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# On the Use of Relaxation Parameters in Hybrid Smoothers 

Ulrike Meier Yang*


#### Abstract

The use of relaxation parameters in hybrid smoothers within algebraic multigrid (AMG) is analyzed both theoretically and practically. Relaxation parameters that are optimal under the assumptions of the theory are determined. The implementation of a procedure to automatically determine outer relaxation parameters for symmetric positive definite smoothers is described. Numerical results are presented, which show significant improvements over AMG with undamped hybrid smoothers.


## 1 Introduction

With the advent of large high performance computers with large number of processors, it has become necessary to design parallel algorithms of all sorts. Particular emphasis has been placed on the development of scalable algorithms, such as multigrid methods. With this in mind, the parallelization of algebraic multigrid, a method that can be applied to a linear system, $A x=b$, without additional knowledge, such as the underlying finite elements or a grid, has become very important. Algebraic multigrid (AMG) proceeds by determining a subset of the original degrees of freedom through a coarsening algorithm, a restriction operator that transfers vectors from the fine space to the coarse space, and an interpolation operator that transfers vectors from the coarser space to the finer space. One important component of AMG is the smoother. A good smoother will reduce the oscillatory error components, whereas the "smooth" error is transferred to the coarser grids. Although the classical approach of AMG focused mainly on the Gauß-Seidel method [7], the use of other iterative solvers as smoothers has been considered [3, 2].

Gauß-Seidel has proven to be an effective smoother for many problems, however its main disadvantage is its sequential nature. On the other hand, highly parallel smoothers such as Jacobi or block-Jacobi often fail, unless an appropriate smoothing parameter is used, and even then their convergence is often slow. Additionally, the user is faced with the challenge on how to choose an appropriate smoothing parameter. Many efforts to parallelize Gauß-Seidel have been made. Possible variants include the use of multi-coloring techniques [1] or hybrid schemes [6]. Multi-coloring techniques are a nuisance to implement and can be inefficient, if too many colors are involved (which is most likely to happen on the coarser levels of AMG). Hybrid schemes use an iterative method, e.g. Gauß-Seidel on each processor, but update in a Jacobi-like approach across boundaries. They are equivalent to block Jacobi methods that use one or more iterations of a smoother within each block instead of a direct solve. Clearly, this approach is very suitable for parallel processing, however, just like the block Jacobi method, it often requires a suitable smoothing parameter for convergence.

[^0]In this paper, we investigate the use of relaxation parameters in hybrid smoothers. There are two types of relaxation parameters: the smoothing parameter, mentioned above, which will be denoted the outer relaxation parameter, $\omega_{J}$; and the so-called inner relaxation parameters, $\omega_{i}$, which occur, if we smooth locally, on processor $i$, using SOR or its symmetric variant, SSOR. For both cases, the question is how to determine good parameters. Additionally, since one deals with a new system on each level of AMG, the development of an automatic procedure to determine such parameters is important. Since the outer relaxation parameter affects the matrix across all processors, it appears that this parameter would be more crucial in improving convergence (or leading to convergence in cases of divergence). Therefore our main focus will be on the determination of an optimal $\omega_{J}$. However, this paper also contains some results on the use of inner relaxation parameters.

In Section 2, we give some basic definitions. In Sections 3 through 5 we present conditions, under which the smoothing properties are fulfilled. Section 4 focuses specifically on the outer relaxation parameters and presents the determination of optimal parameters. Section 5 analyzes the use of inner relaxation parameters. Section 6 describes a procedure to determine outer relaxation parameters automatically, and in Section 7, we present numerical results that show that this approach can lead to significant improvements or even convergence in cases, for which AMG with an undamped hybrid smoother does not converge.

