

Jet quenching on hadron yields and ratios at RHIC and LHC

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1. **Energy loss of jets** in hot, dense non-Abelian plasma:
 - energy loss in a **THICK plasma** - BDMS, LCPI
 - energy loss in a **THIN plasma** - GLV method
 - ⇒ **OPACITY expansion** ($N = 1, 2, 3, \dots$)
2. **Jet quenching on hadron yields (RHIC, LHC)**
 - **E-indep.** vs. **E-dependent** energy loss
 - latest results from RHIC (π^0 in PHENIX)
3. **Baseline: pQCD calculation** in $p + p$, $p + A$, $A + A$:
 - **intrinsic k_T** in $p + p$ ($p + \bar{p}$)
 - p_T broadening in $p + A$ (**saturated Cronin**)
 - **A+A collision by pQCD** (high p_T)
 - ⇒ **G. Papp's QM'01 talk**
4. **Hadron ratios as probes of jet quenching (RHIC)**
 - K^+/π^+ , K^-/π^- ratios at $L/\lambda = 1, 2, 3, \dots$

QM'2001, Brookhaven — Stony Brook
19 January 2001

Bibliography:

- Jet quenching - GLV (opacity expansion, tagged target)
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- Jet quenching vs. Cronin effect at SPS/RHIC
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- pQCD description of $p + p$ and $p + A$ collision
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- Hadron ratios as probes of jet quenching at RHIC
P. Lévai, G. Papp, G. Fai, M. Gyulassy
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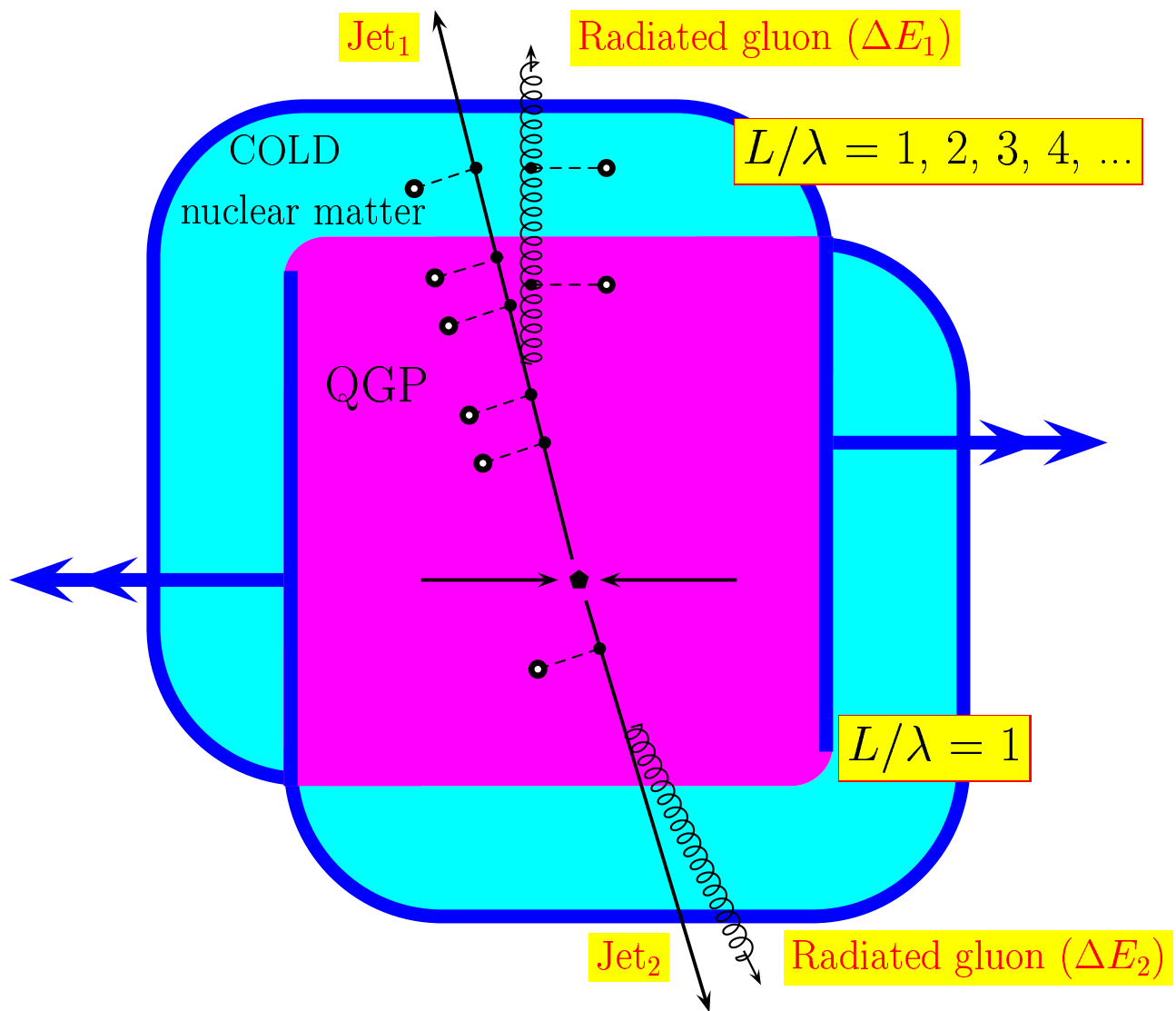
1. **Energy loss of jets** in non-Abelian plasma

A+A Collisions \implies QGP (non-pert. matter)

\implies High p_T jets (perturb. QCD)

Radiative energy loss (jet-quenching): $E_{jet} \longrightarrow E_{jet} - \Delta E$

pQCD probe of the non-pert. QGP phase



Medium induced radiative energy loss - 1

THICK plasma: $L \gg \lambda_g$ [LLA]

