; between 20 and 150 MeV, data tables are used for oxygen and models are used for uranium.

The mix-and-match capability became available with MCNPX2.5.b (November 2002). For the first time, we present here comparisons that calculate radiation transport in materials with various combinations of data charts and model physics. The physics models are poor at low energies (<150 MeV); thus, data tables should be used when available. Our comparisons demonstrate the importance of the mix-and-match capability and indicate how well physics models work in the absence of data tables.

1. MCNPX OVERVIEW

MCNPX¹ is a Monte Carlo N-Particle radiation transport code extending the capabilities of MCNP4C.²

As with MCNP, MCNPX uses nuclear data tables to transport neutrons, photons, and electrons. Unlike MCNP, MCNPX also uses (1) nuclear data tables to transport protons; (2) physics models to transport 30 additional particle types (deuterons, tritons, alphas, pions, muons, etc.); and (3) physics models to transport neutrons and protons when no tabular data are available or when the data are above the energy range (20 to 150 MeV) where the data tables end. The energy range of MCNPX particles is illustrated in Fig. 1.

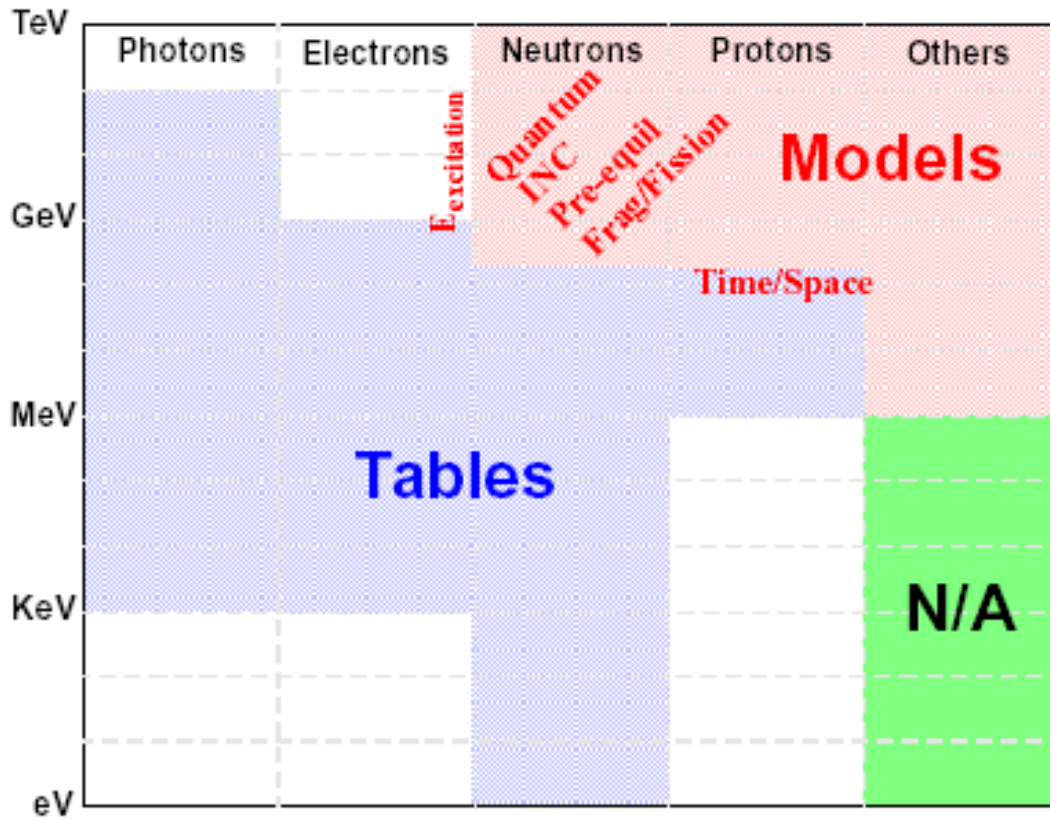


Fig. 1. MCNPX particles and energies.

1.1. MCNPX 2.3.0 Features Not Found in MCNP4C3

MCNPX includes all of the features and capabilities of MCNPX 2.3.0³ (based on MCNP4B), which was released recently to the Radiation Safety Information Computational Center (RSICC). The RSICC⁴ is located in Oak Ridge, Tennessee, and is the principal distributor of MCNP, MCNPX, and other Department of Energy (DOE) computer codes. The following features are not found in MCNP4C3:

- physics for 34 particle types;
- high-energy physics above the tabular data range;
- photonuclear physics;⁵
- neutron, proton, and photonuclear 150-MeV libraries and utilization;
- mesh tallies (tally fluxes, heating, sources, etc., in a superimposed mesh);
- radiography tallies;
- secondary-particle production biasing; and
- autoconfiguration build system for compilation.

In particular, the mesh tallies enable users to plot source particle locations, fluxes, energy deposition, particle tracks, DXTRAN contributions, and other useful quantities on a superimposed tally grid. The mesh tally capability is useful for all transport problems, not just high-energy problems.

1.2. MCNP4C3 Features Not Found in Earlier MCNPX Versions

MCNPX includes all of the features and capabilities of MCNP4C3. Until MCNPX2.4.0 (August 2002), MCNPX was based on MCNP4B (March 1997).⁶ MCNPX now has all the capabilities of MCNP4C^{7,8} (March 2000), MCNP4C2⁹ (January 2001), and MCNP4C3¹⁰ (April 2001). The initials of principal developers are shown in parentheses.* The new features previously unavailable in MCNPX are

- PC enhancements: MCNPX is fully Linux and Windows capable (LLC/GWM);
- easier geometry specification with macrobodies (LLC);
- interactive geometry plotting (JSH);
- improved variance reduction with the superimposed mesh weight window generator (TJE/JAF/JSH);
- superimposed mesh plotting (JSH);
- delayed neutrons (CJW);
- unresolved resonance range probability tables (LLC/RCL);
- perturbations for material-dependent tallies (GWM/LLC/JSH);
- ENDF/B-VI extensions (MCW);
- electron physics enhancements (upgrade to ITS3.0¹¹) (KJA/HGH);
- weight window enhancements (JSH/JAF); and
- distributed memory multiprocessing (GWM).

1.3. MCNPX Capabilities Provided after the MCNP4C/MCNPX Merger

MCNPX has many new capabilities not found in either MCNP4C or MCNPX 2.3.0.¹² The following features became available with MCNPX2.4.0,^{13,14} which was released to the RSICC on August 1, 2002:

- FORTRAN 90 modularity and dynamic memory allocation (GWM);
- distributed memory multiprocessing for the entire energy range of all particles (GWM);
- repeated structures source path improvement (LLC/JSH);
- default dose functions (LSW/JSH);
- light-ion recoil (JSH);
- enhanced color geometry plots (GWM/JSH);
- photonuclear cross-section plots (JSH);
- proton cross-section plots (JSH);

* Gregg W. McKinney (GWM), John S. Hendricks (JSH), Laurie S. Waters (LSW), Leland L. Carter (LLC), Franz X. Gallmeier (FXG), H. Grady Hughes (HGH), Richard E. Prael (REP), Stepan G. Mashnik (SGM), Arnold J. Sierk (AJS), Thomas J. Evans (TJE), Jeffrey A. Favorite (JAF), Christopher J. Werner (CJW), Robert C. Little (RCL), Morgan C. White (MCW), and Kenneth J. Adams (KJA).

- proton reaction multipliers with FM cards (JSH);
- photonuclear reaction multipliers with FM cards (JSH/GWM);
- some speedups (GWM/JSH);
- logarithmic interpolation on input cards (JSH);
- cosine bins that may be specified in degrees (JSH);
- cosine bins that may be specified for F2 flux tallies (JSH);
- source particles that may be specified by descriptors (JSH);
- pause command for tally and cross-section plots (JSH); and
- correction of all known MCNPX and MCNP4C bugs/problems (REP/GWM/FXG/HGH/JSH).

The following additional features became available in MCNPX2.5.b (November 2002):¹⁵

- CEM2k physics^{16,17} (SGM/AJS/FXG);
- mix and match (JSH);
- positron sources (HGH); and
- spontaneous fission (JSH).

1.4. MCNPX Mix-and-Match Capability

MCNPX can mix and match data tables and physics models throughout a problem. For example, MCNPX can model neutron transport in a bismuth germanate (BGO) particle detector by using data tables for bismuth and oxygen and using physics models for germanium. Also, MCNPX can model neutron transport in UO₂, making the best use of physics models and data tables: below 20 MeV, data tables are used; above 150 MeV, physics models are used; between 20 and 150 MeV, data tables are used for oxygen and models are used for uranium.

This report provides comparisons that calculate radiation transport in materials with various combinations of data tables and model physics. The physics models are poor at low energies (<150 MeV); thus, data tables should be used when available. Our comparisons demonstrate the importance of the mix-and-match capability and indicate how well physics models work in the absence of data tables.

2. BGO COMPARISON

The first example of the mix-and-match capability is a BGO detector crystal. An unambiguous description of this example problem is the MCNPX input file provided in the appendix. The crystal is 8.433 cm long and 3.932 cm in radius with a density of 7.13 g/cm³. A monodirectional 100-MeV beam of neutrons impinges on the base of the BGO can, and we tally and plot the neutron and photon leakage, flux, and heating.

The BGO consists of 21% bismuth, 16% germanium, and 63% oxygen. Currently, no germanium nuclear data table is available for either MCNP or MCNPX. With mix and match, the 150-MeV data tables may be used for bismuth and oxygen and a physics

model can be used for germanium. Before the mix-and-match capability, the only way to model the BGO detector was to substitute some other nuclide for germanium, such as arsenic. The available arsenic nuclear data table goes up to 20 MeV, leaving users with two poor choices:

1. use models for all neutron nuclides above 20 MeV, neglecting the 20- to 150-MeV data for bismuth and oxygen; or
2. use models using >150 MeV to take advantage of the 150-MeV bismuth and oxygen data tables but then use the 20-MeV values of the arsenic cross sections in the 20- to 150-MeV range.

In either case, the model for arsenic would have to be used instead of germanium above the arsenic data table energy range. In all cases, data tables are used for the photons.

Figures 2 through 7 show the BGO leakage (penetration through the top of the crystal), flux, and heating for neutrons and photons. In each case, the following are plotted:

- mix/match: data tables are used for bismuth and oxygen, and a physics model is used for germanium.
- with arsenic at 20 MeV (W/As @20): data tables up to 20 MeV are used for bismuth, oxygen, and arsenic (substitute for germanium); above 20 MeV, physics models are used for bismuth, arsenic, and oxygen.
- with arsenic at 150 MeV (W/As @150): data tables are used for bismuth and oxygen up to 20 MeV (above the problem energy range). The arsenic data table is used up to 20 MeV. The 20-MeV arsenic cross section is used above 20 MeV.

These figures illustrate the difference between the new mix-and-match capability in which a model for germanium can be used and the capability before mix and match where Table 1 data had to be either extended or ignored in the mix-and-match energy range.

It is evident that significant differences can be found in the answers, depending on whether a germanium model or arsenic table data below 20 MeV are used and depending on the treatment of arsenic above 20 MeV (either using the model or table extension with 20-MeV values).

The physics models are not as good as evaluated data. Below 20 MeV, they become unreliable. At ~1 MeV, where they are cut off, these models are highly questionable. In particular, the CEM2k physics model is poor below 7 MeV. Whether using a physics model for germanium is better or worse than substituting an arsenic data table will not be known until there is a good data table for germanium. Certainly, extrapolating the 20-MeV data above 20 MeV for arsenic is a very poor approximation.

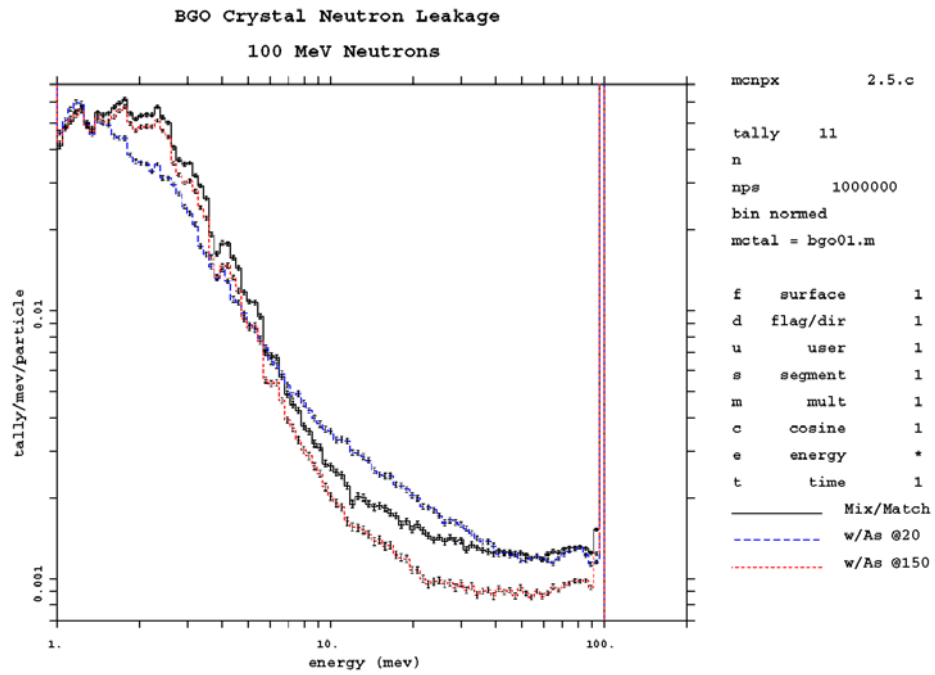


Fig. 2. BGO neutron leakage, new vs old capability.

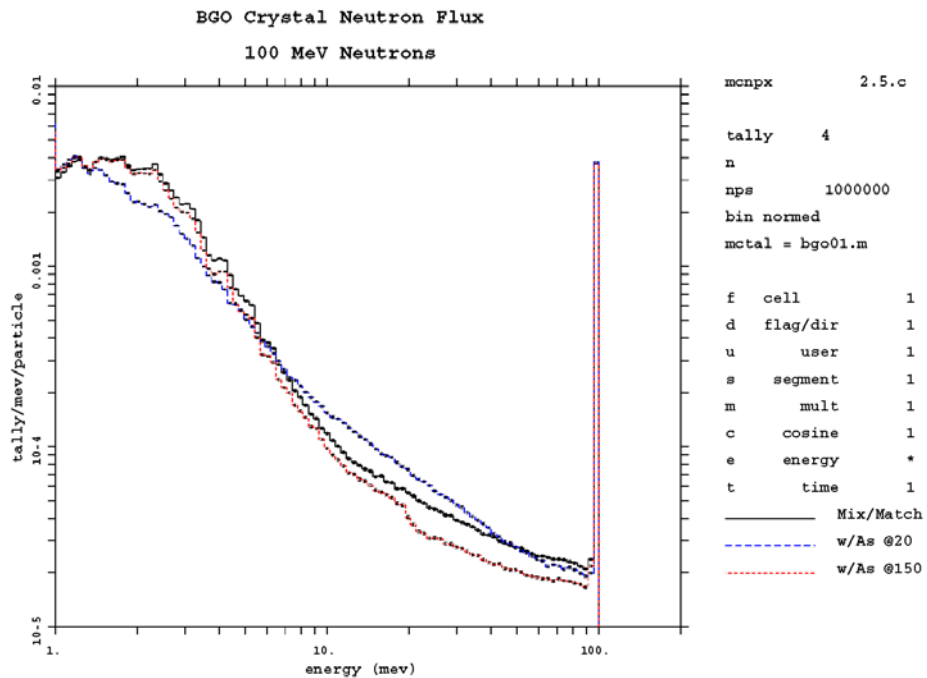


Fig. 3. BGO neutron flux, new vs old capability.

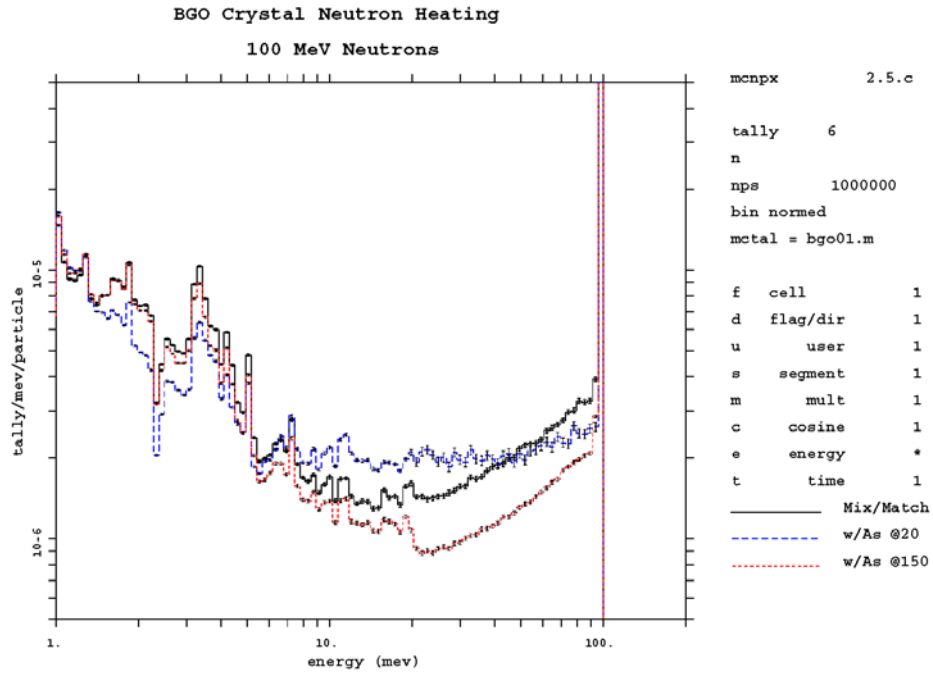


Fig. 4. BGO neutron heating, new vs old capability.