## 2 Definitions

Since our goal is to solve the linear system $A x=b$ on a parallel computer with $p$ processors, we partition the linear system as follows:

$$
A=\left(\begin{array}{llll}
A_{11} & \cdot & \cdot & A_{1 p}  \tag{1}\\
\cdot & \cdot & \cdot & \cdot \\
\cdot & \cdot & \cdot & \cdot \\
A_{p 1} & \cdot & \cdot & A_{p p}
\end{array}\right)\left(\begin{array}{l}
x_{1} \\
\cdot \\
\cdot \\
x_{p}
\end{array}\right)=\left(\begin{array}{l}
b_{1} \\
\cdot \\
\cdot \\
b_{p}
\end{array}\right)
$$

A general definition of a smoother based on a splitting $Q-(Q-A)$ of $A$ is

$$
\begin{equation*}
Q u_{n+1}=b+(Q-A) u_{n} \tag{2}
\end{equation*}
$$

where $Q$ can be any nonsingular matrix. For example, with Gauß-Seidel $Q=D-L$, where $D$ is the diagonal matrix with the diagonal of $A$ and $-L$ the strict lower triangular part of $A$, while for Jacobi $Q=D$. Since we are interested in parallel smoothers, we will only consider splittings of the form

$$
Q=\left(\begin{array}{cccc}
Q_{1} & & &  \tag{3}\\
& Q_{2} & & \\
& & \cdot & \\
& & & Q_{p}
\end{array}\right)
$$

This is equivalent to performing any iterative solver such as Gauß-Seidel or block Gauß-Seidel, etc, locally on each processor, but updating the unknowns that are located on the neighbor processors only after each iteration step. Specific examples for hybrid smoothers are hybrid Gauß-Seidel with $Q_{k}=D_{k}-L_{k}$, or hybrid symmetric Gauß-Seidel with $Q_{k}=\left(D_{k}-L_{k}\right) D_{k}^{-1}\left(D_{k}-L_{k}^{T}\right)$, where $D_{k}$ is the diagonal matrix with the diagonal of $A_{k k}$ and $L_{k}$ the strictly lower triangular part of $-A_{k k}$.

As mentioned in the previous section, there are two types of relaxation parameters, the outer relaxation parameter $\omega_{J}$, and the inner relaxation parameters $\omega_{i}, i=1, \ldots, p$. Therefore consider the following smoothing matrix with outer relaxation parameter $\omega_{J}$

$$
Q=\frac{1}{\omega_{J}} \tilde{Q}=\frac{1}{\omega_{J}}\left(\begin{array}{llll}
\tilde{Q}_{1} & &  \tag{4}\\
& \cdot & & \\
& & \\
& & & \tilde{Q}_{p}
\end{array}\right)
$$

Inner relaxation parameters, $\omega_{k}, k=1, \ldots, p$ occur in hybrid SOR with $Q$ defined as above and

$$
\begin{equation*}
\tilde{Q}_{k}=\frac{1}{\omega_{k}} D_{k}-L_{k} \tag{5}
\end{equation*}
$$

as well as the corresponding symmetric variant, hybrid SSOR, with

$$
\begin{equation*}
\tilde{Q}_{k}=\frac{\omega_{k}}{2-\omega_{k}}\left(\frac{1}{\omega_{k}} D_{k}-L_{k}\right) D_{k}^{-1}\left(\frac{1}{\omega_{k}} D_{k}-L_{k}^{T}\right) \tag{6}
\end{equation*}
$$

Further notations needed include the smallest eigenvalue of a matrix $A, \lambda_{\min }(A)$, and the spectral radius of $A, \rho(A)$, which is defined as the absolute largest eigenvalue of $A$.

## 3 Smoothing Properties

Denote by $A^{(m)}$ the operator on the $m$-th level of AMG, $P^{(m)}$ the interpolation operator that interpolates from the $m+1$ st level to the $m$ th level, and $R^{(m)}$ the restriction operator, that transfers from the $m$ th to the $m+1$ st level, which in general, if $A^{(m)}$ is symmetric, is defined as $R^{(m)}=\left(P^{(m)}\right)^{T}$.