GW: QED - analogy $\implies \Delta E_{GW}/dx = 1-2 \text{ GeV/fm}$

M. Gyulassy, X.N. Wang '94

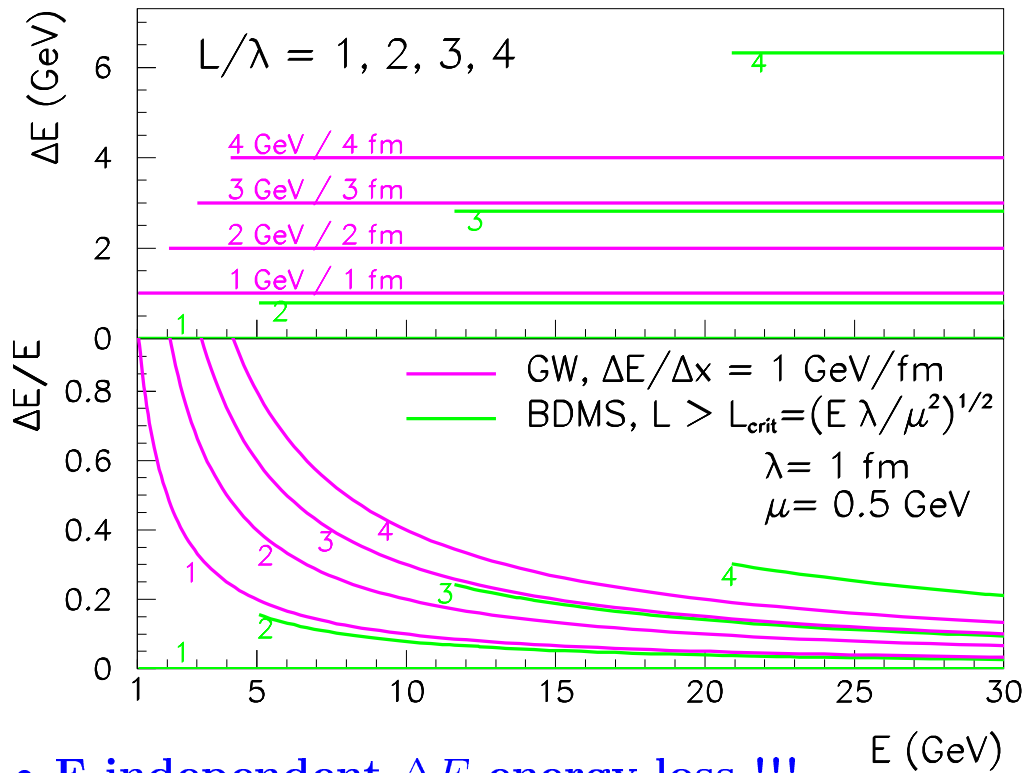
BDMS: time-ordered pQCD (Feynman diagram.) \implies

$$\Delta E_{BDMS} = \frac{C_R \alpha_s L^2 \mu^2}{4 \lambda_g} \log \frac{L}{\lambda_g}$$

Baier-Dokshitzer-Mueller-Schiff '97, '98

LCPI: light-cone path integral formalism $\implies \equiv$ BDMS

B.G. Zakharov '96, '98



• **E-independent ΔE energy loss !!!**

Medium induced radiative energy loss - 2

THIN plasma: $L \sim \lambda_g$

GLV: time-ordered pQCD (Feynman diagrammatic)

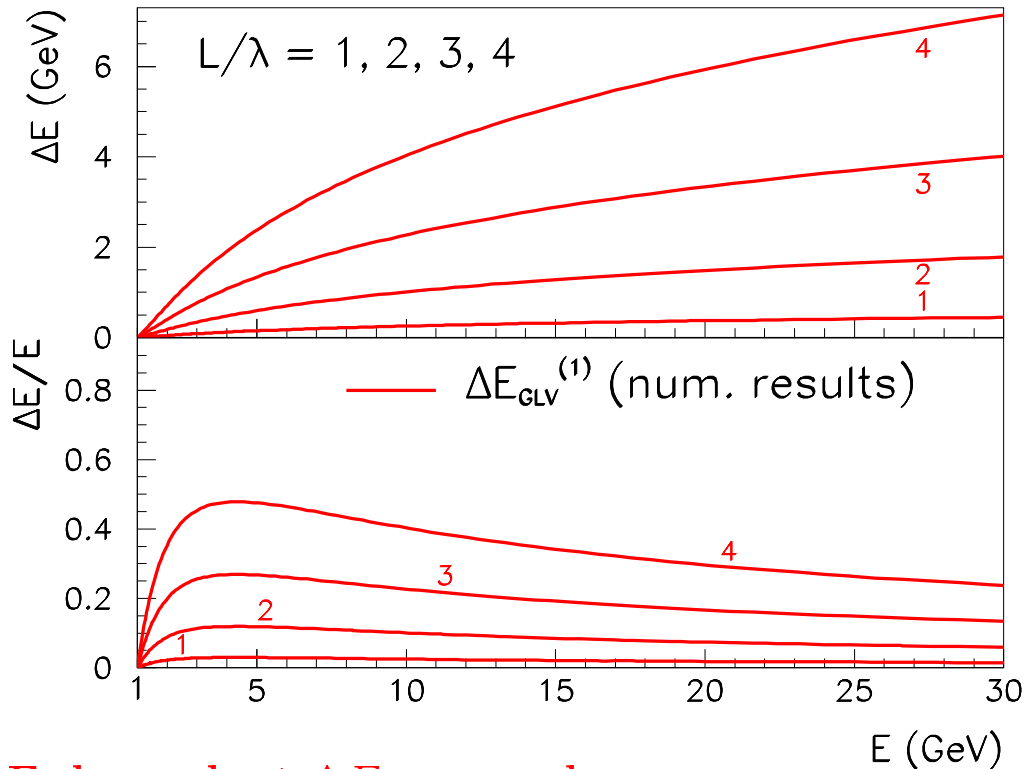
+ OPACITY expansion ($N = 1, 2, 3, \dots$)

+ kinematical cuts \Rightarrow U. Wiedemann, QM'01 talk

$$\Rightarrow \Delta E_{GLV} \approx \Delta E_{GLV}^{(1)} \approx \frac{C_R \alpha_s L^2 \mu^2}{N(E) \lambda_g} \log \frac{E}{\mu}$$

$N(E)$ is a numerical function, $N(E) \rightarrow 4$ at $E \rightarrow \infty$.

Gyulassy-Levai-Vitev, PRL85,5535; NPB594,371



- **E-dependent ΔE energy loss**
- **\approx E-independent $\Delta E/E$ in $3 < E < 10$ GeV**

Medium induced radiative energy loss - 3

THIN plasma: $L \sim \lambda_g$

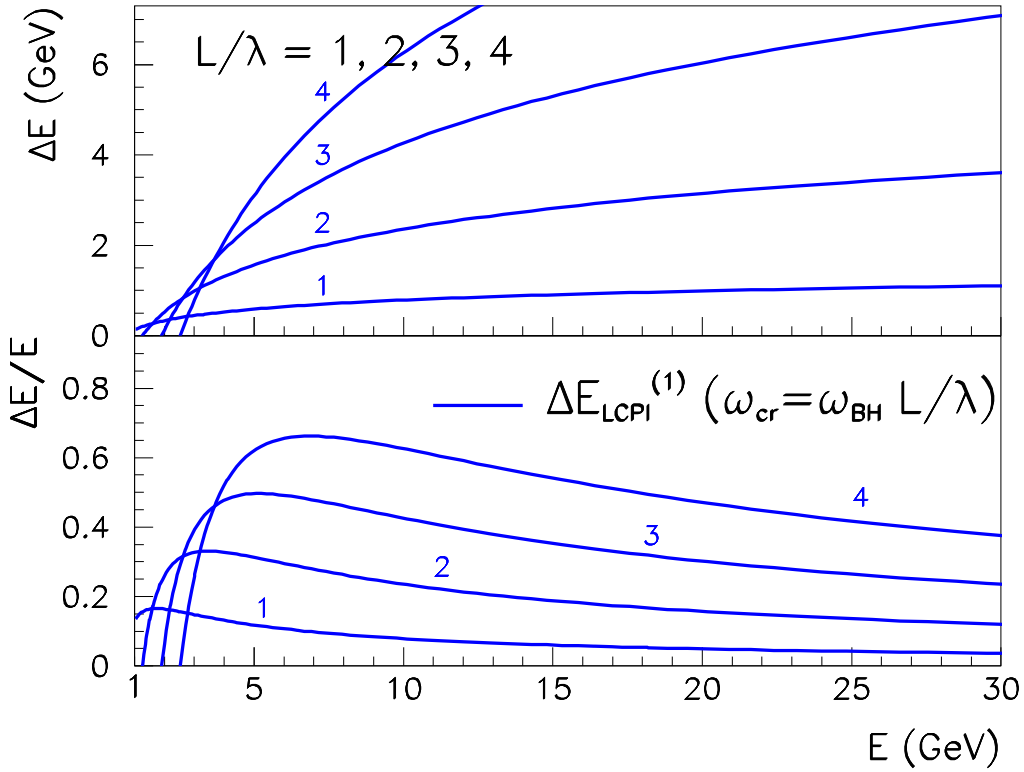
LCPI: light-cone path integral formalism at $N = 1$