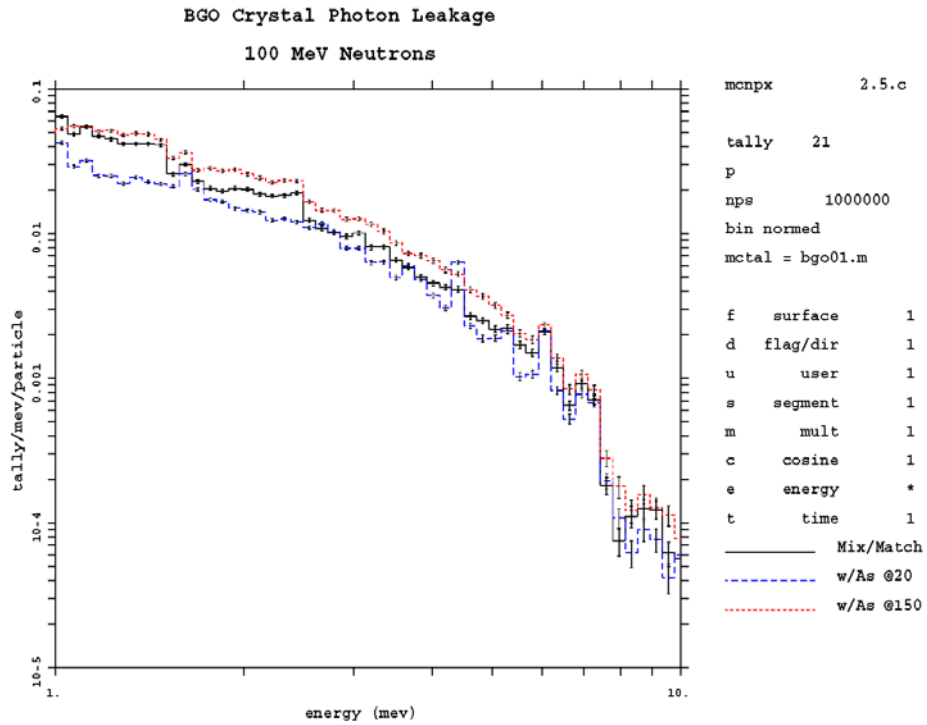


Fig. 5. BGO photon leakage, new vs old capability.

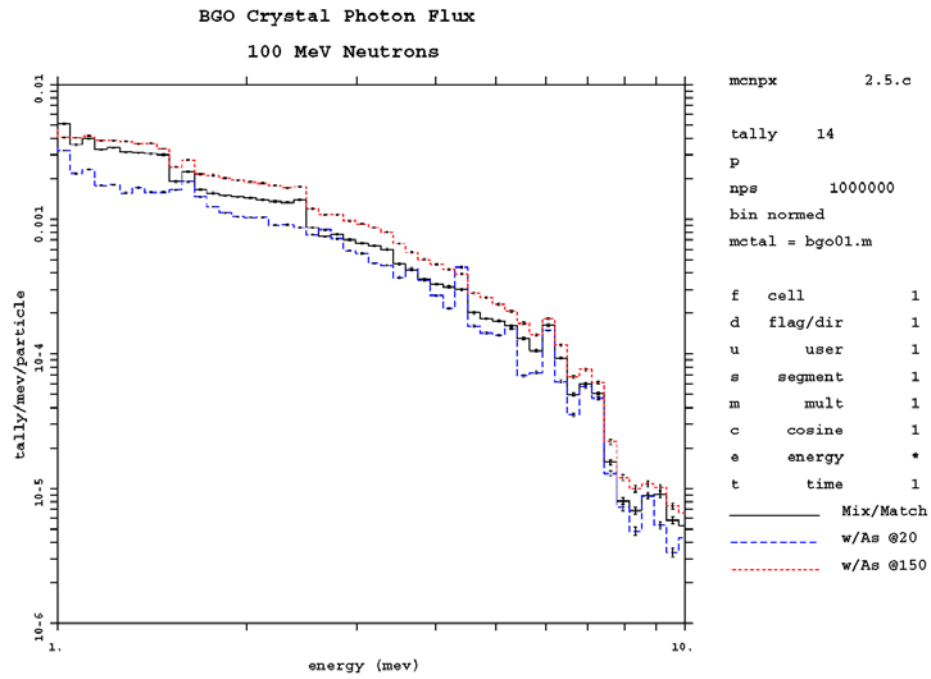


Fig. 6. BGO photon flux, new vs old capability.

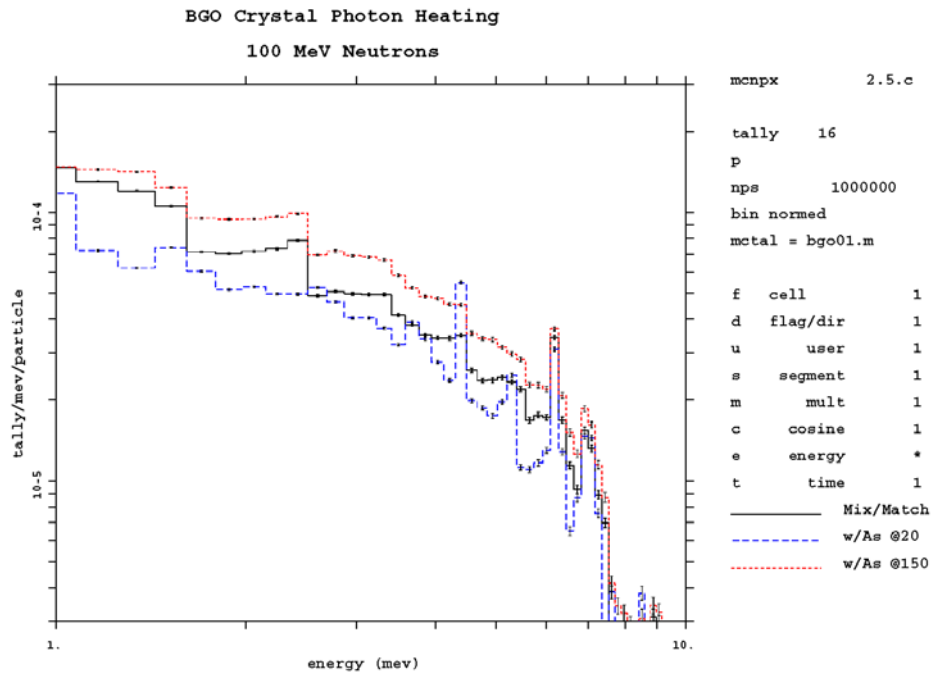


Fig. 7. BGO photon heating, new vs old capability.

Table 1
BGO Crystal Integral Quantities

Quantity	Mix/Match	W/As @20 MeV	W/As @150 MeV
Neutron Leakage	9.9287E-01	9.8895E-01	9.0710E-01
Neutron Flux	3.5566E-02	3.4562E-02	3.3783E-02
Neutron Heating	3.1818E-03	2.8093E-03	2.5052E-03
Photon Leakage	1.1231E-01	8.7039E-02	1.2856E-01
Photon Flux	8.9828E-03	6.9487E-03	1.0474E-02
Photon Heating	4.6665E-04	3.7081E-04	5.5945E-04

Using an arsenic model instead of germanium is a crude approximation. Using arsenic photon table data rather than germanium whenever arsenic is substituted for the neutron table is also a poor approximation.

An indication of how poor the neutron physics model is at low energies can be found by comparing the default model, CEM2k model, and data table for arsenic. Figures 8 through 13 show the arsenic leakage, flux, and heating for neutrons and photons for a 20-MeV neutron source. The crystal density and dimensions are unchanged.

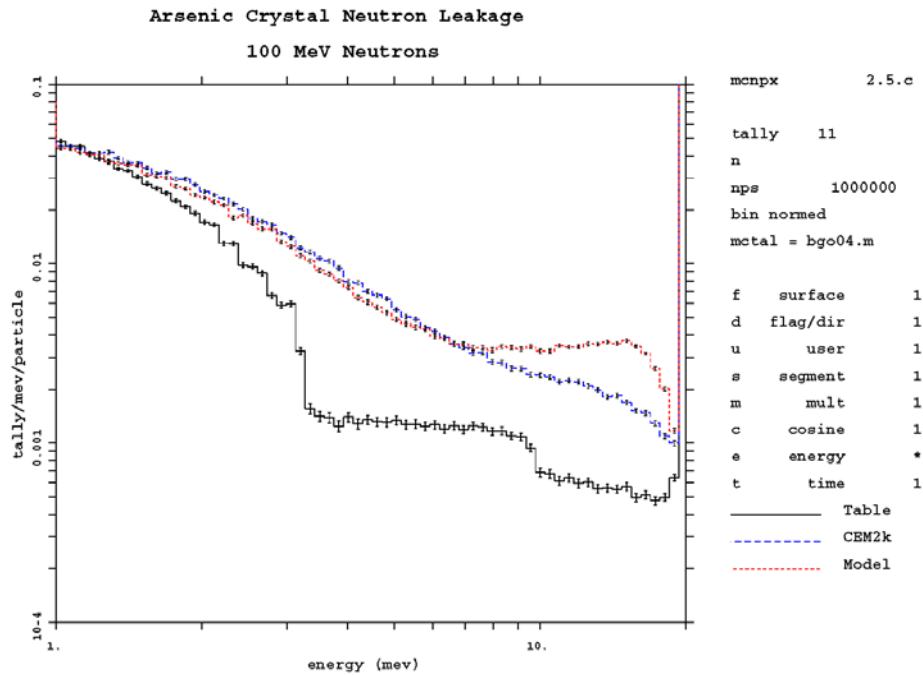


Fig. 8. Arsenic neutron leakage, table vs model.

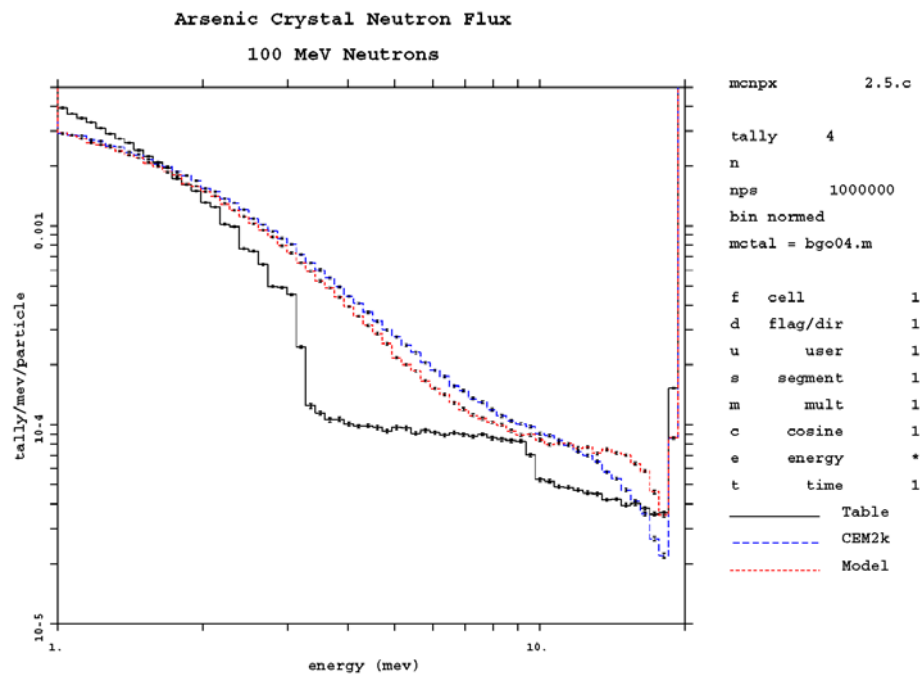


Fig. 9. Arsenic neutron flux, table vs model.

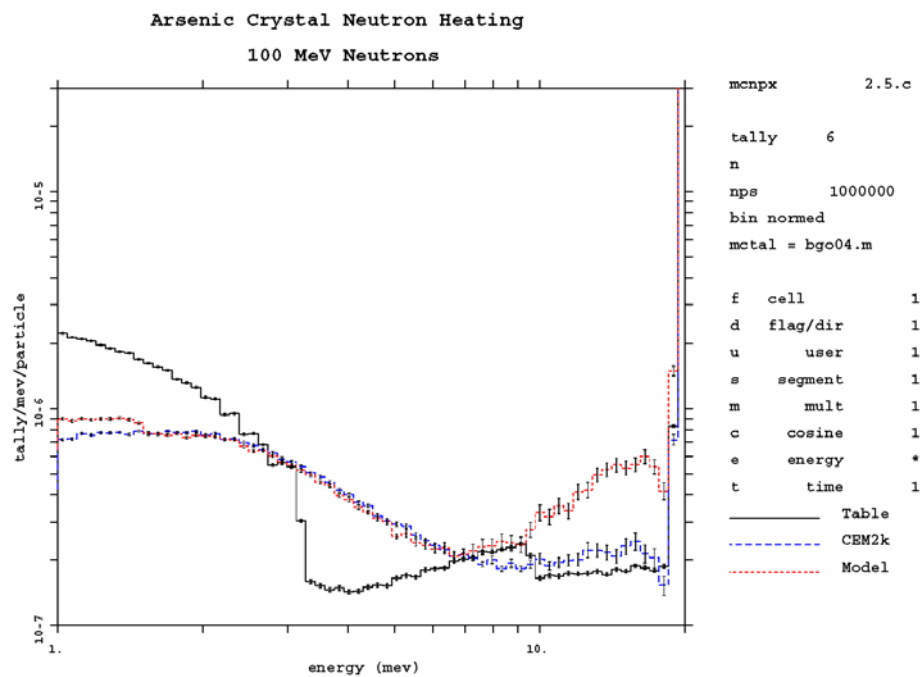


Fig. 10. Arsenic neutron heating, table vs model.

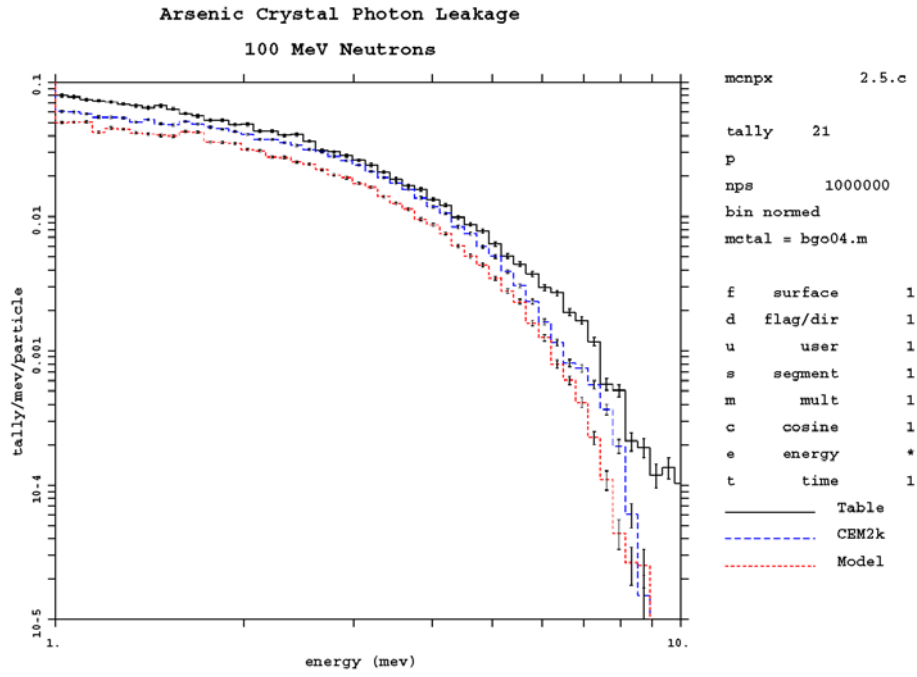


Fig. 11. Arsenic photon leakage, table vs model.