Then the coarse grid correction operator is given by

$$
T^{(m)}=I-P^{(m)}\left(A^{(m+1)}\right)^{-1} R^{(m)} A^{(m)}
$$

We define the smoothing operator on the $m$ th level of AMG as

$$
G^{(m)}=I-\left(Q^{(m)}\right)^{-1} A^{(m)},
$$

where $Q^{(m)}$ is the matrix defined by the relaxation process $Q^{(m)} u_{n+1}=b^{m}+\left(Q^{(m)}-\right.$ $\left.A^{(m)}\right) u_{n}$.

Important conditions for convergence of algebraic multigrid methods are the smoothing properties

$$
\begin{array}{rc}
\left\|G^{(m)} e^{m}\right\|_{1}^{2} \leq\left\|e^{m}\right\|_{1}^{2}-\sigma_{1}\left\|e^{m}\right\|_{2}^{2}, & \sigma_{1}>0, \\
\left\|G^{(m)} e^{m}\right\|_{1}^{2} \leq\left\|e^{m}\right\|_{1}^{2}-\sigma_{2}\left\|G^{(m)} e^{m}\right\|_{2}^{2}, & \sigma_{2}>0,  \tag{8}\\
\text { for any } e^{m} \\
e^{m}
\end{array}
$$

with the following norms

$$
\begin{equation*}
\|x\|_{1}=\left(x^{T} A^{(m)} x\right)^{\frac{1}{2}}, \quad\|x\|_{2}=\left(x^{T}\left(A^{(m)}\right)^{T}\left(D^{(m)}\right)^{-1} A^{(m)} x\right)^{\frac{1}{2}} \tag{9}
\end{equation*}
$$

where $D^{(m)}$ is the diagonal part of $A^{(m)}$. Here (7) refers to postsmoothing, (8) to presmoothing. In conjunction with the approximation property

$$
\begin{equation*}
\left\|T^{(m)} e^{m}\right\|_{1}^{2} \leq \beta\left\|e^{m}\right\|_{2}^{2}, \quad \text { for any } e^{m} \tag{10}
\end{equation*}
$$

either (7) or (8) imply two-level convergence with the convergence factor bounded above by $\sqrt{1-\delta_{1}}$ for (7) and by $1 / \sqrt{1+\delta_{2}}$ for (8), with $\delta_{1}=\sigma_{1} / \beta$ and $\delta_{2}=\sigma_{2} / \beta$. This shows that larger $\sigma_{1}$ and $\sigma_{2}$ will lead to better smoothing and ultimately better convergence. For a more detailed discussion on convergence see [7, 9].

Obviously, the approximation property is determined by the choice of interpolation and restriction, and therefore of no concern for us in this context, in which we only focus on smoothing.

Now according to [7], (7) and (8) are equivalent to the following inequalities, which are somewhat easier to deal with (for simplicity, the indices $m$ will be omitted in the remainder of the paper)

$$
\begin{align*}
\sigma_{1} e^{T} Q^{T} D^{-1} Q e & \leq e^{T}\left(Q^{T}+Q-A\right) e  \tag{11}\\
\sigma_{2} e^{T}(Q-A)^{T} D^{-1}(Q-A) e & \leq e^{T}\left(Q^{T}+Q-A\right) e \tag{12}
\end{align*}
$$

If $A$ is a symmetric positive definite matrix, then the matrix $Q+Q^{T}-A$ is symmetric positive definite if and only if $A=Q-(Q-A)$ is a splitting that leads to a convergent iterative method, i.e. $\rho\left(I-Q^{-1} A\right)<1$, see [10]. Using this fact it is easy to show that both smoothing properties can always be fulfilled.