(no kinematical cuts) \implies

$$\Delta E_{LCPI}^{(1)} = \frac{C_R \alpha_s L^2 \mu^2}{4 \lambda_g} \log \frac{E}{\omega_{cr}}$$

$$\omega_{cr} = \max (nC_3 L^2/4; L\mu^2/2 = \omega_{BH} \cdot L/\lambda)$$

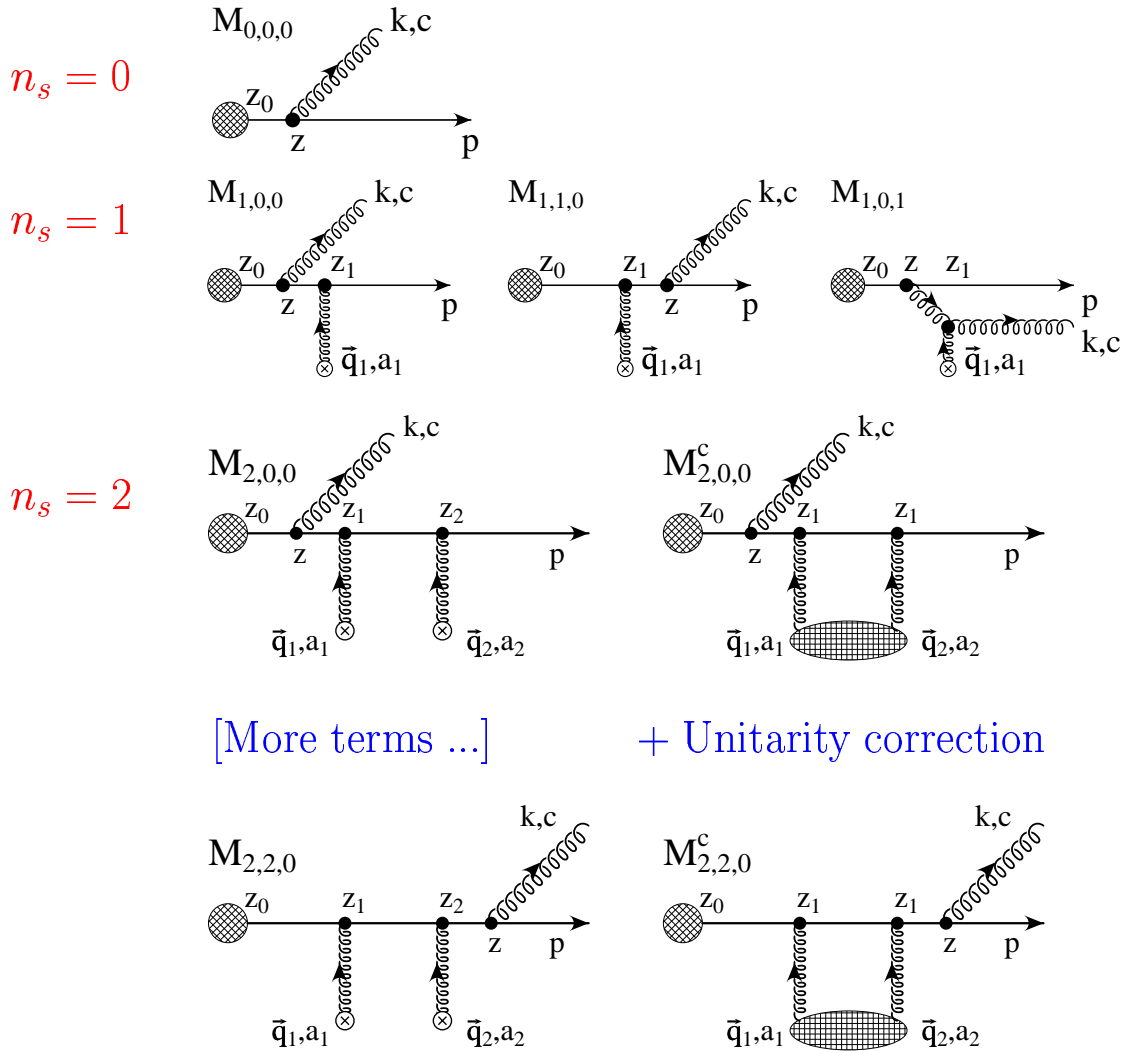
B.G. Zakharov, (hep-ph/0012360)



$$\iff \Delta E_{GLV}^{(1)} \approx \frac{C_R \alpha_s L^2 \mu^2}{N'(E) \lambda_g} \log \frac{E}{\omega_{BH}}$$

Opacity expansion in the GLV method - 1

Time-ordered perturbative calculation of gluon radiation
pQCD Feynman diagrams with n_s scattering centers:



- $n_s = 2$: 7 tree + 4 unitarity correction diagrams
- $n_s = 3$: 15 tree + 19 unitarity correction diagrams!!
- $n_s = 4$: 31 tree + ...

Opacity expansion in the GLV method - 2

Average number of jet scatterings, $\bar{n} = L/\lambda$

Gluon radiation amplitudes $I^{(N)} \iff (\sigma_{el}/A_{\perp})^N \propto (L/\lambda)^N$

OPACITY EXPANSION: $I = I^{(1)} + I^{(2)} + I^{(3)} + \dots$

$I^{(0)}$: $(L/\lambda)^0 \propto (\sigma_{el}/A_{\perp})^0$ contributions

- 1 $(n_s = 0) \times (n_s = 0)^{\dagger}$

$I^{(1)}$: $(L/\lambda)^1 \propto (\sigma_{el}/A_{\perp})^1$ contributions

- 4 $(n_s = 2) \times (n_s = 0)^{\dagger}$
- 9 $(n_s = 1) \times (n_s = 1)^{\dagger}$

$I^{(2)}$: $(L/\lambda)^2 \propto (\sigma_{el}/A_{\perp})^2$ contributions

- 13 $(n_s = 4) \times (n_s = 0)^{\dagger}$
- 57 $(n_s = 3) \times (n_s = 1)^{\dagger}$
- 65 $(n_s = 2) \times (n_s = 2)^{\dagger}$