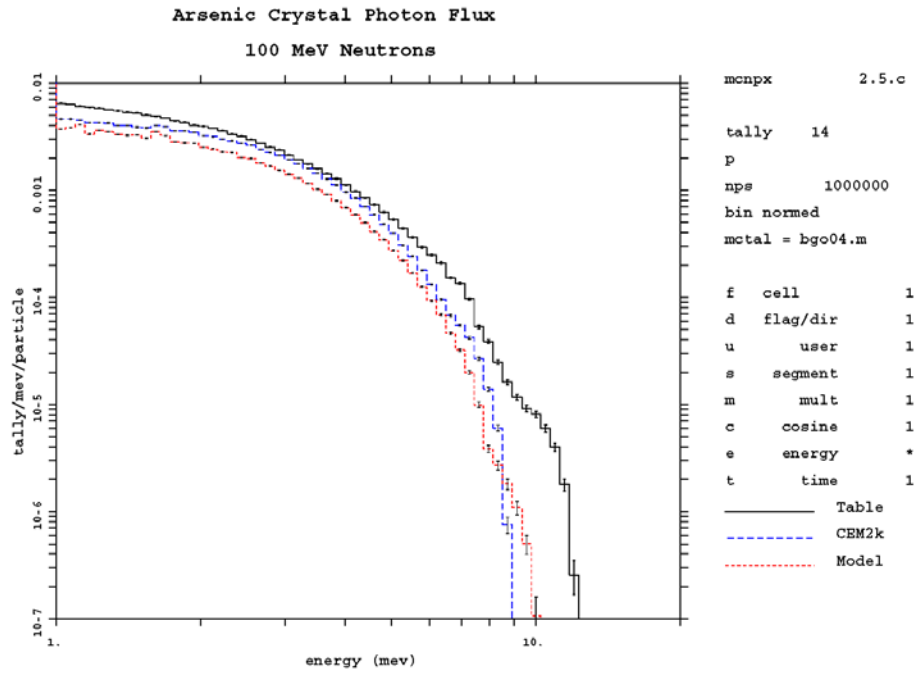


Fig. 12. Arsenic photon flux, table vs model.

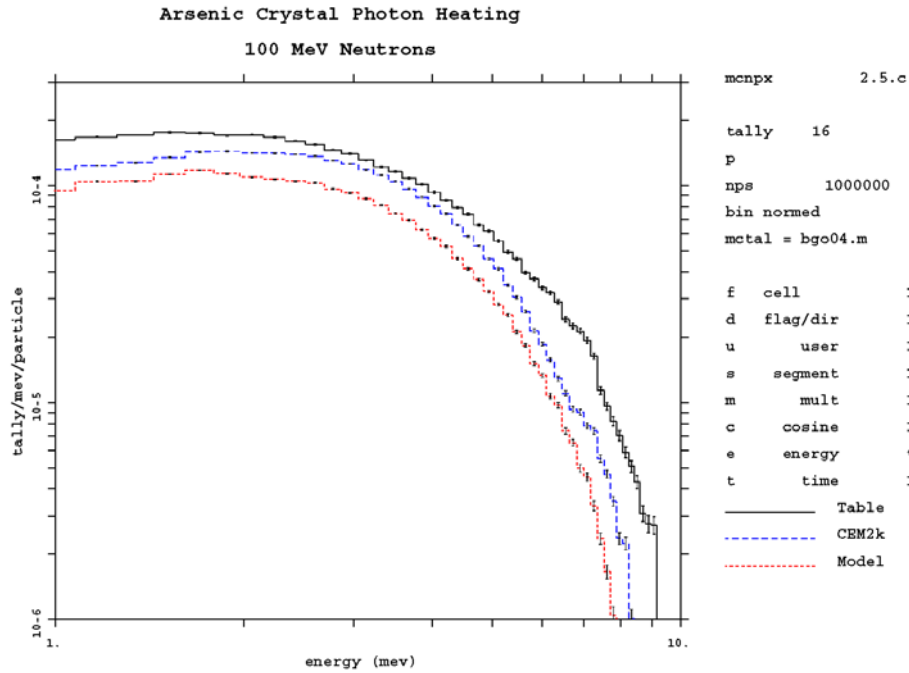


Fig. 13. Arsenic photon heating, table vs model.

The arsenic crystal comparison plots illustrate the serious deficiencies of models at low energies. All neutron heating is dumped at 20 MeV with the physics models. The neutron flux and leakage are an order of magnitude higher for the models than for the arsenic data table below 20 MeV. Apparently, the LAHET portion of the MCNPX model physics has built-in assumptions that neutron cross sections will not be used below 20 MeV, and in some, or perhaps all cases, the cross section is set to zero.

For photons, the same photoatomic data table is used in all three cases: neutron table, neutron CEM2k, and neutron default model. The differences between the various photon plots are caused by scaling due to the neutron production of photons. Evidently, the models produce about an order-of-magnitude fewer photons.

Although the arsenic crystal plots make the models look poor, the difference between the data table and physics models is much smaller in an integral sense. Table 2 shows the arsenic crystal data integrated over energy.

How much of the difference between the data table and the models is due to a bad arsenic data table? Although arsenic is probably not as well characterized as more important nuclides, the quality of the arsenic table data is suggested by its cross-section plot in Figs. 14 and 15, which show the total (mt = -1), absorption (mt = -2), and elastic (mt = -3) table cross section over the entire energy range and over the 1- to 20-MeV energy range of the arsenic crystal problem.

Table 2
Arsenic Crystal Integral Quantities, 20-MeV Source

Quantity	Data Table	CEM2k	Model
Neutron Leakage	5.7065E-01	6.0530E-01	6.4494E-01
Neutron Flux	2.5913E-02	2.3656E-02	2.5626E-02
Neutron Heating	9.0549E-05	1.5594E-04	2.5703E-04
Photon Leakage	2.7146E-01	2.4545E-01	2.0415E-01
Photon Flux	2.4901E-02	2.2346E-02	1.8625E-02
Photon Heating	8.3046E-04	7.0777E-04	5.5700E-04

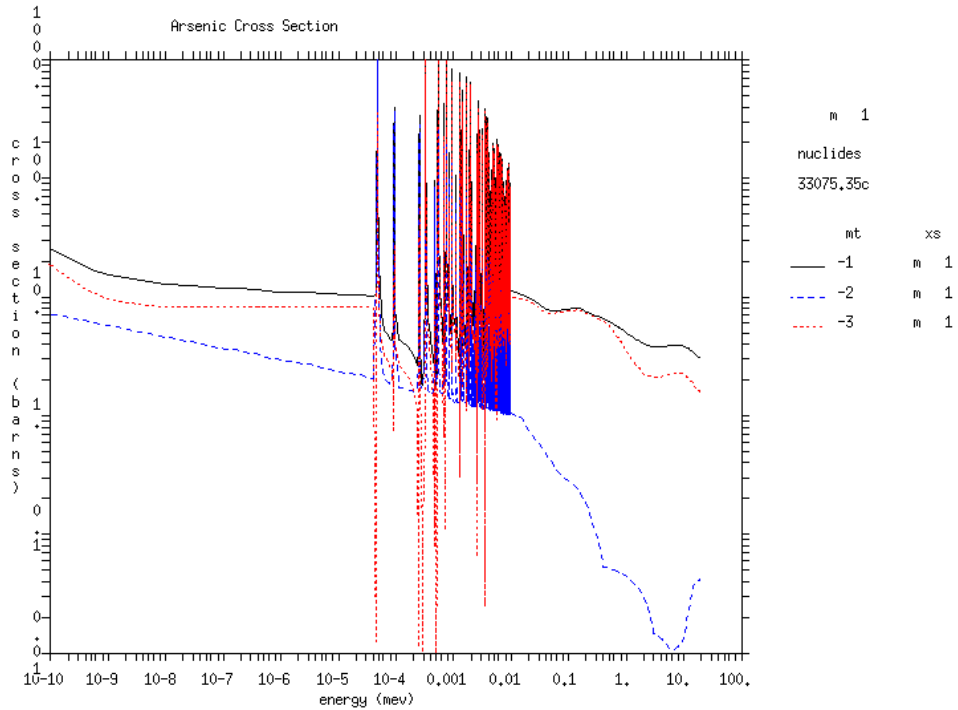
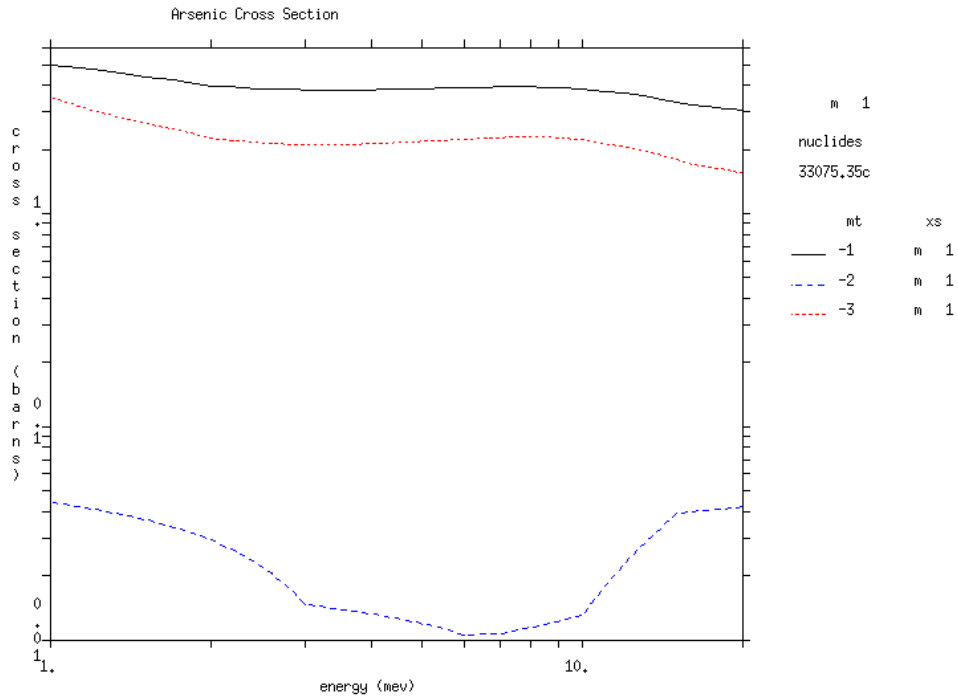


Fig. 14. Arsenic cross section, 0 to 20 MeV.



*Where the reaction MT numbers -1, -2, and -3 shown in the right margin are the total, absorption, and elastic cross sections, respectively.

Fig. 15. Arsenic cross section, 1 to 20 MeV.*

The BGO results demonstrate that the MCNPX mix-and-match capability enables the use of physics models where no data tables are available. Fortunately, neutron data tables are available for most nuclides. Photon and electron data tables are available for all elements and thus for all nuclides.

3. WATER COMPARISON

The worst-case scenario for models is light nuclei at low energy. Water is a very important material composed of light nuclei. The water comparison considers neutron table data and physics models for a 20-cm-radius sphere of water. Leakage, flux, and heating are calculated for both a 150-MeV neutron source and a 2-MeV neutron source. These results are presented in Figs. 16 to 21 and Tables 3 and 4.

The 150-MeV neutron source in the center of the 20-cm-radius water sphere shows fairly decent agreement between using data tables and physics models. Note that the neutron heating jumps when the data tables are used. The 150-MeV neutron data tables were constructed by adding 20- to 150-MeV data to the 20-MeV data tables without correcting for discontinuities at 20 MeV.

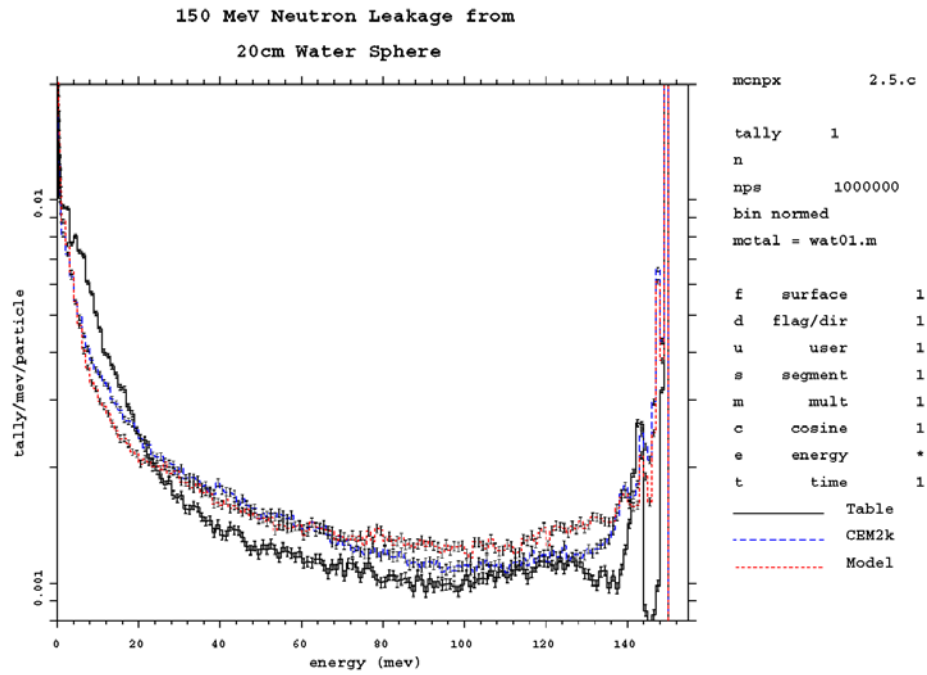


Fig. 16. The 150-MeV neutron water leakage, table vs model.

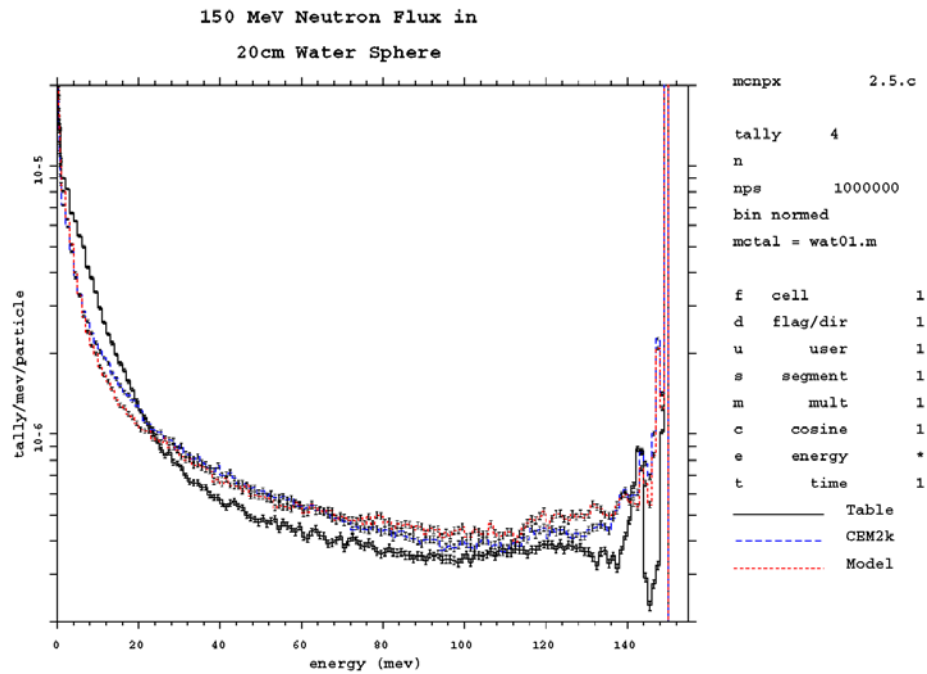


Fig. 17. The 150-MeV neutron water flux, table vs model.

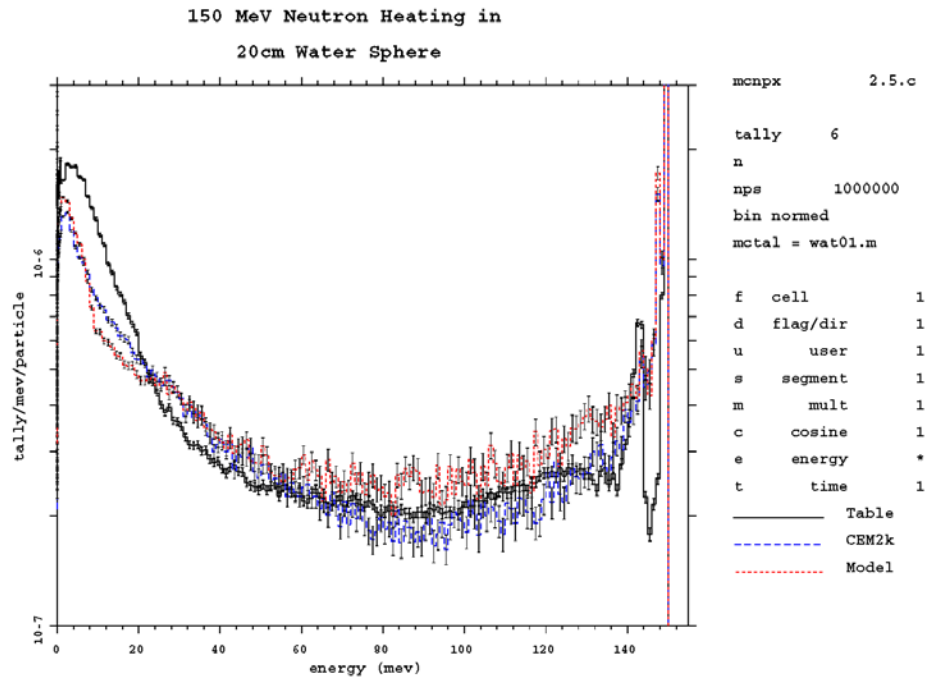


Fig. 18. The 150-MeV neutron water heating, table vs model.