Proposition 1 Assume $A$ and $D$ are symmetric positive definite and $Q$ is a matrix, for which $\rho\left(I-Q^{-1} A\right)<1$. If

$$
\begin{equation*}
\sigma_{1} \leq \frac{\lambda_{\min }\left(Q+Q^{T}-A\right)}{\rho\left(Q^{T} D^{-1} Q\right)} \tag{13}
\end{equation*}
$$

then the first smoothing property (11) holds. If

$$
\begin{equation*}
\sigma_{2} \leq \frac{\lambda_{\min }\left(Q+Q^{T}-A\right)}{\rho\left((Q-A)^{T} D^{-1}(Q-A)\right)} \tag{14}
\end{equation*}
$$

then the second smoothing property (12) is fulfilled.
Proof:
Since $\rho\left(I-Q^{-1} A\right)<1, Q+Q^{T}-A$ is symmetric positive definite, and consequently $\lambda_{\min }\left(Q+Q^{T}-A\right)>0$. The matrices $Q^{T} D^{-1} Q$ and $(Q-A)^{T} D^{-1}(Q-A)$ are symmetric positive semidefinite and thus have real nonnegative eigenvalues. Therefore the following inequality holds for any $e$ with $Q e \neq 0$

$$
\begin{aligned}
\frac{e^{T}\left(Q+Q^{T}-A\right) e}{e^{T} Q^{T} D^{-1} Q e} & =\frac{e^{T}\left(Q+Q^{T}-A\right) e}{e^{T} e} \frac{e^{T} e}{e^{T} Q^{T} D^{-1} Q e} \\
& \geq \frac{\lambda_{\min }\left(Q+Q^{T}-A\right)}{\rho\left(Q^{T} D^{-1} Q\right)} \geq \sigma_{1}>0
\end{aligned}
$$

Note that if $Q e=0$, (11) holds for any $\sigma_{1}$. Analogously, for $\sigma_{2}$ : for any vector $e$ with $(Q-A) e \neq 0$

$$
\begin{aligned}
\frac{e^{T}\left(Q+Q^{T}-A\right) e}{e^{T}(Q-A)^{T} D^{-1}(Q-A) e} & =\frac{e^{T}\left(Q+Q^{T}-A\right) e}{e^{T} e} \frac{e^{T} e}{e^{T}(Q-A)^{T} D^{-1}(Q-A) e} \\
& \geq \frac{\lambda_{\min }\left(Q+Q^{T}-A\right)}{\rho\left((Q-A)^{T} D^{-1}(Q-A)\right)} \geq \sigma_{2}>0
\end{aligned}
$$

Note that if $(Q-A) e=0,(12)$ holds for any $\sigma_{2}$.
q.e.d.

Thus, if the underlying iterative scheme is convergent, the smoothing properties are fulfilled. Further, it is easy to come up with a sample vector $e$, for which the properties are not fulfilled, if the method does not converge. The proposition also explains, why hybrid smoothers fail: for many problems, as can be seen in Section $7, \rho\left(I-Q^{-1} A\right)>1$, or, equivalently, $\rho\left(Q^{-1} A\right)>2$.

However, the estimated bounds for $\sigma_{1}$ and $\sigma_{2}$ are not very useful if one wants to determine good relaxation parameters. Therefore, we investigate symmetric positive definite $Q$ and derive more meaningful bounds for $\sigma_{1}$ and $\sigma_{2}$. Note that the best values for $\sigma_{1}$ and $\sigma_{2}$ in all theorems are obtained when they are taken to be equal to their upper limit. Smaller values provide less sharp bounds.