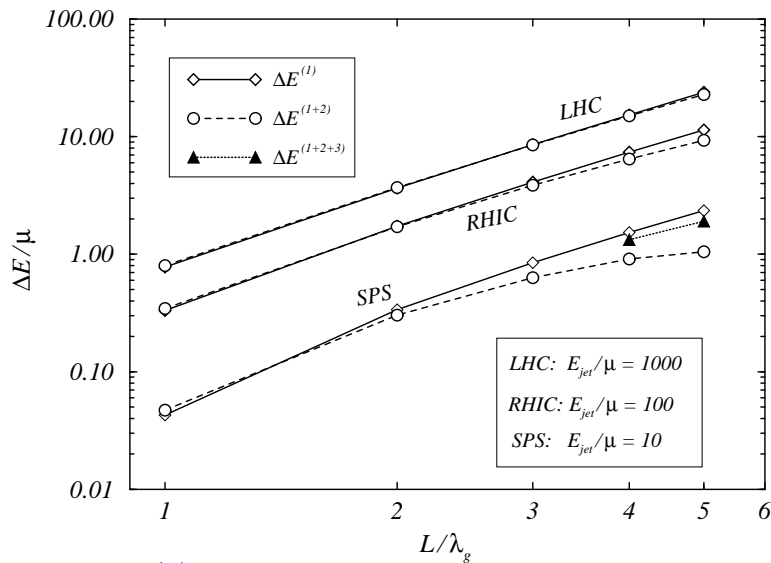
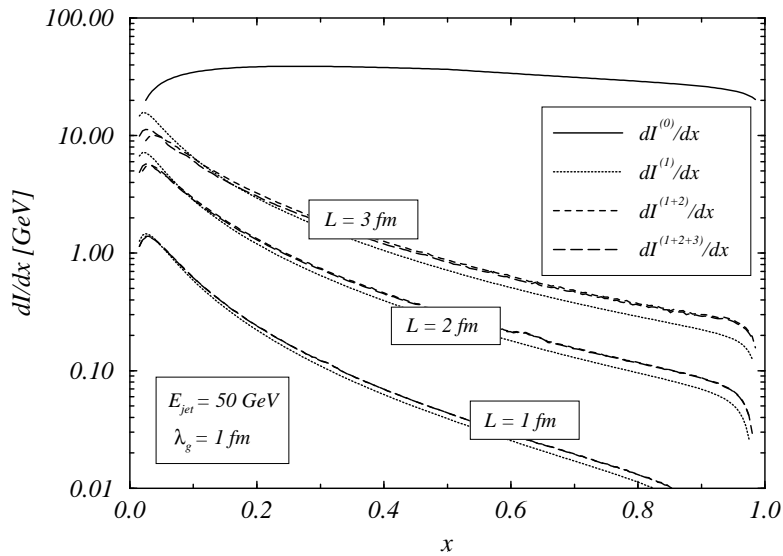
After evaluating these amplitudes [see NPB594(2001)371]

e.g. $xdI^{(1)}/dxdk_{\perp}^2 \implies \Delta E^{(1)}$:

$$\Delta E_{GLV}^{(1)} = \frac{2C_R\alpha_s EL}{\pi \lambda_g} \int_0^1 dx \int_0^{k_{max}^2} \frac{d\mathbf{k}_{\perp}^2}{\mathbf{k}_{\perp}^2} \int_0^{q_{max}^2} \frac{d^2\mathbf{q}_{\perp} \mu_{eff}^2}{\pi(\mathbf{q}_{\perp}^2 + \mu^2)^2} \cdot \frac{2\mathbf{k}_{\perp} \cdot \mathbf{q}_{\perp} (\mathbf{k} - \mathbf{q})_{\perp}^2 L^2}{16x^2 E^2 + (\mathbf{k} - \mathbf{q})_{\perp}^4 L^2}$$

Opacity expansion in the GLV method - 3

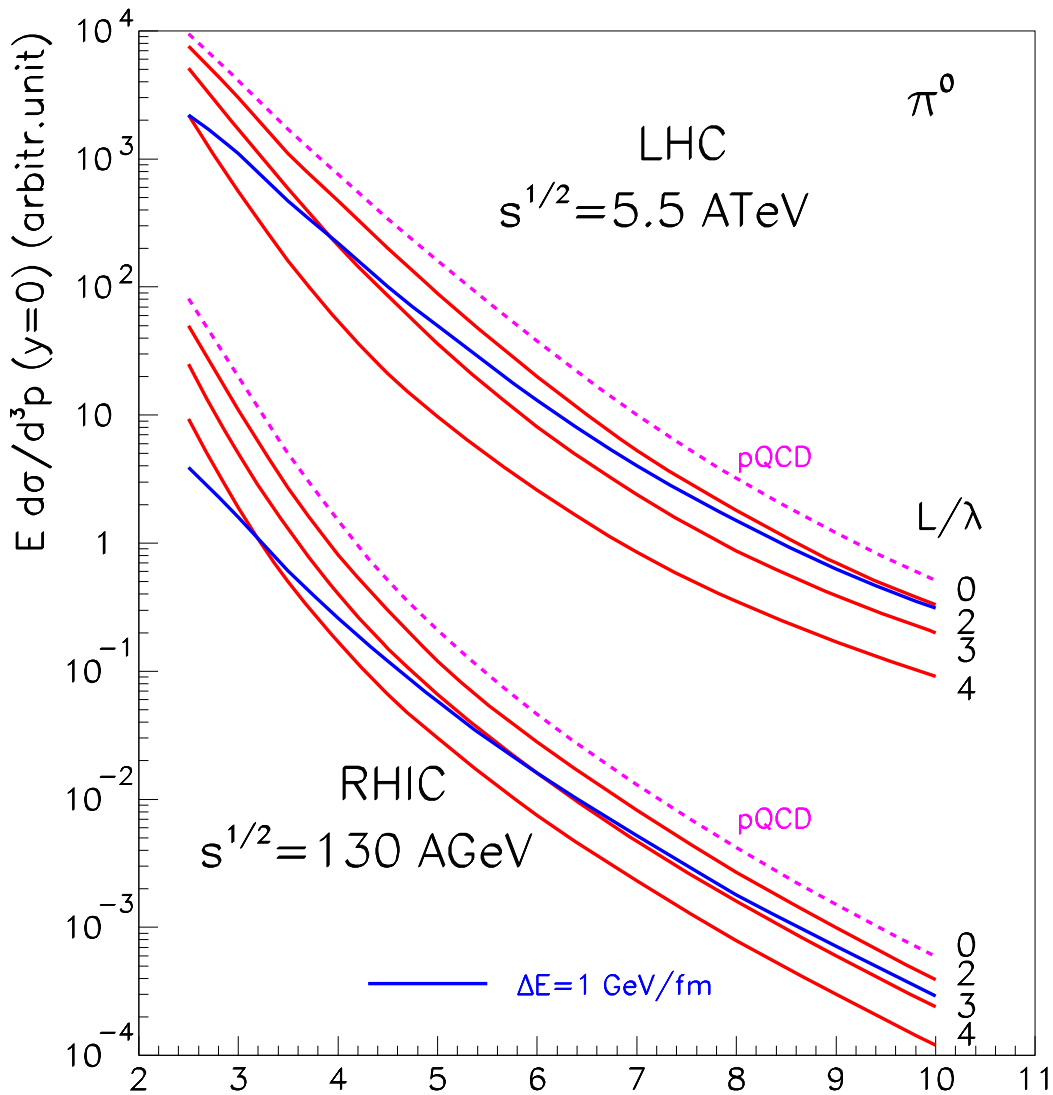
OPACITY EXPANSION: FAST CONVERGENCE !!