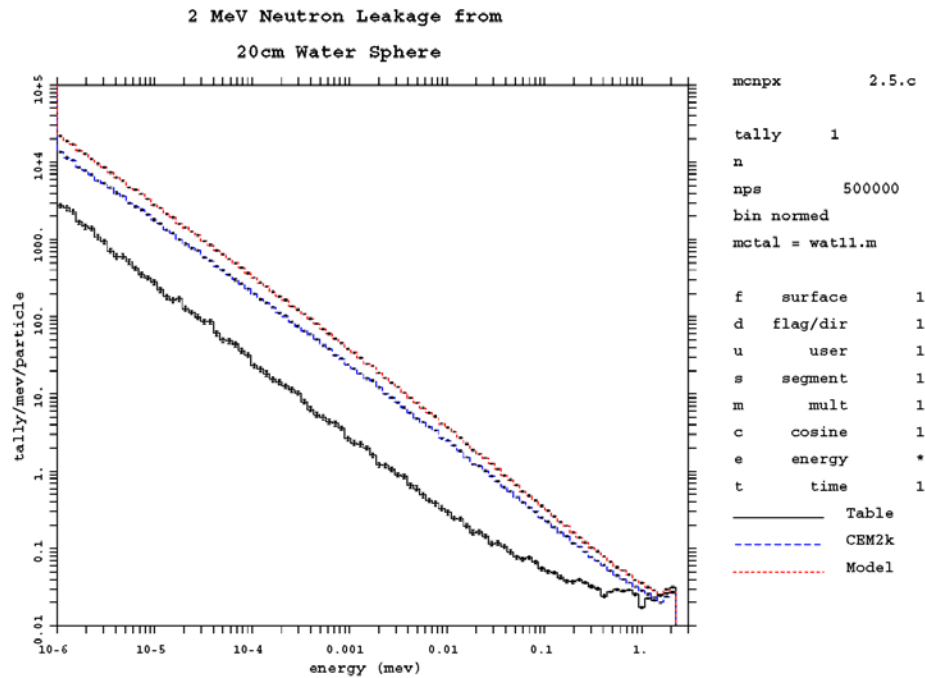


Fig. 19. The 2-MeV neutron water leakage, table vs model.

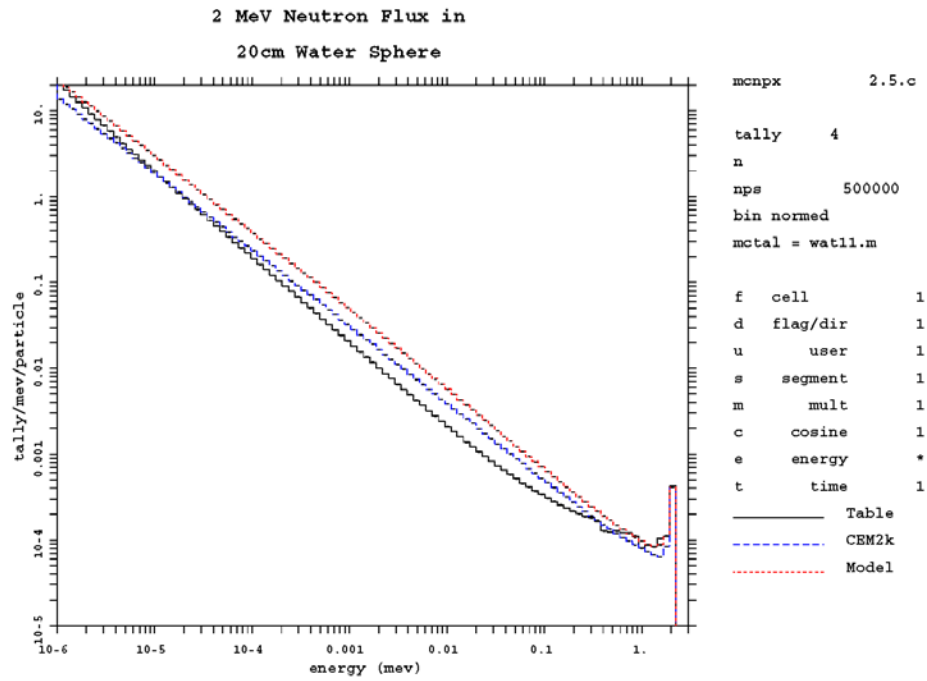


Fig. 20. The 2-MeV neutron water flux, table vs model.

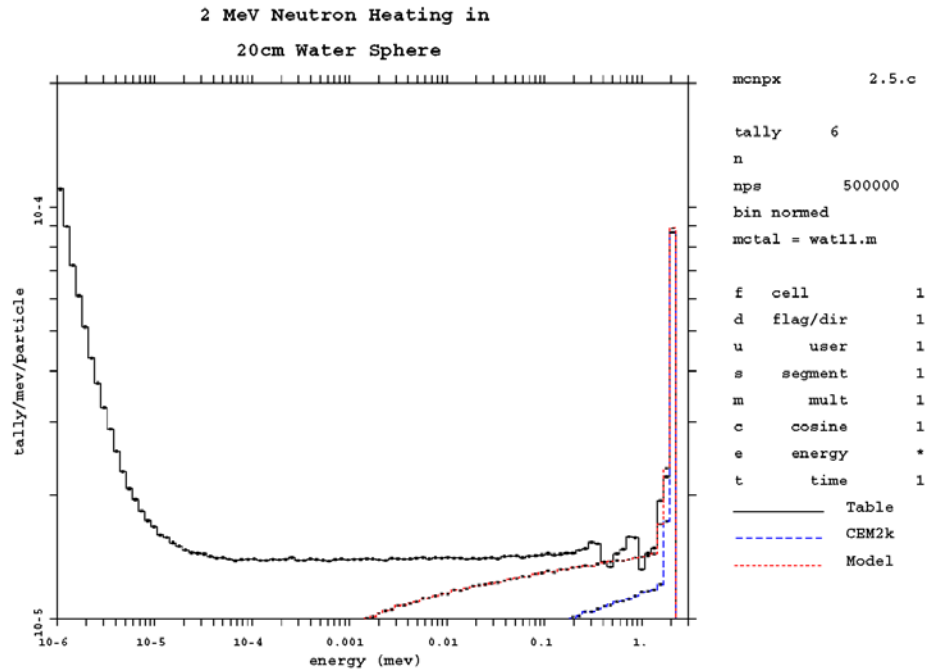


Fig. 21. The 2-MeV neutron water heating, table vs model.

Table 3
150-MeV Neutrons in the 20-cm-Radius Water Sphere

Quantity	Data Table	CEM2k	Model
Neutron Leakage	1.0662E+00	1.0781E+00	1.0968E+00
Neutron Flux	8.1373E-04	6.9696E-04	7.1651E-04
Neutron Heating	4.7144E-04	4.0048E-04	4.6065E-04

Table 4
2-MeV Neutrons in the 20-cm-Radius Water Sphere

0 to 2 MeV

Quantity	Data Table	CEM2k	Model
Neutron Leakage	1.7597E-01	4.4548E-01	6.7341E-01
Neutron Flux	1.8921E-03	7.7193E-04	1.0697E-03
Neutron Heating	5.7769E-05	4.9643E-05	5.7015E-05

1 eV to 2 MeV

Quantity	Data Table	CEM2k	Model
Neutron Leakage	9.2692E-02	3.3109E-01	4.8552E-01
Neutron Flux	5.9237E-04	6.6370E-04	8.9802E-04
Neutron Heating	5.7717E-05	4.9643E-05	5.7015E-05

The integral results for the 20-MeV neutron source in the 20-cm-radius water sphere are shown in Table 3. The physics models differ from the data tables by 10% to 20% in the range of 20 to 150 MeV.

Integral results for the 2-MeV neutron source in the 20-cm-radius water sphere are given in Table 4. Results are presented for the ranges of both 0 MeV to 20 MeV and 1 eV to 20 MeV. It is evident that the low-energy results strongly affect the comparison. Certainly, physics models are known to be very poor at low energies, and the agreement between the tables and models in this worst-case scenario is impressive: energy deposition agrees to 15%, fluxes above 1 eV agree to within 50%, and leakage is off by only a factor of 5.

Light nuclei at low energies is the worst-case scenario for substituting neutron physics models for data tables. The 150-MeV neutron source in the 20-cm-radius sphere of water showed 10% to 20% differences between data tables and physics models. The 2-MeV source showed a factor of 5 difference for leakage, 50% difference for flux, and only 15% difference for heating.

4. STEEL COMPARISON

Stainless steel is another critical material. Our first stainless-steel problem consists of a 200-MeV monodirectional neutron beam impinging on 10-cm-thick steel. The second problem is a 20-MeV monodirectional neutron beam impinging on 1-cm-thick steel. The 8-g/cm³ steel isotopic composition is presented in the MCNPX input deck in the appendix. Calculated quantities include penetration, backscatter, flux, and heating.

Note that ¹⁵N, ³²S, ⁵⁵Mn, and ⁵⁸Fe do not have neutron data tables for 150 MeV. The mix-and-match capability makes it possible to use physics models for these nuclides above 20 MeV. In Figs. 22 to 25, the mix-and-match results are compared with results using (1) a 20-MeV cutoff between table and model regions; and (2) a 150-MeV cutoff between table and model regions, which was required before mix and match. The 20-MeV cutoff ignores the table data for most of the stainless-steel nuclides above 20 MeV. The 150-MeV cutoff is better, but for the trace elements listed above, the 20-MeV cross section is used in the 20- to 150-MeV energy range. In all cases, physics models are used above 150 MeV.

The differences between the mix-and-match capability and having a 150-MeV boundary between data tables and physics models is tiny because only the few trace elements previously mentioned use the 20-MeV data in the range of 20 to 100 MeV. Using a 20-MeV boundary between data tables and physics models makes a small but noticeable difference for the neutron heating.

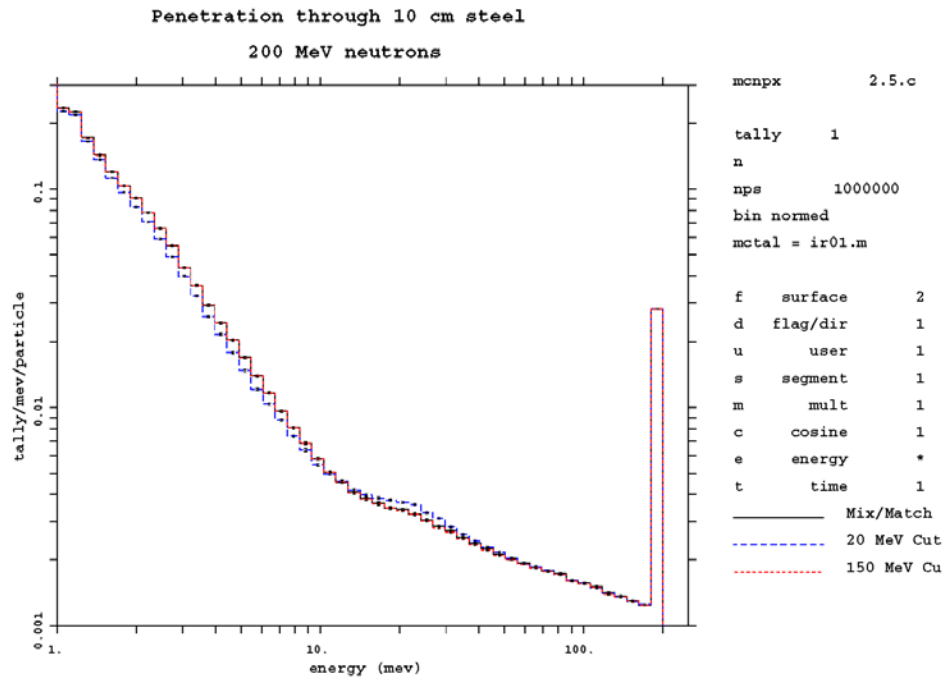


Fig. 22. The 200-MeV neutron steel penetration, new vs old.

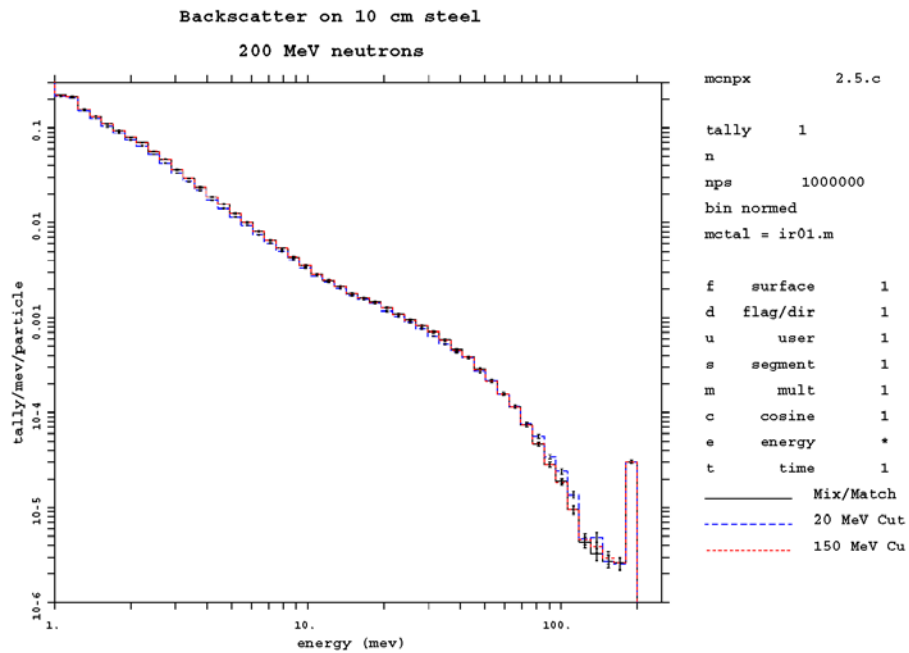


Fig. 23. The 200-MeV neutron steel backscatter, new vs old.

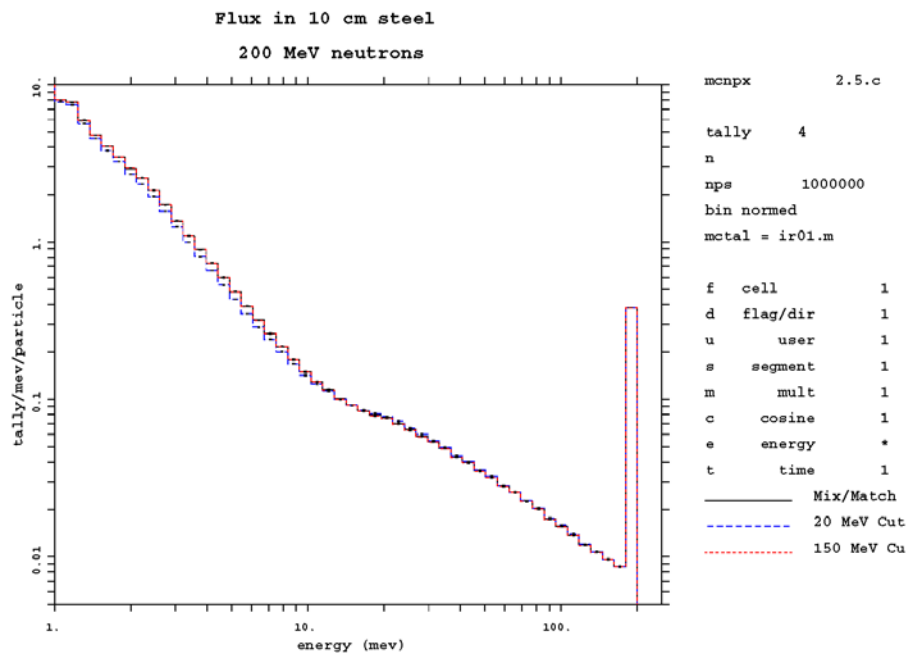


Fig. 24. The 200-MeV neutron steel flux, new vs old.

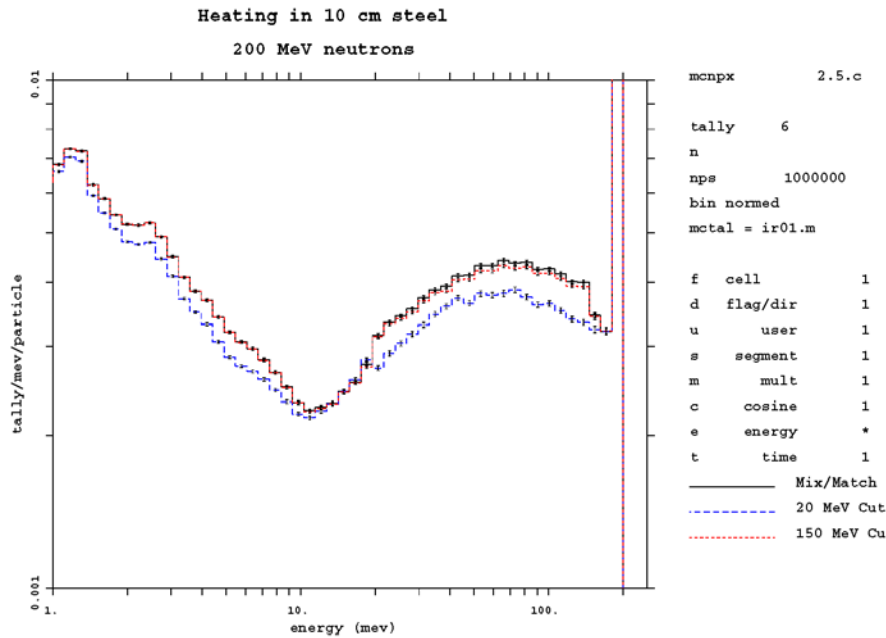


Fig. 25. The 200-MeV neutron steel heating, new vs old.

Figures 26 through 29 show the differences between using the mix-and-match tables for stainless steel and using strictly models instead.