Theorem 2 Assume that $A, Q$ and $D$ are symmetric positive definite and

$$
\rho\left(Q^{-1} A\right)<2
$$

Assume

$$
\begin{equation*}
\sigma_{1} \leq \frac{2-\rho\left(Q^{-1} A\right)}{\rho\left(D^{-1} Q\right)} \tag{15}
\end{equation*}
$$

then the smoothing property

$$
\sigma_{1} e^{T} Q^{T} D^{-1} Q e \leq e^{T}\left(Q+Q^{T}-A\right) e
$$

is fulfilled.
Proof: Assume $e \neq 0$. Note that for $e=0$ the smoothing property is fulfilled for any $\sigma_{1}$. Since $Q$ is symmetric positive definite, and since for any positive definite matrix $C, \rho(C A)=\rho\left(C^{\frac{1}{2}} A C^{\frac{1}{2}}\right)$,

$$
\begin{aligned}
\frac{e^{T}\left(Q^{T}+Q-A\right) e}{e^{T} Q^{T} D^{-1} Q e} & =\frac{e^{T} Q^{\frac{1}{2}}\left(2 I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right) Q^{\frac{1}{2}} e}{e^{T} Q e} \frac{e^{T} Q e}{e^{T} Q^{\frac{1}{2}}\left(Q^{\frac{1}{2}} D^{-1} Q^{\frac{1}{2}}\right) Q^{\frac{1}{2}} e} \\
& \geq \frac{2-\rho\left(Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right)}{\rho\left(Q^{\frac{1}{2}} D^{-1} Q^{\frac{1}{2}}\right)} \\
& =\frac{2-\rho\left(Q^{-1} A\right)}{\rho\left(D^{-1} Q\right)}>0
\end{aligned}
$$

because $\rho\left(Q^{-1} A\right)<2$. Thus if $\sigma_{1}$ is chosen as in (15), (11) is fulfilled.
q.e.d.

A similar result is obtained for the second smoothing property (12).
Theorem 3 Assume that $A, Q$ and $D$ are symmetric positive definite and

$$
\rho\left(Q^{-1} A\right)<2
$$

Assume that

$$
\begin{equation*}
\sigma_{2} \leq \frac{2-\rho\left(Q^{-1} A\right)}{\rho\left(D^{-1} Q\right)\left[\rho\left(I-Q^{-1} A\right)\right]^{2}} \tag{16}
\end{equation*}
$$

Then the smoothing property

$$
\sigma_{2} e^{T}(Q-A)^{T} D^{-1}(Q-A) e \leq e^{T}\left(Q+Q^{T}-A\right) e
$$

is fulfilled.
Proof: Assume $(Q-A) e \neq 0$. Note that for $(Q-A) e=0$ the smoothing property is fulfilled for any $\sigma_{2}$. First let us consider $\rho\left(\left(I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right)^{T} Q^{\frac{1}{2}} D^{-1} Q^{\frac{1}{2}}(I-\right.$ $\left.\left.Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right)\right)$. Also, $\|\cdot\|$ denotes here the spectral norm or 2-norm.

$$
\begin{aligned}
\rho\left(\left(I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right) Q^{\frac{1}{2}} D^{-1} Q^{\frac{1}{2}}\left(I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right)\right) & =\left\|D^{-\frac{1}{2}} Q^{\frac{1}{2}}\left(I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right)\right\|^{2} \\
& \leq\left\|D^{-\frac{1}{2}} Q^{\frac{1}{2}}\right\|^{2}\left\|I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right\|^{2} \\
& =\rho\left(Q^{\frac{1}{2}} D^{-1} Q^{\frac{1}{2}}\right)\left[\rho\left(I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right)\right]^{2} \\
& =\rho\left(D^{-1} Q\right)\left[\rho\left(I-Q^{-1} A\right)\right]^{2} .
\end{aligned}
$$

Using this result, we get

$$
\begin{aligned}
\frac{e^{T}\left(Q^{T}+Q-A\right) e}{e^{T}(A-Q)^{T} D^{-1}(A-Q) e} & =\frac{e^{T} Q^{\frac{1}{2}}\left(2 I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right) Q^{\frac{1}{2}} e}{e^{T} Q^{\frac{1}{2}}\left(I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right) Q^{\frac{1}{2}} D^{-1} Q^{\frac{1}{2}}\left(I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right) Q^{\frac{1}{2}} e} \\
& \geq \frac{2-\rho\left(Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right)}{\rho\left(\left(I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right)^{T} Q^{\frac{1}{2}} D^{-1} Q^{\frac{1}{2}}\left(I-Q^{-\frac{1}{2}} A Q^{-\frac{1}{2}}\right)\right)} \\
& \geq \frac{2-\rho\left(Q^{-1} A\right)}{\rho\left(D^{-1} Q\right)\left[\rho\left(I-Q^{-1} A\right)\right]^{2}}>0,
\end{aligned}
$$

because $\rho\left(Q^{-1} A\right)<2$.
q.e.d.