$\Rightarrow \Delta E_{GLV}^{(1)}$ is a good approximation;
 \Rightarrow Jet energy loss is small at CERN SPS;

Jet quenching on hadron yields - 1

- Au+Au collision at RHIC, $\sqrt{s} = 130$ AGeV.
- Pb+Pb collision at LHC, $\sqrt{s} = 5500$ AGeV.

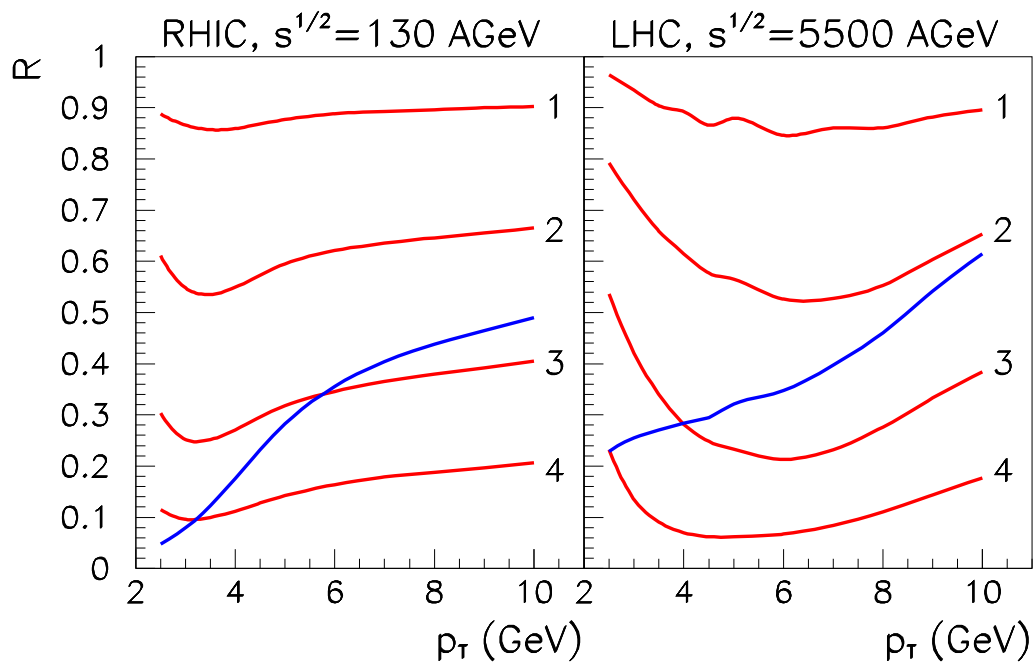


$L/\lambda = 0, 1, 2, 3, 4$ at $\mu = 0.5$ GeV, $\lambda = 1$ fm

Jet quenching on hadron yields - 2

- Au+Au collision at RHIC, $\sqrt{s} = 130$ AGeV.
- Pb+Pb collision at LHC, $\sqrt{s} = 5500$ AGeV.

”QUENCHED / NO QUENCH” ratios



Constant $\Delta E \implies$ Varying slope

Constant $\Delta E/E \implies$ Parallel shift

Slope analysis does not help

Compare to RHIC data:

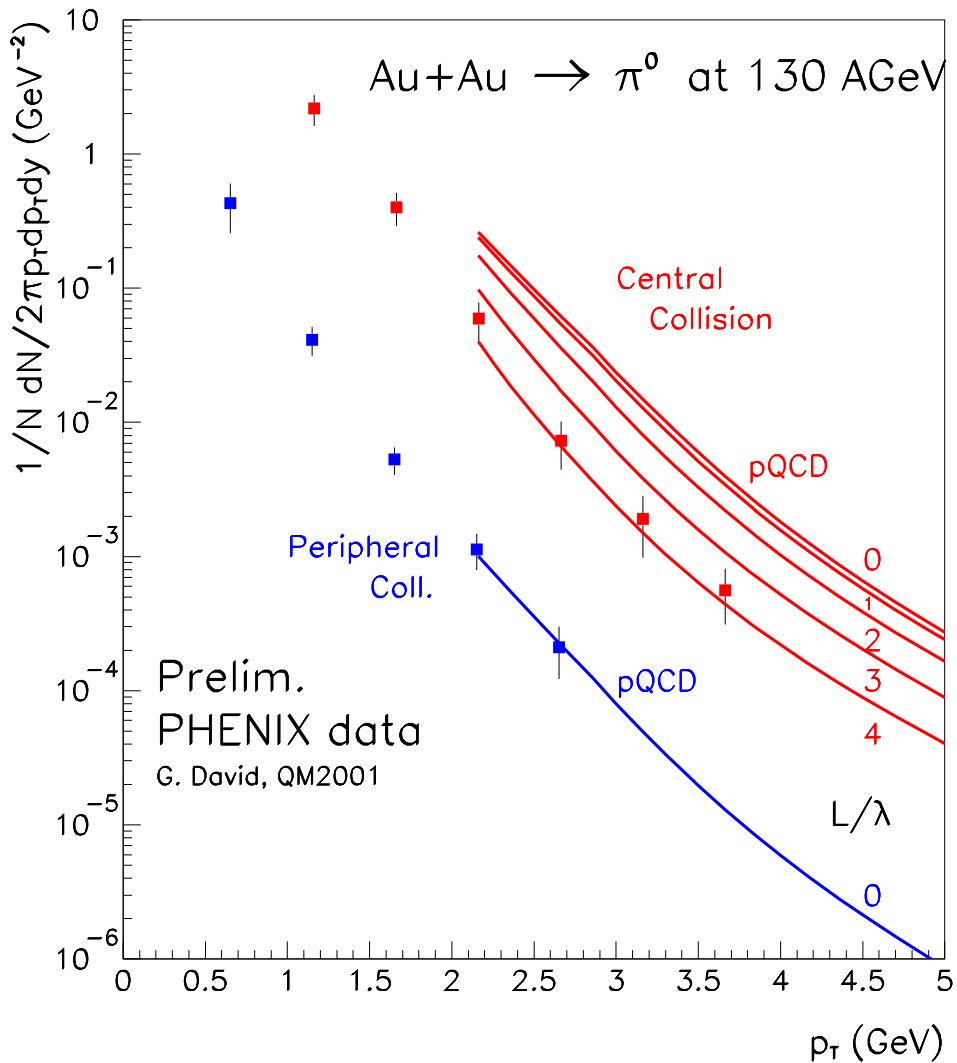
1. Total yields (π^0 in central and perif. collision)
2. Hadron ratios (K^+/π^+ , K^-/π^- ...)