The integral results over all energies are tabulated in Table 5. The neutron data tables and default physics model agree excellently. The CEM2k results are not as good.

When the CEM2k and default model calculations are rerun using a “zoo” of additional particles, the integral results change, as illustrated in Table 6. The “zoo” of particles added is

mode n h / z p d t a

Thus, in addition to neutrons, there are protons, positive pions, neutral pions, photons, deuterons, tritons, and alphas.

Note that the neutron heating in the neutron-only data table problem correctly includes the heating from all daughter particles, which is presumed to be deposited without further transport. The two model physics calculations, CEM2k and the default model, now agree better with the data table total heating calculation because the heating from all particles is included in all cases.

The total heating (for all particles) is illustrated in Fig. 30. As shown in Table 6, the integral results agree well; however, as shown in Fig. 30, the neutron-only (data table) calculation deposits most of the energy in the 180- to 200-MeV energy bin. The model physics calculations with all of the progeny particles deposit the energy along the progeny particle tracks.

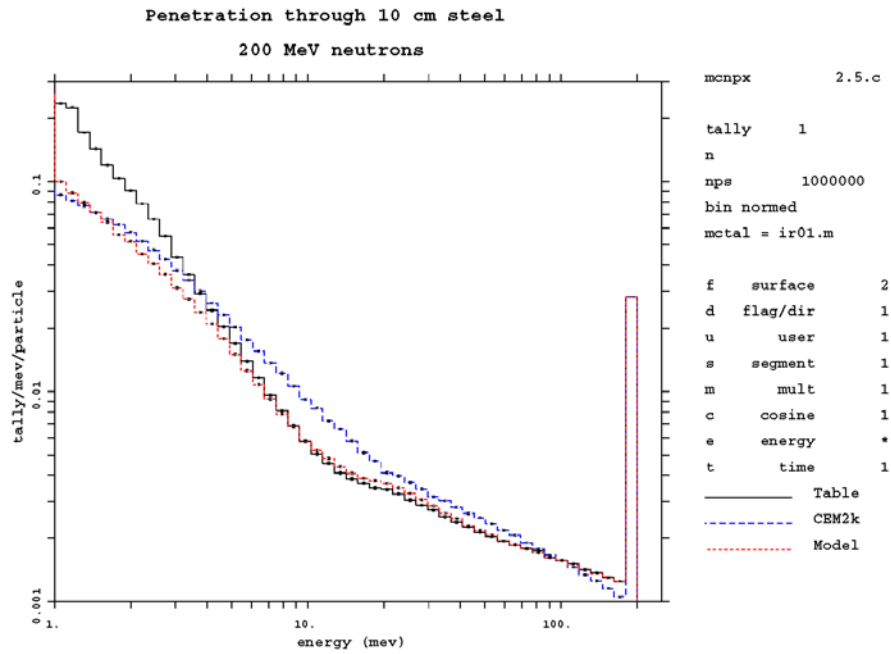


Fig. 26. The 200-MeV neutron steel penetration, table vs model.

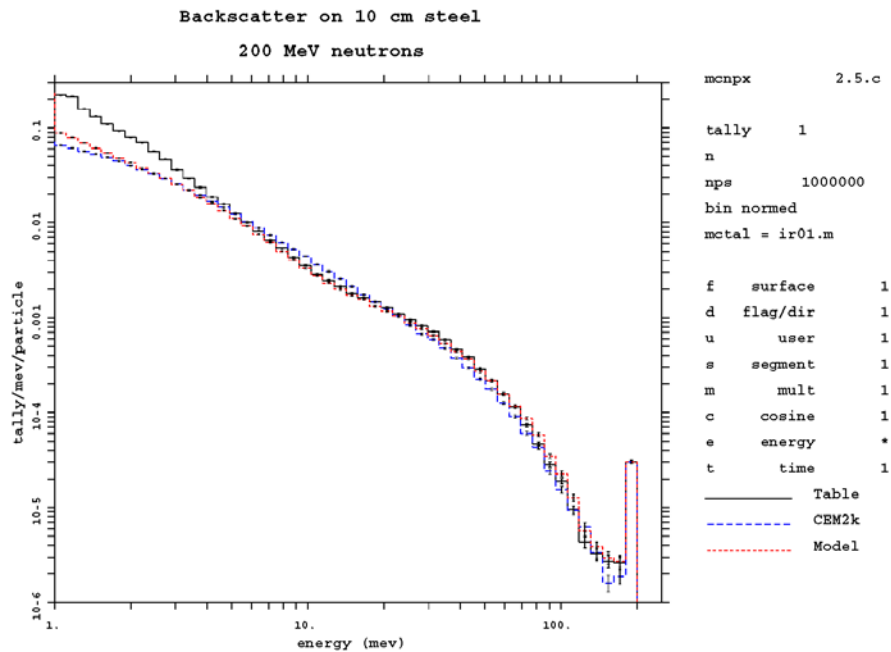


Fig. 27. The 200-MeV neutron steel backscatter, table vs model.

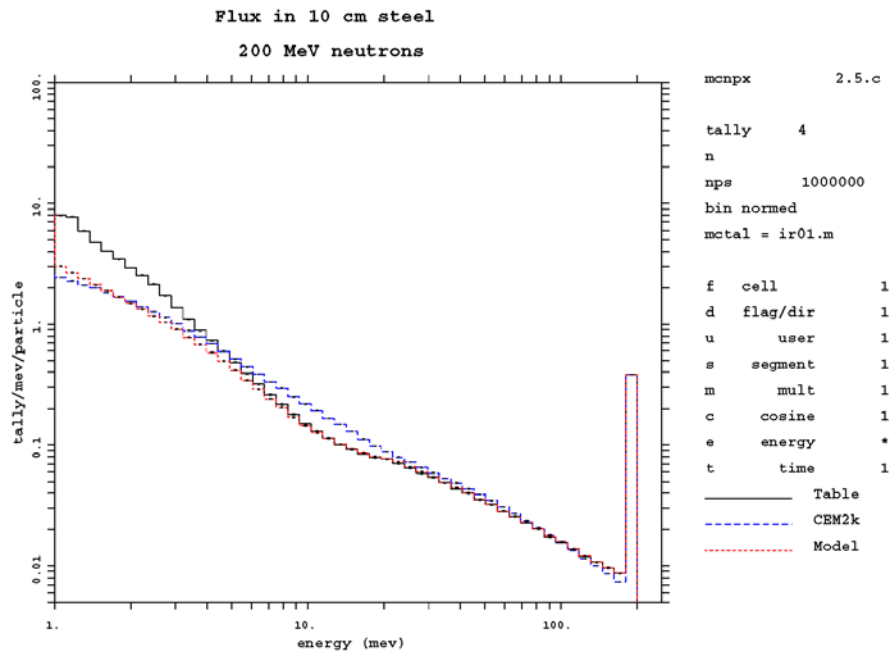


Fig. 28. The 200-MeV neutron steel flux, table vs model.

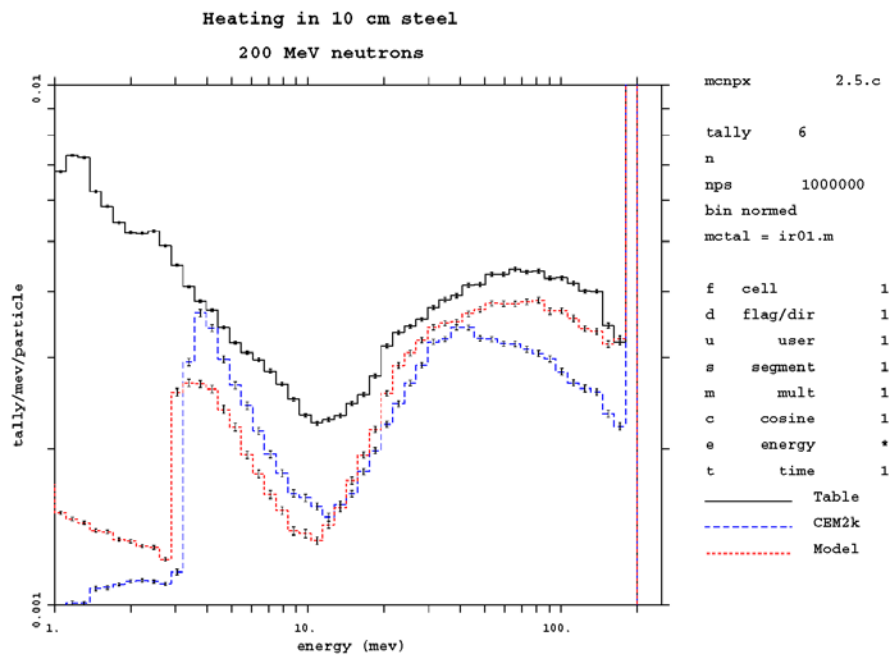


Fig. 29. The 200-MeV neutron steel heating, table vs model.

Table 5
200-MeV Neutrons Impinging on 10-cm-Thick Stainless-Steel Slab: 0 to 200 MeV

Quantity	Data Table	CEM2k	Model
Neutron Penetration	1.5595E+00	1.2567E+00	1.5664E+00
Neutron Backscatter	6.5962E-01	2.7833E-01	6.5341E -01
Neutron Flux	3.5544E+01	2.1934E+01	3.3541E+01
Neutron Heating	4.0758E+00	3.3203E+00	3.9753E+00

Table 6
200-MeV Neutrons Impinging on 10-cm-Thick Stainless-Steel Slab: 0 to 200 MeV

Quantity	Data Table	CEM2k+zoo	Model+zoo
Neutron Penetration	1.5595E+00	1.2650E+00	1.5875E+00
Neutron Backscatter	6.5962E-01	2.8266E-01	6.6924E-01
Neutron Flux	3.5544E+01	2.2095E+01	3.4103E+01
Neutron Heating	4.0758E+00	1.9220E-01	1.4086E-01
Total Heating	4.0758E+00	4.0309E+00	4.2283E+00

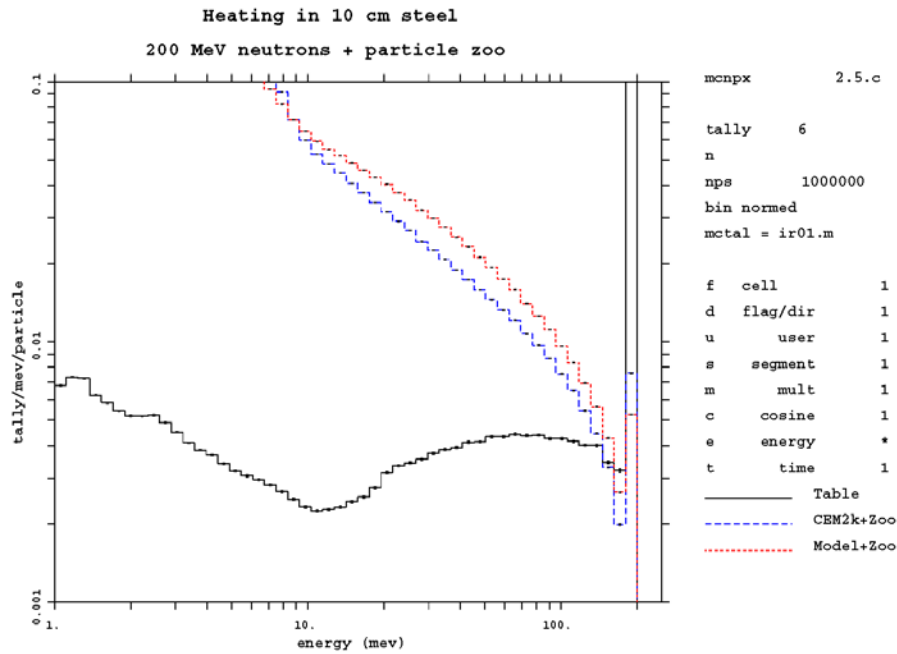


Fig. 30. The 200-MeV neutron steel heating, table vs model, zoo.

A closer examination of the low-energy results indicates that the physics models are really quite good until the neutrons fall below 1 MeV. Figures 31 through 34 refer to a 20-MeV monodirectional source normal to a 1-cm-thick stainless-steel infinite slab.

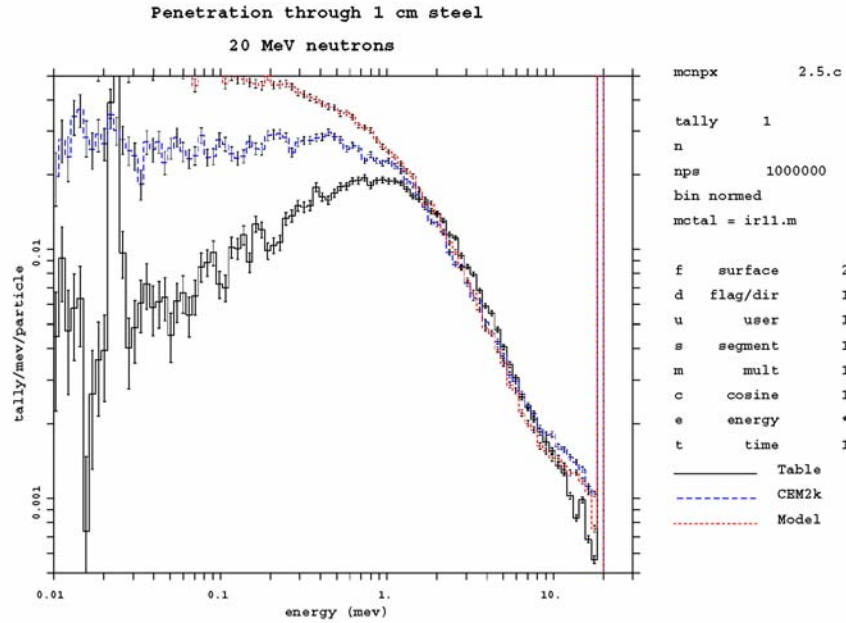


Fig. 31. The 20-MeV neutron steel penetration, table vs model.

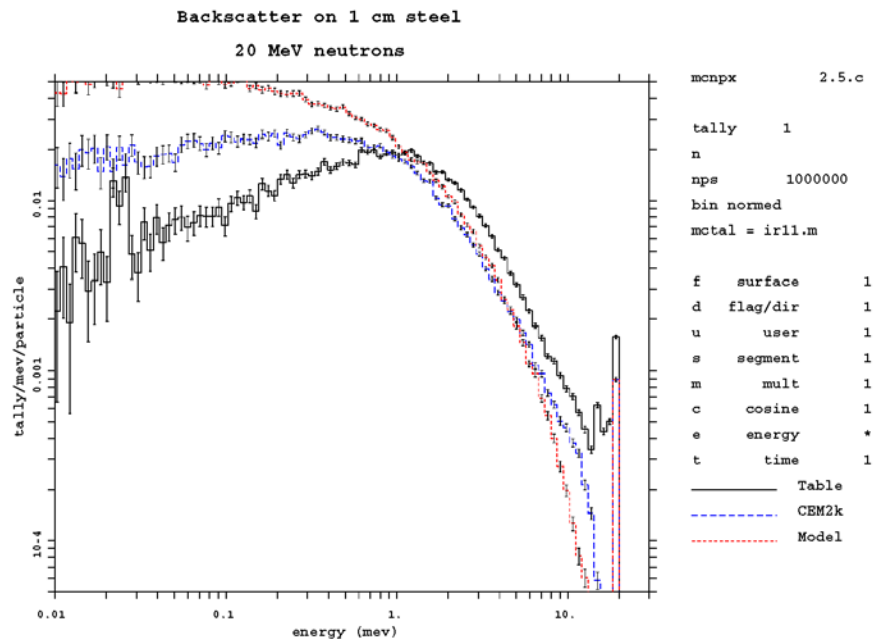


Fig. 32. The 20-MeV neutron steel backscatter, table vs model.

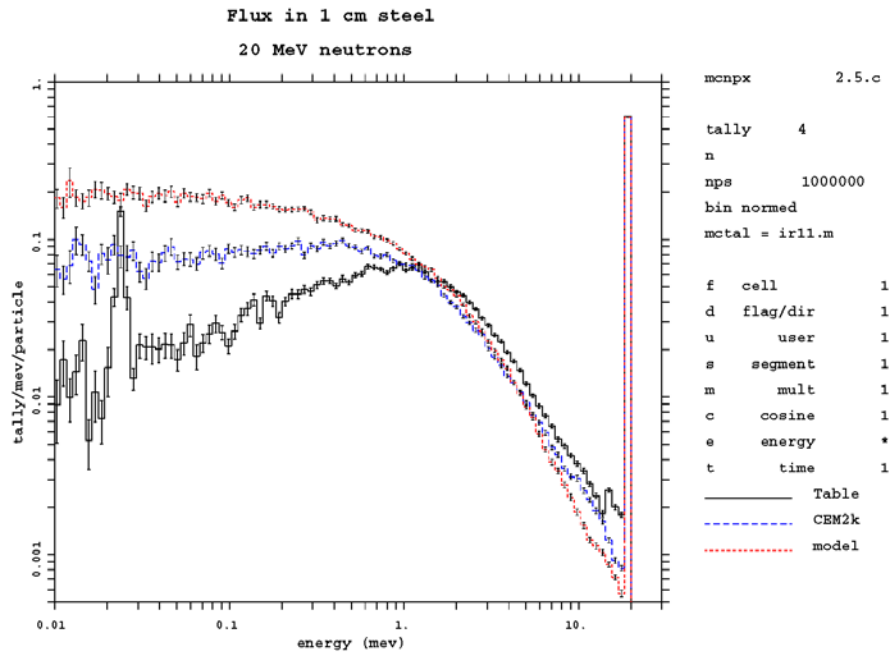


Fig. 33. The 20-MeV neutron steel flux, table vs model.

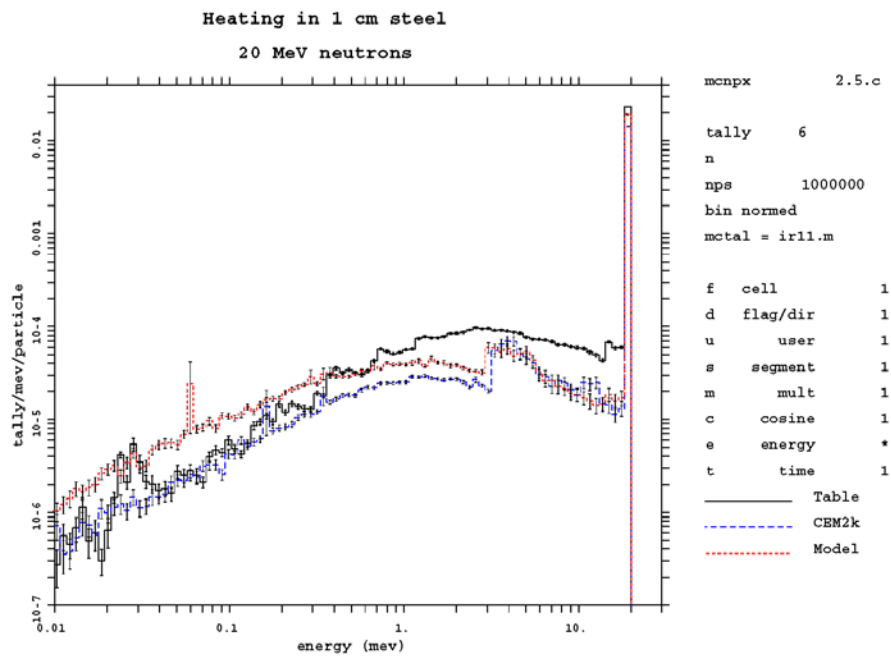


Fig. 34. The 20-MeV neutron steel heating, table vs model.

The integral results are presented in Table 7.

The integral results from 0 to 20 MeV show good agreement between the data table and the physics model calculations. Although the neutron penetration looks similar, the attenuations (fraction of source particles not penetrating the thin slab) for the data table and the default-model physics calculations are significantly different. The fluxes are in excellent agreement. Note that when only neutrons are run, the heating from all particles is put into the neutron heating. When the zoo of particles is run, the neutron heating is small and the total heating now compares to the total neutron heating (20% difference). Note further that the neutron heating number does not include photons or electrons; these are included only by using all the particles in the zoo (electron heating is included in the photon heating of the zoo); thus, the zoo heating is higher than the neutron heating in the mode N problems.

Figure 35 compares the total heating from the data table calculation and the model + zoo calculation. Note the constant heating for the model + zoo below 1 MeV because the heating data run out at 1 MeV; however, in this problem the default energy cutoff for model particles was dropped from the default 1 MeV to the lowest allowed, 1 keV.

The stainless-steel slab results show that the mix-and-match capability enables the use of data tables up to their maximum table energies before switching over to models. It is very advantageous not to have a global cutoff of, for example, 20 MeV, below which all particles in the problem use data tables and above which all particles in the problem use physics models.

Second, the stainless-steel problems demonstrate that for higher Z materials, neutron models can be used when no tables are available. The physics models do not really fail until the energies fall below ~1 MeV.

Table 7
20-MeV Neutrons Impinging on 1-cm-Thick Stainless-Steel Slab: 0 to 20 MeV

Quantity	Data Table	CEM2k	Model	Model+zoo
Neutron Penetration	9.6976E-01	9.7752E-01	9.8769E-01	9.8783E-01
Neutron Backscatter	6.8292E-02	5.4757E-02	7.0302E-02	7.0162E-02
Neutron Flux	1.2093E+00	1.1852E+00	1.2327E+00	1.2323E+00
Neutron Heating	3.8544E-02	2.3226E-02	3.1178E-02	6.1903E-03
Total Heating	3.8544E-02	2.3226E-02	3.1178E-02	4.6914E-02

Fig. 35. The 20-MeV neutron steel heating, table vs model + zoo.

5. TUNGSTEN COMPARISON

Tungsten is an important proton target for neutron conversion. The tungsten comparison is for a 20-cm-high can having a radius of 5 cm and a density of 19.3 g/cm^3 . The 200-MeV monodirectional protons start at the base center, and the penetration, backscatter, flux, and heating of protons and neutrons are calculated. Comparisons are between

- tungsten using 150-MeV data tables and neutron and proton models for ^{180}W ;
- tungsten using 150-MeV data tables with ^{180}W lumped into ^{182}W so that no models are used (this was the best capability before mix and match);
- tungsten with CEM2k neutron and proton models; and
- tungsten with the default models.

The default models are used above 150 MeV except for the CEM2k calculation, where the CEM2k models are used from 0 to 200 MeV. The lower-energy cutoff is set to zero so that the default model cutoff of 1 MeV is overridden; the lowest model energy then is

1 keV. The results of these calculations are compared in Figs. 36 to 43. Note that proton backscatter is omitted because it was small and statistically noisy.

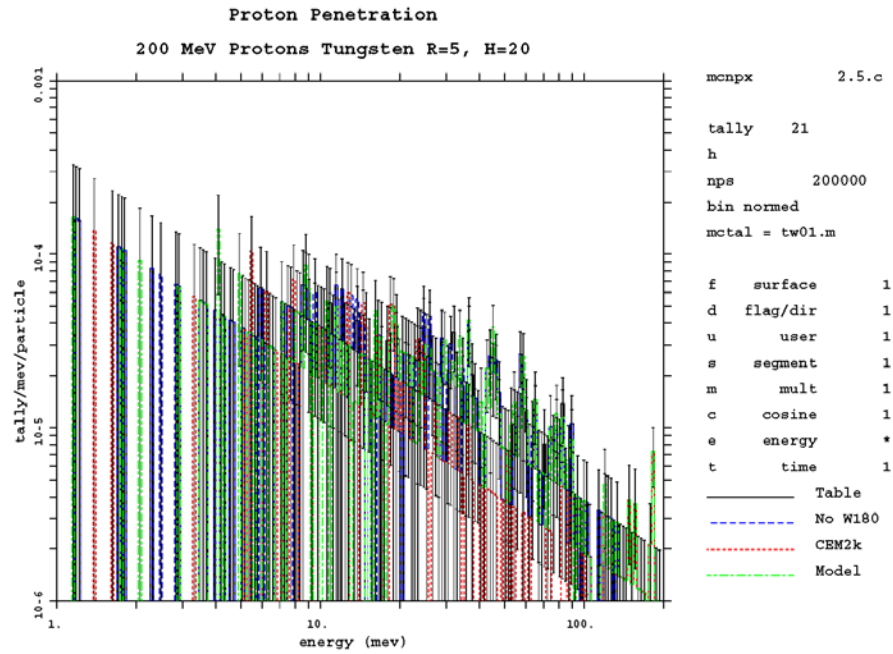


Fig. 36. The 200-MeV proton tungsten penetration, table vs model.

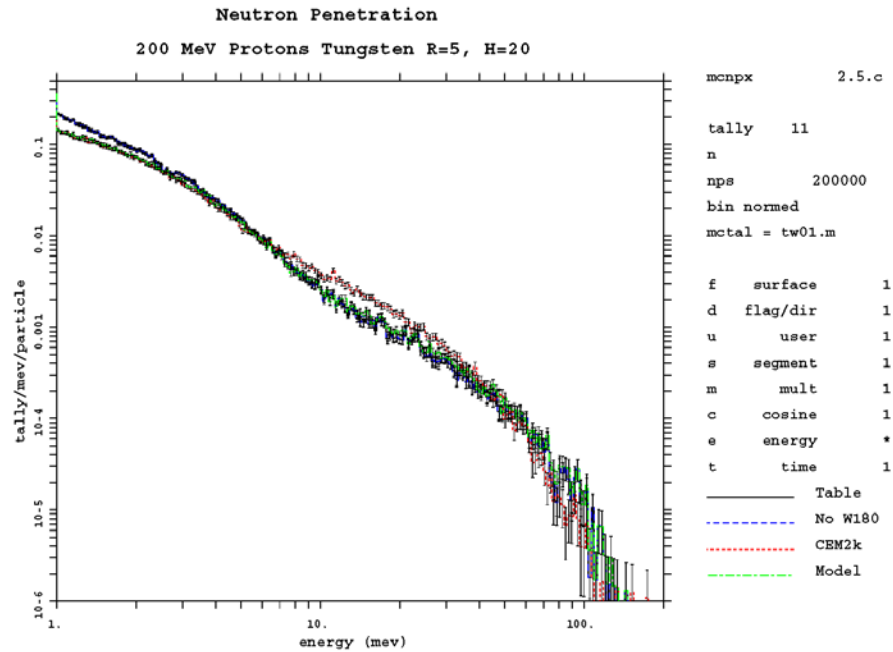


Fig. 37. Tungsten neutron penetration from 200-MeV protons, table vs model.

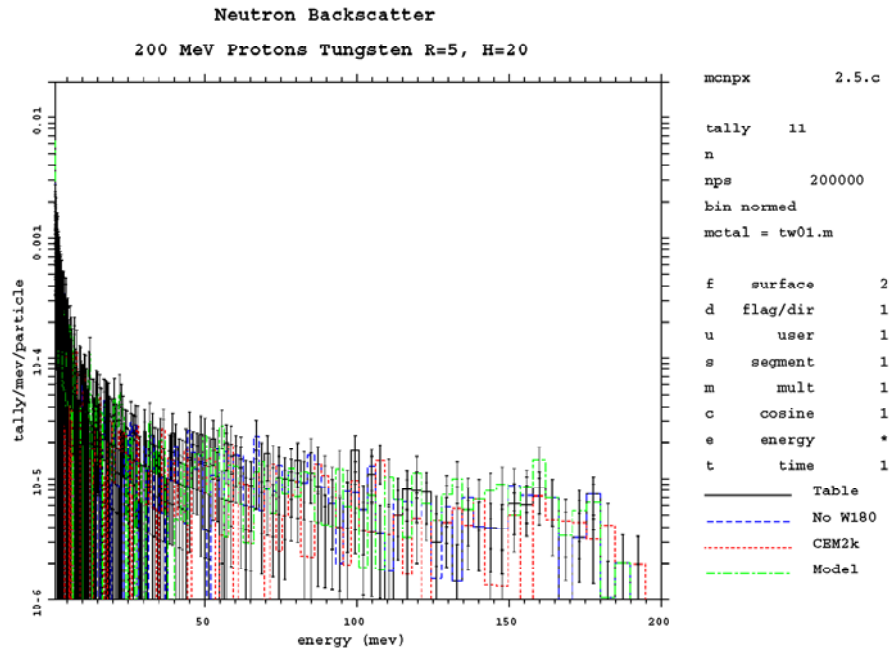


Fig. 38. Tungsten neutron backscatter from 200-MeV protons, table vs model.

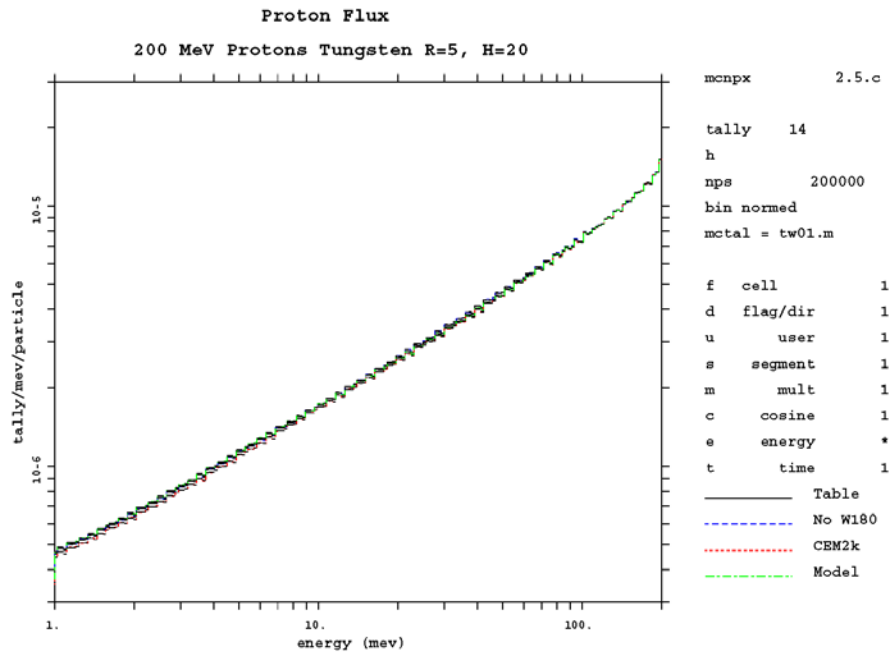


Fig. 39. The 200-MeV proton tungsten flux, table vs model.

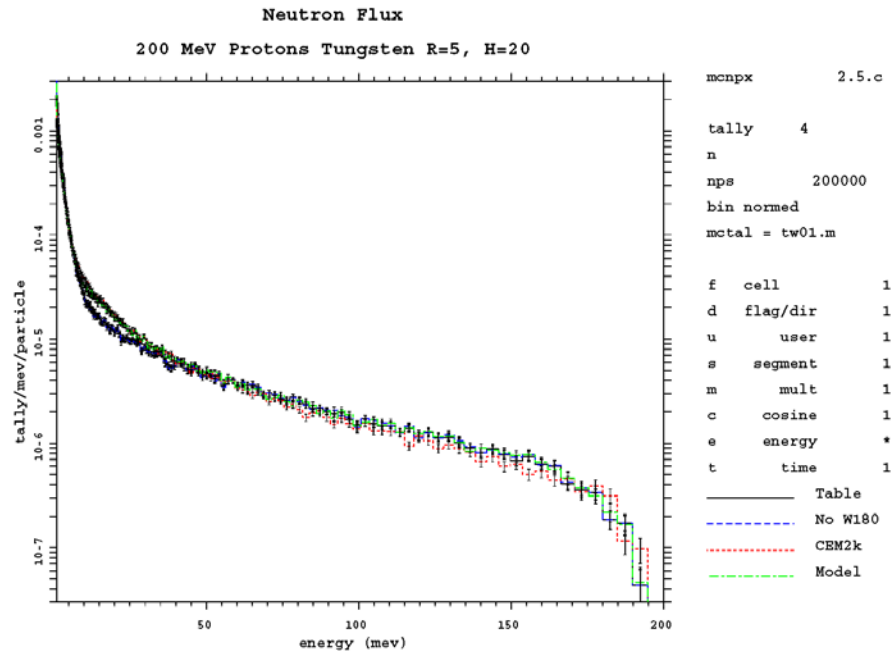


Fig. 40. Tungsten neutron flux from 200-MeV protons, table vs model.

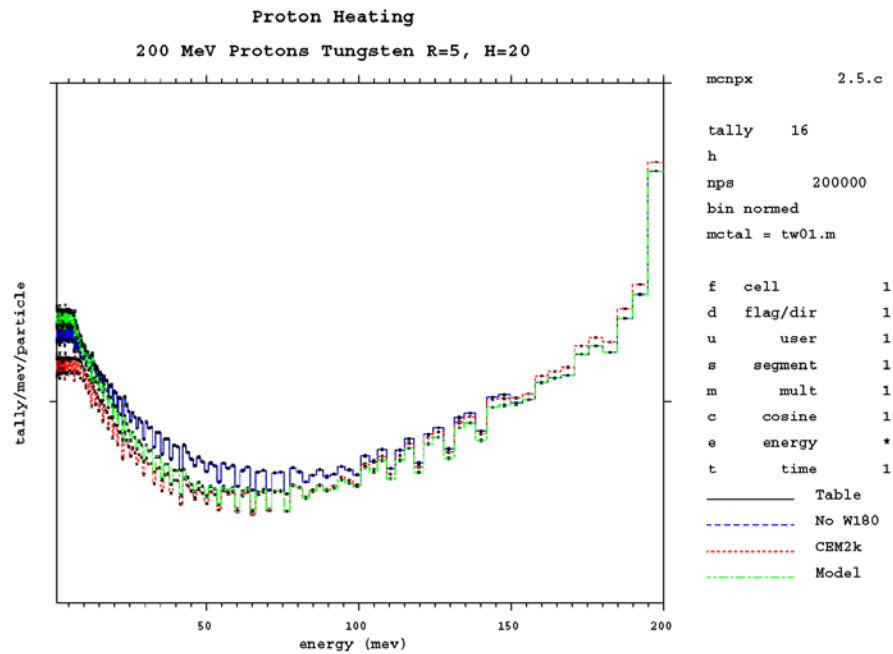


Fig. 41. The 200-MeV proton tungsten heating, table vs model.