Jet quenching on hadron yields - 3

PHENIX prel. results at $\sqrt{s} = 130$ GeV G. David's talk

Peripheral collisions (75-92 %; $\langle N_{coll} \rangle = 5.5$)

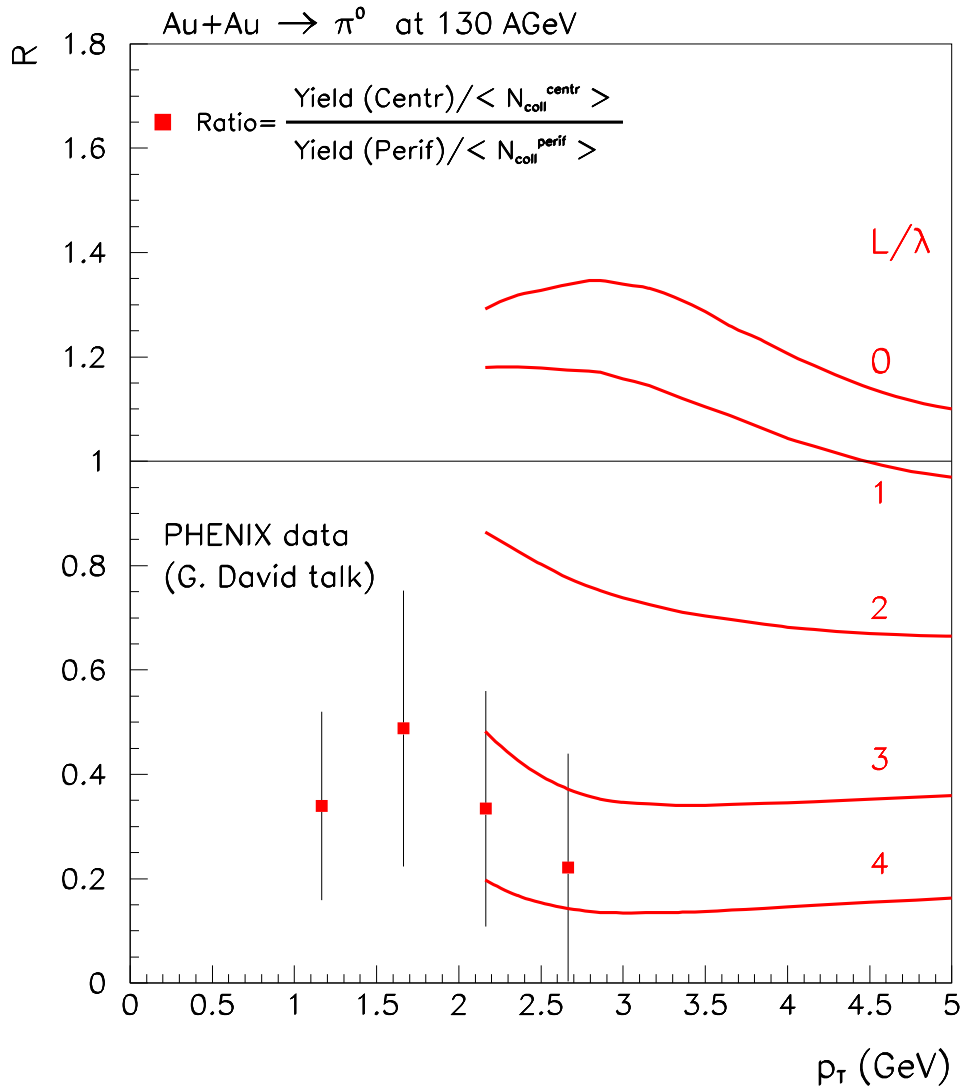
Central collisions (10 %; $\langle N_{coll} \rangle = 857$)



Factor 8-10 shift in central coll. \iff jet-quenching !!!!

Jet quenching on hadron yields - 4

Central/periferial ratios:

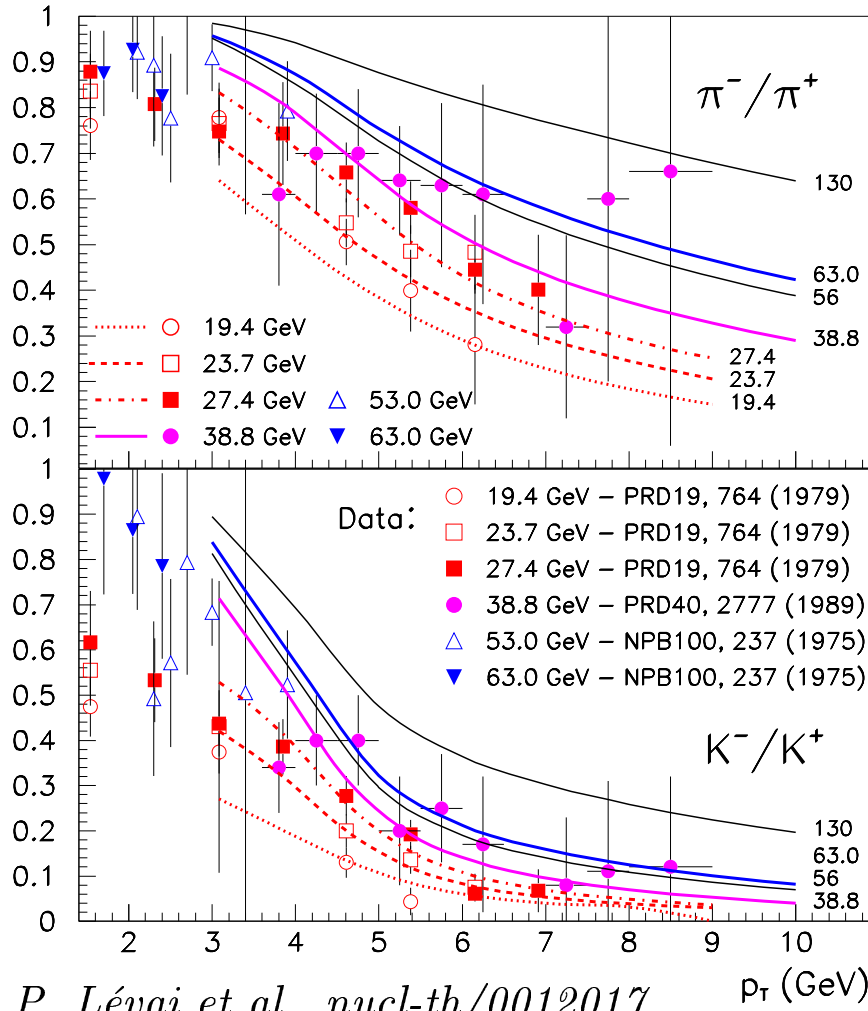


Hadron production at high p_T by pQCD - 1

Pion and Kaon production in pp collisions:
 \longrightarrow intrinsic k_T

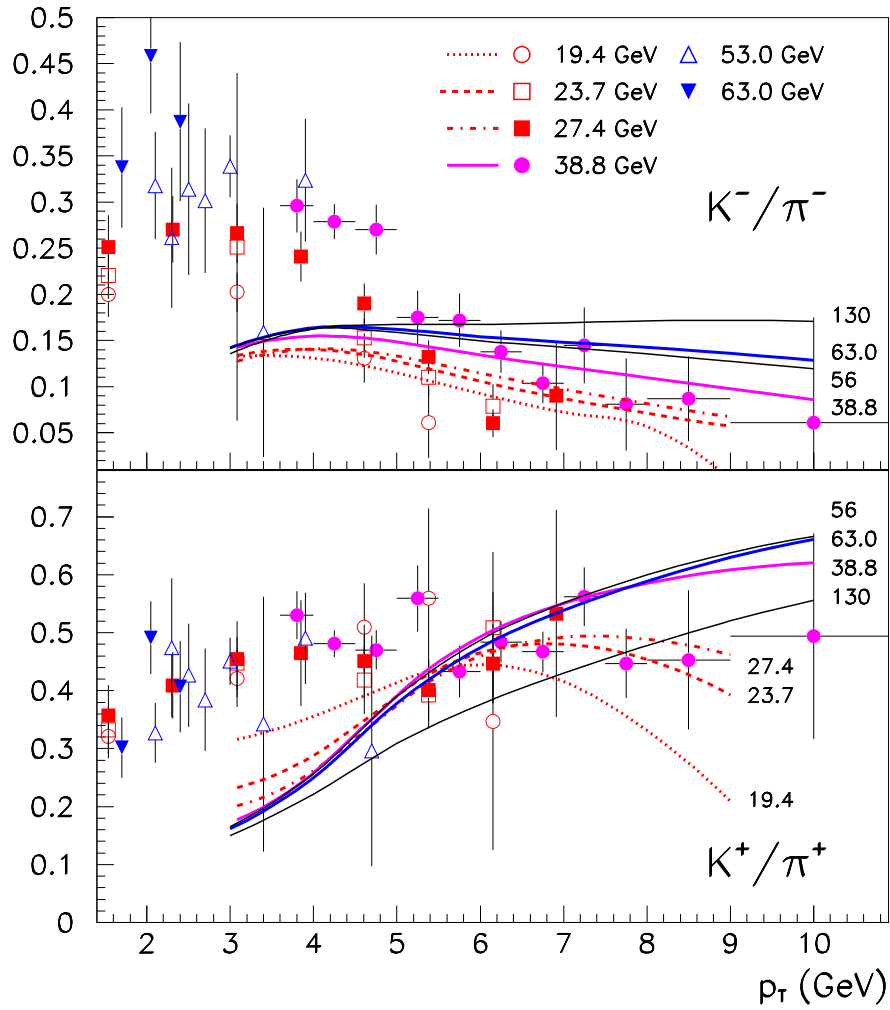
$$E_h \frac{d\sigma_h^{pp}}{d^3p} = \sum_{abcd} \int dx_1 dx_2 d^2k_{T,a} d^2k_{T,b} g(\vec{k}_{T,a}) g(\vec{k}_{T,b}) \times$$

$$f_{a/p}(x_1, Q^2) f_{b/p}(x_2, Q^2) \frac{d\sigma D_{h/c}(z_c, \hat{Q}^2)}{d\hat{t}} \frac{1}{\pi z_c}$$



Hadron production at high p_T by pQCD - 2

Pion and Kaon ratios in pp collisions:



Hadron ratios as probe of jet quenching - 1

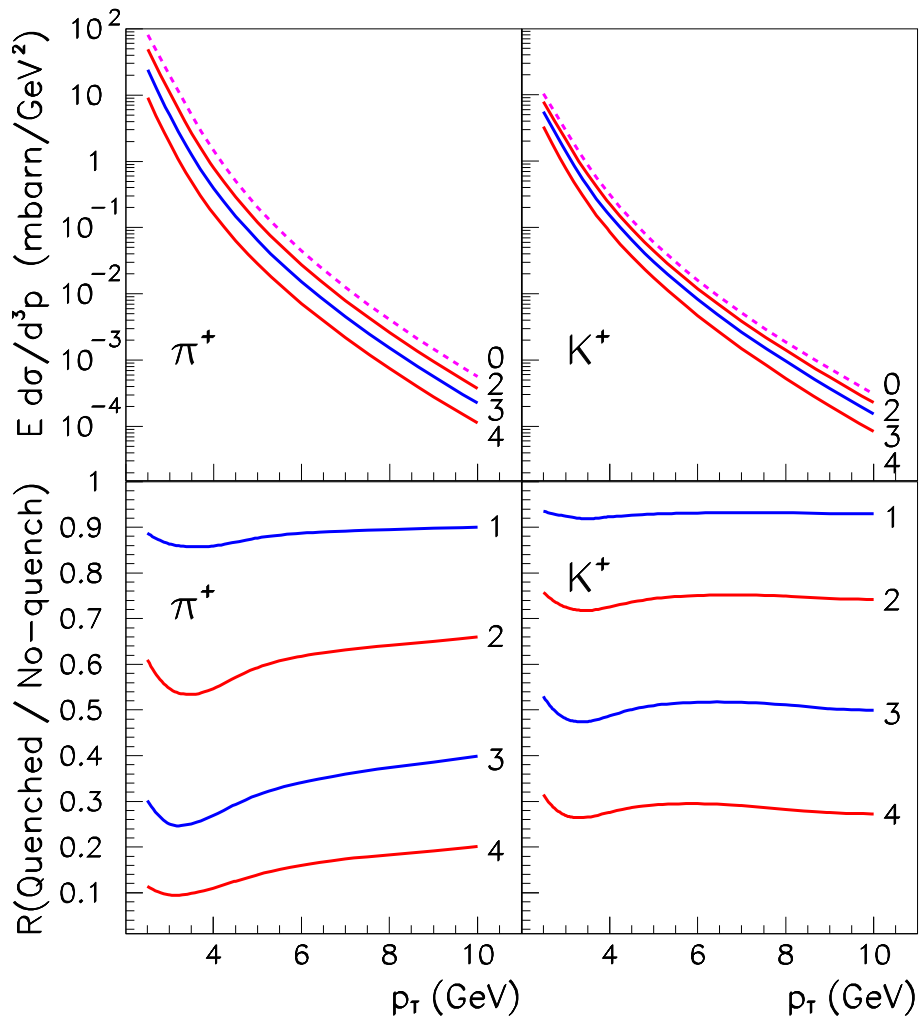
Hadron production in A+A collision (LO):

$$E_h \frac{d\sigma_h^{AA}}{d^3p} = B_{AA}^{(0)} \sum_{abcd} \int dx_1 dx_2 d^2k_{T,a} d^2k_{T,b} g(\vec{k}_{T,a}) g(\vec{k}_{T,b})$$