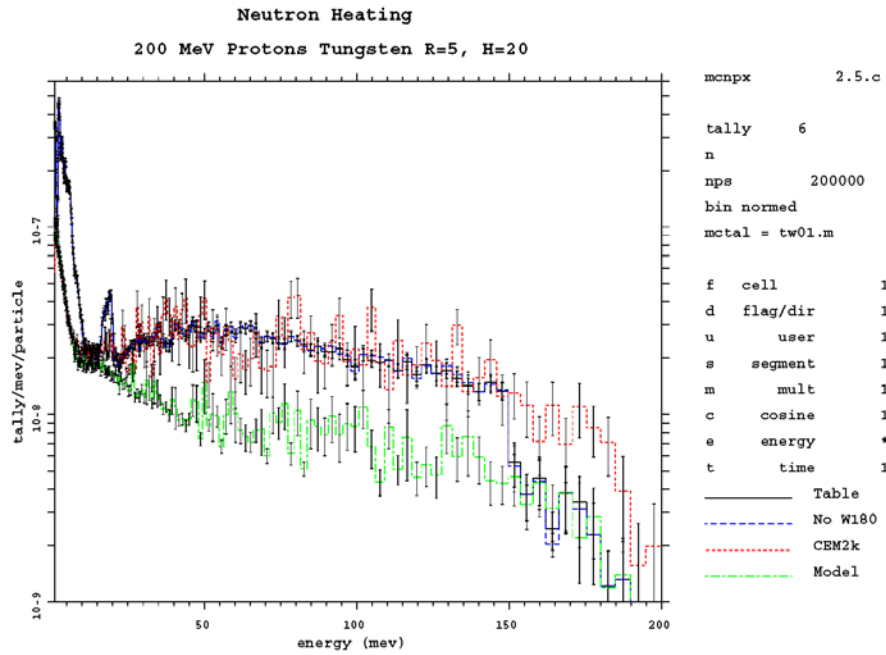


Fig. 42. Tungsten neutron heating from 200-MeV protons, table vs model.

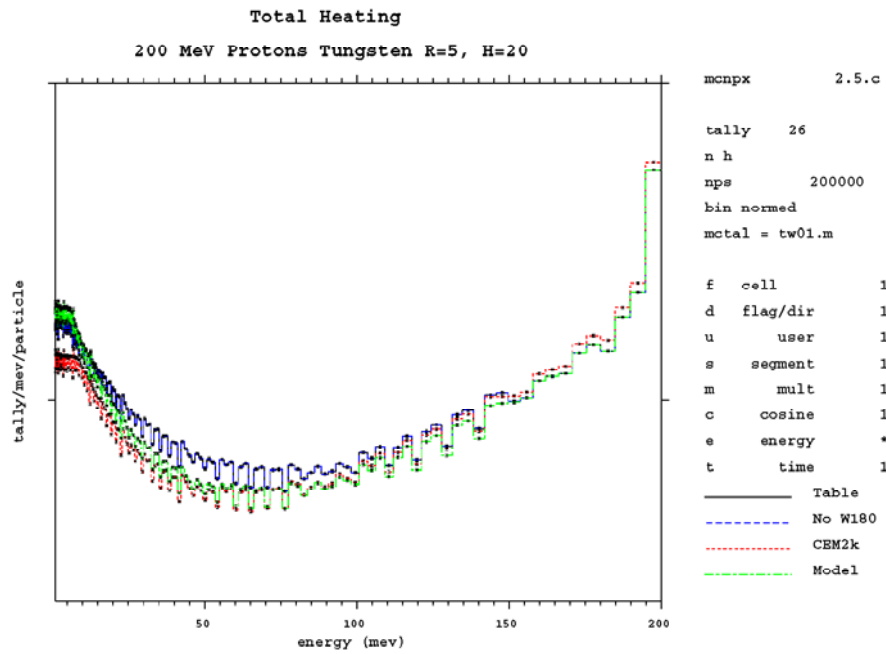


Fig. 43. Tungsten total heating from 200-MeV protons, table vs model.

The proton penetration is not plotted because it was not converged.

Table 8 presents the integral results from 0 to 200 MeV and from 1 to 200 MeV. Proton penetration and CEM2k proton backscatter were not converged. Neutron heating had a relative error of 3%, and proton backscatter had a relative error of 6%. All other quantities were converged to relative errors <1%.

The tungsten calculations demonstrate that the proton physics models have excellent agreement with the proton data tables. Further, a proton physics model for the trace isotope, ^{180}W , may be mixed with the proton data tables for the other isotopes of tungsten. Also, the approximation of lumping the trace isotope ^{180}W into ^{182}W is excellent.

Table 8
200-MeV Protons Impinging on Tungsten Can

0 to 200 MeV

Quantity	Data Table	NoW180	CEM2k	Model
Proton Penetration	1.5000E-05*	1.5000E-05*	1.0000E-05*	1.0000E-05*
Proton Backscatter	1.5400E-03	1.5500E-03	7.9400E-04*	1.5300E-03
Neutron Penetration	6.7376E-03	6.5435E-03	4.2100E-03	9.8350E-03
Neutron Backscatter	6.4091E-01	6.4052E-01	4.6088E-01	6.5053E-01
Proton Flux	1.5120E-03	1.5120E-03	1.5022E-03	1.5054E-03
Neutron Flux	7.3386E-03	7.3245E-03	4.4633E-03	7.5520E-03
Proton Heating	5.9487E-03	5.9484E-03	5.8872E-03	5.8968E-03
Neutron Heating	4.9915E-06	5.0039E-06	4.0502E-06	1.9285E-06
Total Heating	5.9536E-03	5.9534E-03	5.8912E-03	5.8988E-03

1 to 200 MeV

Quantity	Data Table	NoW180	CEM2k	Model
Proton Penetration	0.0000E-00*	5.0000E-06*	5.0000E-06*	1.5000E-05*
Proton Backscatter	1.5100E-03	1.5050E-03	8.1000E-04*	1.5050E-03
Neutron Penetration	3.3498E-03	3.2566E-03	2.6600E-03	3.9250E-03
Neutron Backscatter	3.0590E-01	3.0604E-01	2.7196E-01	2.6414E-01
Proton Flux	1.5116E-03	1.5117E-03	1.5019E-03	1.5050E-03
Neutron Flux	3.0868E-03	3.0876E-03	2.6837E-03	2.7537E-03
Proton Heating	5.9487E-03	5.9167E-03	5.8871E-03	5.8963E-03
Neutron Heating	4.7211E-06	4.7211E-06	4.1821E-06	1.9255E-06
Total Heating	5.9534E-03	5.9214E-03	5.8913E-03	5.8982E-03

*Unconverged.

Below 1 MeV, the neutron physics model results are unreliable and disagree with the neutron data table results. Comparing the 0- to 200-MeV results to the 1- to 200-MeV results shows that half of the neutron penetration, neutron backscatter, and neutron flux is <1 MeV. Consequently, the neutron physics models disagree with the data table results by as much as a factor of 3 over the 0- to 1-MeV energy range. Above 1 MeV, the agreement between physics models and data tables is much better, to within ~10% for penetration, backscatter, and flux. One reason why the neutron physics model results do not agree better with the neutron data table results is that the small differences between the proton physics models and data tables can result in differences in the neutron production, thus degrading the comparison of neutron physics models and data tables. The default physics model neutron heating is significantly lower than the neutron data table heating because the data table heating numbers include all of the progeny particle heating and the model neutron heating does not.

6. CARBON COMPARISON

Carbon is an important material for proton beamstops and other applications.

Comparisons are made for neutrons and protons for a 200-MeV monodirectional source impinges on a 100-cm-thick slab of 2.267 g/cm³ carbon (graphite). Neutron and proton penetration, backscatter, flux, and heating are compared. The calculations are the following:

- Data tables for 150 MeV are used for mixing natural carbon for neutrons and ¹²C for protons.
- Data tables use the ENDF/B-V ¹²C for neutrons and LA150 ¹²C for protons. Before mix and match, natural carbon could not be mixed with isotopic ¹²C. The only available 150-MeV libraries, the LA150 libraries,¹⁸ have ENDF/B-VI natural carbon for neutrons and isotopic carbon for protons. To have isotopic carbon for neutrons, the old ENDF/B-V evaluation must be used, which goes up to only 20 MeV.
- CEM2k physics models are used for neutrons and photons.
- Default physics models are used for neutrons and photons.

The results are displayed in Figs. 44 to 51.

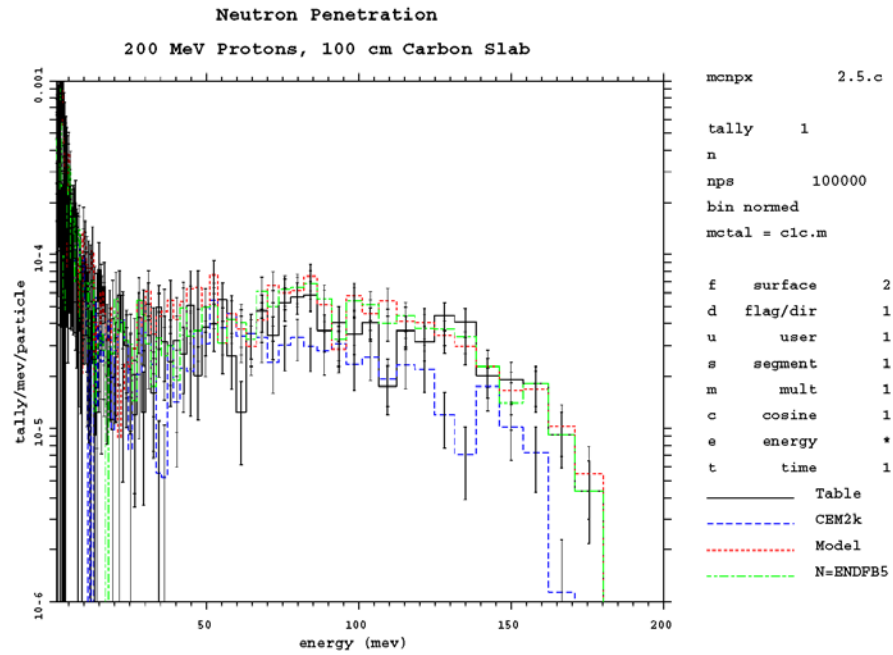


Fig. 44. Carbon neutron penetration from 200-MeV protons, table vs model.

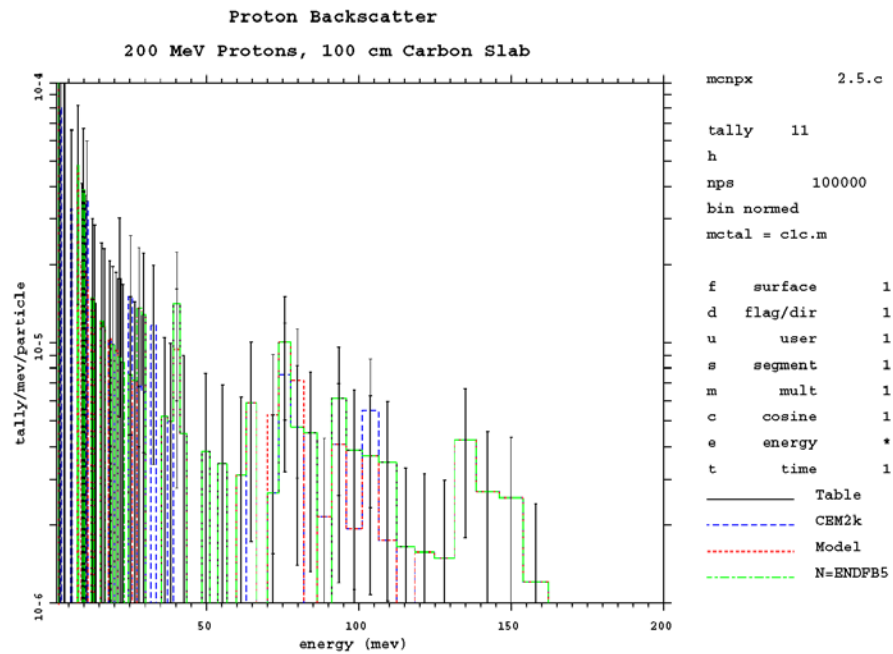


Fig. 45. Carbon proton backscatter from 200-MeV protons, table vs model.

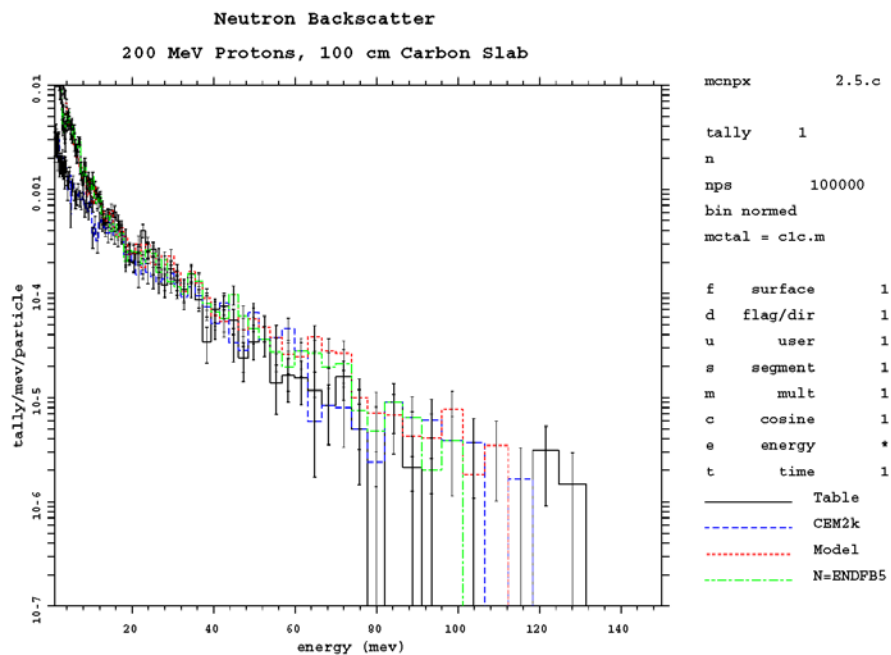


Fig. 46. Carbon neutron backscatter from 200-MeV protons, table vs model.

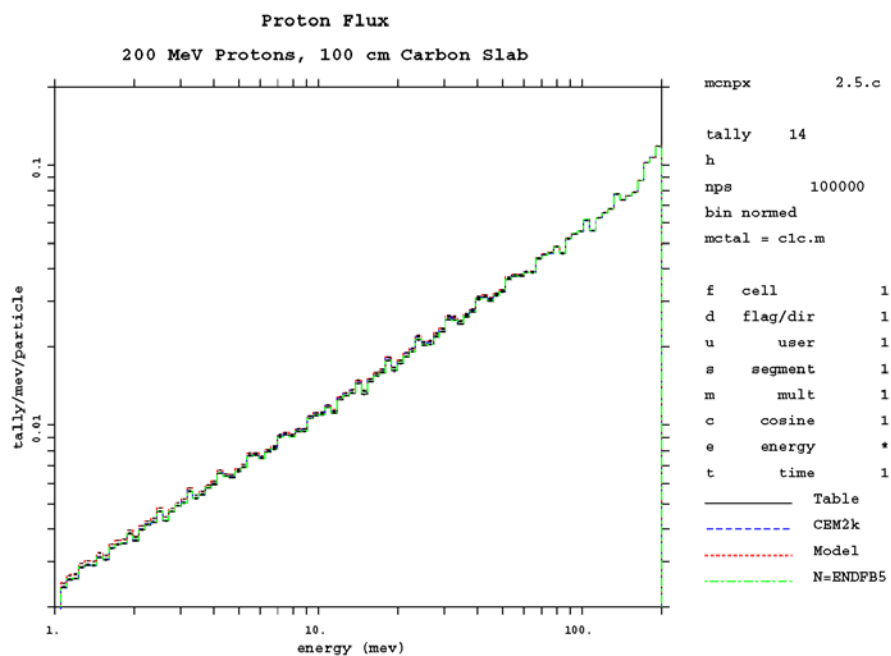


Fig. 47. Carbon proton flux from 200-MeV protons, table vs model.

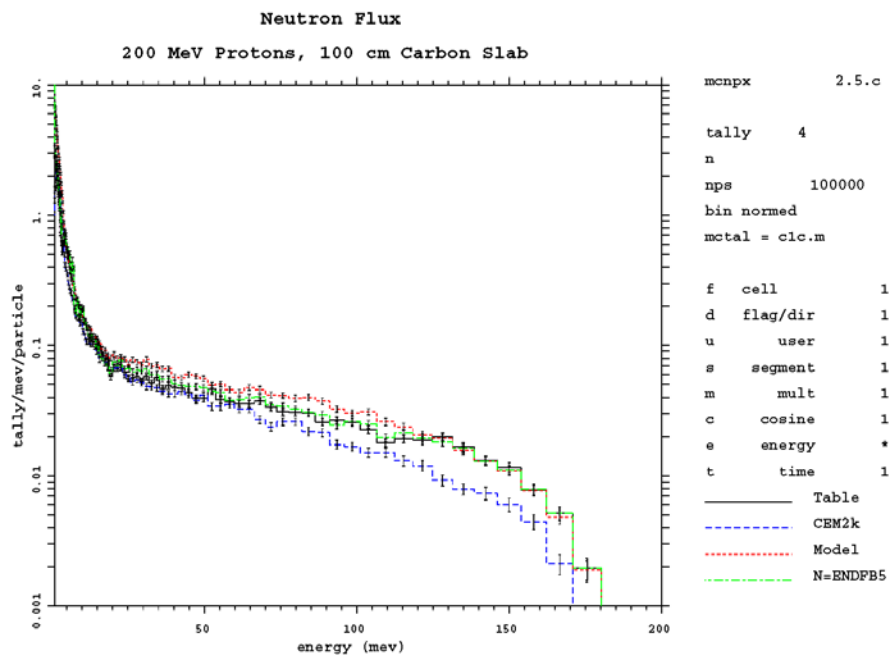


Fig. 48. Carbon neutron flux from 200-MeV protons, table vs model.

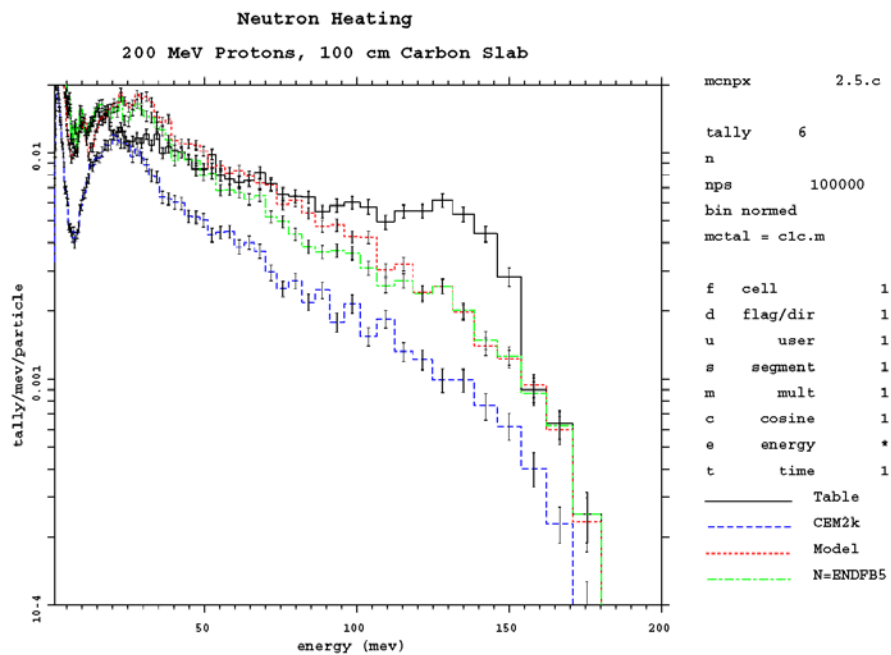


Fig. 49. Carbon neutron heating from 200-MeV protons, table vs model.

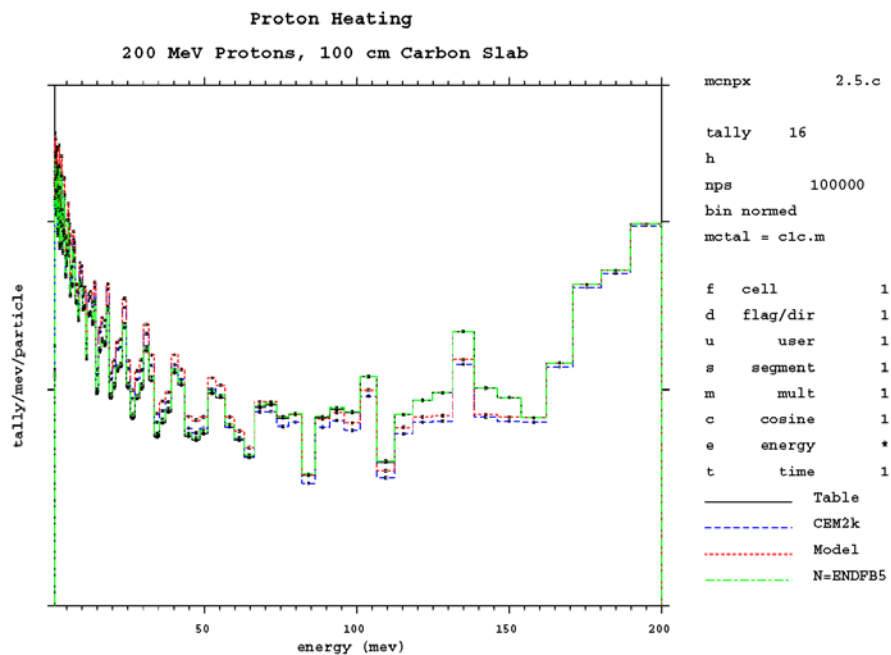


Fig. 50. Carbon proton heating from 200-MeV protons, table vs model.

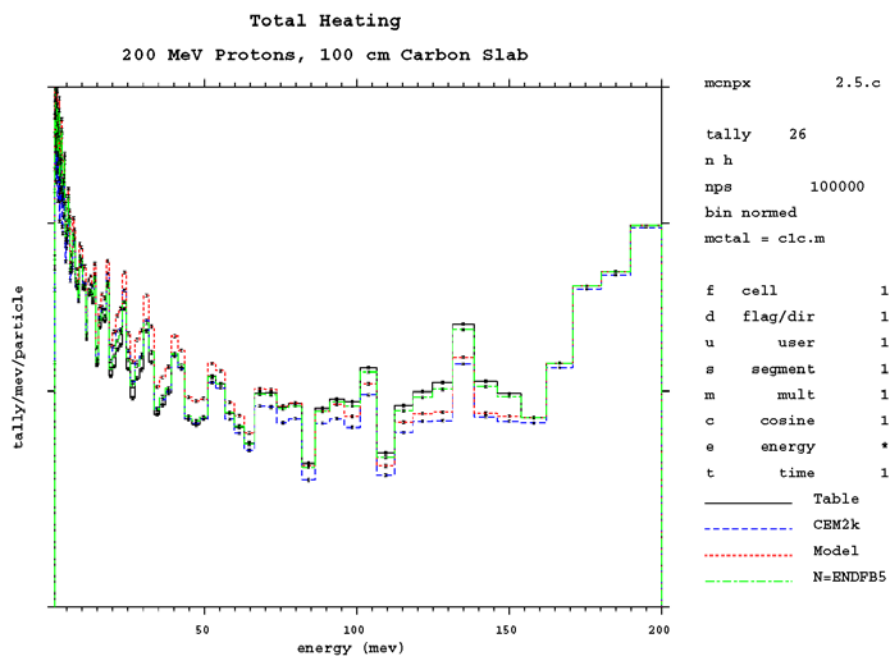


Fig. 51. Carbon total heating from 200-MeV protons, table vs model.

The integral tabular results from 0 to 200 MeV are presented in Table 9. All results are converged to <1% relative error, except for the proton penetration and backscatter, which are not converged because carbon absorbs protons very well.

The carbon slab comparisons illustrate the ability to mix different nuclides, specifically natural carbon for neutrons and ^{12}C for protons, so that 150-MeV libraries can be used for both neutrons and protons.

Proton models agree excellently with the proton tabulated data in the LA150 library, even for this low-Z material. Proton models may be used with confidence for nuclides where data tables are unavailable.

The total heating is essentially the proton heating. In the neutron heating plot, Fig. 49, the “ENDFB5” results agree with the “Model” results better than with the “Table” results. This is because in the 20- to 150-MeV range, the neutron default physics model is used for both “ENDFB5” and “Model” and only “Table” uses neutron data tables. In other plots, the “ENDFB5” results agree better with the “Table” results in the 20- to 150-MeV range because both use data tables for protons in this range.

The differences in the neutron penetration, backscatter, flux, and heating in Table 9 occur for several reasons. The proton transport is somewhat different for the data table and physics model calculations, and the subsequent differences in neutron production can make the neutron calculations different. Most of the neutron penetration, backscatter, and flux are from neutrons below 1 MeV, where the models are known to be unreliable. Table 10 shows the results for 20-MeV neutrons impinging on a 10-cm-thick carbon slab with a 1-MeV energy cutoff. The neutron default physics model agrees with the data table penetration and backscatter to within 30%, and the flux and heating agree within 15%.

Table 9
200-MeV Protons Impinging on 100-cm-Thick Carbon Slab: 0 to 200 MeV

Quantity	Data Table	N=ENDF5	CEM2k	Model
Proton Penetration	2.9994E-05*	6.0000E-05*	4.0000E-05*	7.0000E-05*
Proton Backscatter	5.9000E-04	5.9000E-04	3.9000E-04*	5.4000E-04
Neutron Penetration	3.0052E-02	3.6079E-02	9.0900E-03	3.5150E-02
Neutron Backscatter	1.4915E-01	1.6572E-01	2.1520E-02	1.5253E-01
Proton Flux	1.1516E+01	1.1525E+01	1.1532E+01	1.1564E+01
Neutron Flux	1.4309E+02	1.6859E+02	1.1397E+01	5.8180E+01
Proton Heating	8.1770E+01	8.1902E+01	8.1747E+01	8.2505E+01
Neutron Heating	1.3508E+00	1.2148E+00	6.2820E-01	1.2773E+00
Total Heating	8.3121E+01	8.3116E+01	8.2375E+01	8.3782E+01

*Unconverged.

Table 10
20-MeV Neutrons Impinging on 10-cm-Thick Carbon Slab: 1 to 20 MeV

Quantity	Data Table	CEM2k	Model
Neutron Penetration	6.2304E-01	7.6540E-01	7.4447E-01
Neutron Backscatter	1.2405E-01	4.4490E-02	1.6191E-01
Neutron Flux	1.2499E+01	1.1424E+01	1.3988E+01
Neutron Heating	2.0572E+00	1.4415E+00	1.7758E+00

Table 11 shows the results for 20-MeV neutrons impinging on a 10-cm-thick carbon slab with an 8-MeV energy cutoff. Other than backscatter, the default physics model results are within 2% to 12% of the data table results and the CEM2k results are 4% to 25%.

7. SUMMARY

The MCNPX mix-and-match capability makes it possible to substitute physics models for data tables at will. The ability to intermingle data tables and models is especially valuable for proton and photonuclear physics, where the models are best. Previously, either all physics models or all data tables for proton and photonuclear physics had to be used. Now data tables can be used whenever they are available. Further, if elements are required for one particle type, such as 6000.24c for neutrons, isotopic data may be used for another data type, such as 6012.24h, by nuclide substitution. Previously, 6000.24c and 6012.24h could not be used in the same calculation; frequently, the best choice was to totally ignore the proton table data and use proton physics models exclusively. For photonuclear calculations, only a limited set of nuclides was available previously and models could not be used at all.

The ability to intermingle data tables and models for neutrons enables the use of models for nuclides where no data tables are available. The neutron models are not as good as the neutron data tables, particularly at low energies and for low-Z materials. Below 1 MeV, they are unreliable. From 1 to 8 MeV, the CEM2k model is known to be unreliable and the default model is poor. Integral results from 1 to 20 MeV and from 8 to 20 MeV show the disagreement between physics models and data tables to be ~10% to 20%, and plots show differences of a factor of 2 to 10 at specific energies. Above 20 MeV, the physics models are good, although data tables should be used whenever available.

The mix-and-match capability also enables full utilization of neutron data tables with different upper-energy bounds. (Currently, all photonuclear and proton data libraries have the same 150-MeV upper-energy bound.) The matching of energies eliminates the need to ignore table data above 20 MeV in some problems and to extrapolate the 20-MeV data all the way up to 150 MeV in other problems.

The MCNPX mix-and-match capability is clearly a significant step forward for a wide variety of important applications.

Table 11
20-MeV Neutrons Impinging on 10-cm-Thick Carbon Slab: 8 to 20 MeV

Quantity	Data Table	CEM2k	Model
Neutron Penetration	5.2888E-01	5.6228E-01	5.6530E-01
Neutron Backscatter	3.5856E-02	1.6740E-02	2.1660E-02
Neutron Flux	9.3482E+00	9.0042E+00	9.1407E+00
Neutron Heating	1.9541E+00	1.4201E+00	1.7056E+00

APPENDIX

MCNPX Input Files

BGO Crystal Response

```
BGO Crystal Response
1 1 -7.130 -1 imp:n=1
2 0 1 imp:n=0

1 RCC 0 0 0 0 0 8.433 3.932

sdef sur=1.3 vec=0 0 1 dir=1 erg=100
nps 1000000
prdmp 2j -1
print -161 -162
m1 83209.24c 0.2105 8016.24c 0.6316 32000 0.1579
c m1 83209.24c 0.2105 8016.24c 0.6316 33075 0.1579
c phys:n 101 3j 20 $ Arsenic model above 20 MeV
c LCA 8j 1 $ Use CEM2k
mode n p
f4:n 1
f14:p 1
f6:n 1
f16:p 1
e16 0. 50i 9.143
f11:n 1.2
e0 1 100log 100
f21:p 1.2
```

Arsenic Crystal: Data Table vs Models

```
Comparison of Arsenic table vs model
1 1 -7.130 -1 imp:n=1
2 0 1 imp:n=0

1 RCC 0 0 0 0 0 8.433 3.932

sdef sur=1.3 vec=0 0 1 dir=1 erg=20
nps 1000000
prdmp 2j -1
print -161 -162
m1 33075 1
phys:n 4j 150
c phys:n 4j 0 $ Replace to use models
c LCA 8j 1 $ Add for CEM2k
mode n p
f4:n 1
f14:p 1
f6:n 1
f16:p 1
e16 0. 50i 9.143
f11:n 1.2
```

```
e0 1 100log 100
f21:p 1.2
```

Water Sphere

Test of Mix and Match Water

```
1 101 -1. -1 imp:n=1
2 0 1 imp:n=0
```

```
1 so 20
```

```
m101 1001.24c 2 8016.24c 1
mt101 lwtr $ Remove for models
phys:n 200 $ 5th entry = 0 for models
sdef erg=150 $ erg=2 for 2 MeV problems
LCA 8J 1 $ Remove for default model
print
prdump 2j -1
e0 1e-6 100LOG 1 148i 150
f1:n 1
f4:n 1
f6:N 1
NPS 1000000
```

Stainless-Steel Slab

Test of mix and and match: iron

```
1 0 -1
2 1 -8. -2 1
3 0 -2
```

```
1 pz 0
2 pz 10
```

```
imp:n 0 1 0
sdef dir=1 vec 0 0 1 erg=200 sur=1
f1:n 1 2
e0 1e-6 5log 1 49log 200
fq1 e f
vol 1 1 1
f4:n 2
f6:n 2
+f16 2
print -161 -162
c Plate, sheet, strip, flat bar 304L stainless steel
m1 6000.24c -.0003 $ .0003 carbon
7014.24c -.0009963 7015.60c -.0000037 $ .001 nitrogen
14028.24c -.00691575 14029.24c -.0003525 $ silicon
14030.24c -.00023175 $ .0075 silicon
15031.24c -.00045 $ .00045 phosphorous
16032.60c -.0003 $ .0003 sulphur
24050.24c -.00862 24052.24c -.16752 $ chromium
```

```

24053.24c -.0191 24054.24c -.00476 $ .20 chromium
25055.60c -.02 $ .02 manganese
26054.24c -.03785619 26056.24c -.5962025 $ iron
26057.24c -.01424486 26058.60c -.002146485 $ .65045 iron
28058.24c -.0814176 28060.24c -.0314736 $ nickel
28061.24c -.0014256 $ nickel
28062.24c -.0043896 28064.24c -.0012936 $ .12 nickel
prtmp 2j -1
nps 1000000
phys:n 200 3j -1

```

Tungsten Can

```

Tungsten
1 1 -19.3 -1 imp:n=1
2 0 1 imp:n=0

1 RCC 0 0 0 0 0 20 5

mode n h
sdef sur=1.3 vec=0 0 1 dir=1 erg=200 par h
nps 200000
prtmp 2j -1 2
print
m1 74180 .0014 74182 .2641 74183 .1440 74184 .3064 74186 .2841
  nlib=24c
e0 1 200log 200
f11:n 1.3 1.2
fq11 e f
f21:h 1.3 1.2
fq21 e f
f4:n 1
f14:h 1
f6:n 1
f16:h 1
+f26 1
cut:n j 0
cut:h j 0
c LCA 8j 1
phys:n 201 3j -1
phys:h 201 j -1

```

100-cm-Thick Carbon Slab

```

200-MeV Protons on 100-cm Carbon
1 0 -1
2 1 -2.267 -2 1
3 0 -2

1 pz 0
2 pz 100

```



```

imp:n,h 0 1 0
print
mode h n
vol 1 1 1
sdef dir=1 vec 0 0 1 erg=200 sur=1 par=h
m1 6012.24 1 nlib=24c hlib=24h
mx1:n 6000
f1:n 1 2
e0 1 100log 200
fq1 e f
f4:n 2
f6:n 2
f16:h 2
+f26 2
f11:h 1 2
fq11 e f
f14:h 2
prdmp 2j -1 2
nps 100000
c LCA 8j 1
phys:n 201 3j -1
phys:h 201 j -1

```

10-cm-Thick Carbon Slab

1-20-MeV Neutrons on 10-cm Carbon

```

1 0 -1
2 1 -2.267 -2 1
3 0 -2

```

```

1 pz 0
2 pz 10

```

```

imp:n 0 1 0
print
vol 1 1 1
sdef dir=1 vec 0 0 1 erg=20 sur=1
m1 6000.24c 1
e0 1 100log 20
f4:n 2
f6:n 2
f11:n 1
f21:n 2
prdmp 2j -1 2
nps 100000
c LCA 8j 1
phys:n 21 3j -1
cut:n j 1

```

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