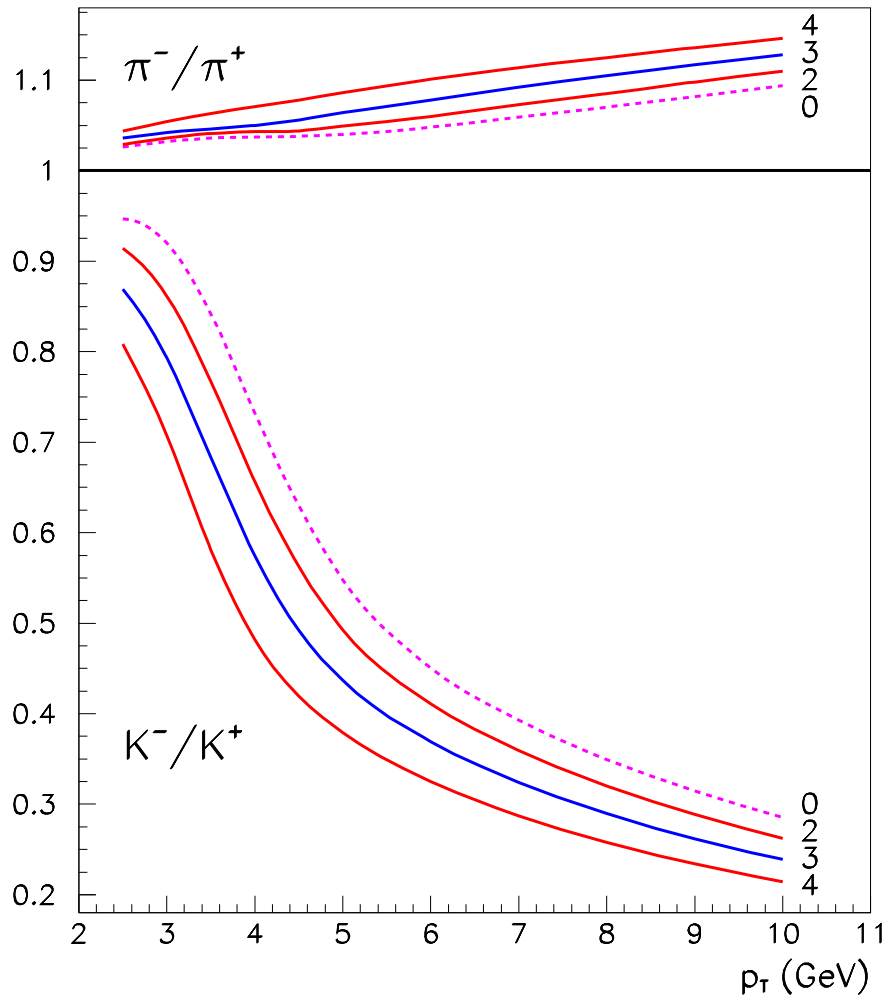
$$f_{a/A}(x_1, Q^2) f_{b/A}(x_2, Q^2) \frac{d\sigma}{d\hat{t}} \frac{z_c^*}{z_c} \frac{D_{h/c}(z_c^*, \hat{Q}^2)}{\pi z_c}$$

- $f_{a/A}(x_1, Q^2)$ in-medium PDF (shadowing).
- $B_{AA}^{(0)} = 2\pi \int_0^{b_{max}} b db T_{AA}(b)$
- $T_{AA}(b)$ is the thickness function
- $z_c^*/z_c \iff$ in-medium fragmentation function (renorm.)

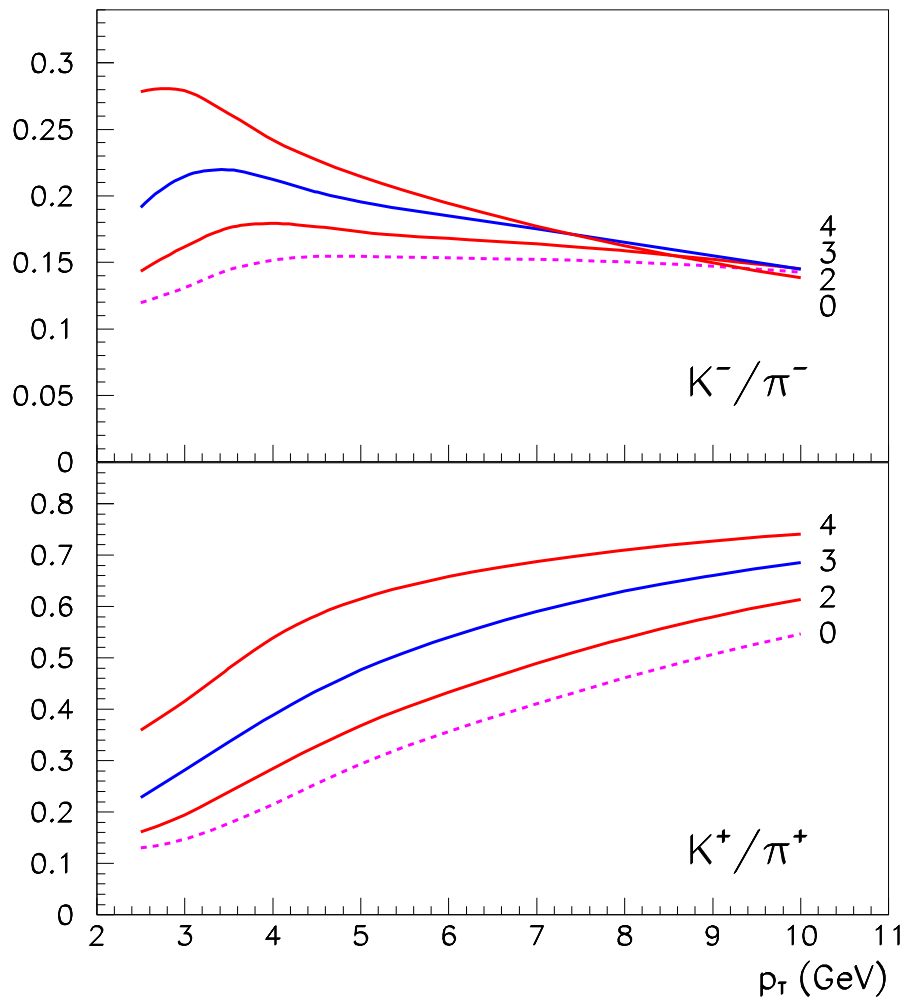
Hadron ratios as probe of jet quenching - 2



Hadron ratios as probe of jet quenching - 3



Hadron ratios as probe of jet quenching - 4



Summary

**0. PHENIX result indicates strong jet-quenching:
~ 6 – 8 fold suppression in the π^0 p_T -spectra
in central Au+Au collisions**

1. Suppression pattern understood:

- appropriate description of jet energy loss
GLV: thin matter, non-abelian case
- baseline pQCD predictions for AA collision
precise pQCD description in pp
multiscattering & shadowing in pA .

$\Rightarrow L/\lambda \approx 3 - 4$ in an effective static medium

expansion implies very high initial density

$$dN_g/dy \approx 1000$$

2. Flavour dependence of jet-quenching

K^+/π^+ , K^-/π^- ratios are sensitive to energy loss.

3. Energy loss in expanding matter

Gyulassy-Wang-Vitev (nucl-th/0012092)

\Rightarrow Poster P113 — I. Vitev