## 4 Determination of an Optimal Outer Relaxation Parameter

In this section, outer relaxation parameters are determined that are optimal in the sense that they maximize the upper bounds for $\sigma_{1}$ and $\sigma_{2}$ that have been derived in the previous section. Since the conditions derived in the previous section are just sufficient conditions for satisfying the smoothing properties, this does not prove optimality in the absolute sense. Nevertheless, there is empirical evidence that the use of this theory leads to very good results in practice, as can be seen in Section 7.

Using (4), (15) can be expressed as

$$
\sigma_{1} \leq \frac{\omega_{J}\left(2-\omega_{J} \rho\left(\tilde{Q}^{-1} A\right)\right)}{\max _{1 \leq k \leq p} \rho\left(D_{k}^{-1} \tilde{Q}_{k}\right)}
$$

and (16) as

$$
\sigma_{2} \leq \frac{\omega_{J}\left(2-\omega_{J} \rho\left(\tilde{Q}^{-1} A\right)\right)}{\max _{1 \leq k \leq p} \rho\left(D_{k}^{-1} \tilde{Q}_{k}\right)\left[\rho\left(I-\omega_{J} \tilde{Q}^{-1} A\right)\right]^{2}}
$$

Since $\sigma_{1}$ and $\sigma_{2}$ are positive, one obtains the following condition for $\omega_{J}$ :

$$
0<\omega_{J}<\frac{2}{\rho\left(\tilde{Q}^{-1} A\right)}
$$

Note that in the case of inner relaxation parameters $\omega_{k}$, for $1 \leq k \leq p$, the value of $\omega_{J}$ depends on the choice of these parameters.

In order to maximize the upper bound for $\sigma_{1}$ in (11), one must choose

$$
\begin{equation*}
\omega_{J}=\frac{1}{\rho\left(\tilde{Q}^{-1} A\right)}, \tag{17}
\end{equation*}
$$

which leads to

$$
\sigma_{1} \leq \frac{1}{\rho\left(D^{-1} \tilde{Q}\right) \rho\left(\tilde{Q}^{-1} A\right)}
$$

Consider now the second smoothing inequality (12). Note that since $\omega_{J}>0$ and both Q and A are symmetric positive definite, so are $\tilde{Q}$ and $\tilde{Q}^{-\frac{1}{2}} A \tilde{Q}^{-\frac{1}{2}}$. Therefore

$$
\begin{aligned}
\rho\left(I-\omega_{J} \tilde{Q}^{-1} A\right) & =\rho\left(I-\omega_{J} \tilde{Q}^{-\frac{1}{2}} A \tilde{Q}^{-\frac{1}{2}}\right) \\
& = \begin{cases}1-\omega_{J} \lambda_{\min }\left(\tilde{Q}^{-1} A\right) & \text { for } \quad 0<\omega_{J}<\frac{2}{\rho\left(\tilde{Q}^{-1} A\right)+\lambda_{\min }\left(\tilde{Q}^{-1} A\right)} \\
\left|1-\omega_{J} \rho\left(\tilde{Q}^{-1} A\right)\right| & \text { for } \quad \omega_{J} \geq \frac{2}{\rho\left(\tilde{Q}^{-1} A\right)+\lambda_{\min }\left(\tilde{Q}^{-1} A\right)} .\end{cases}
\end{aligned}
$$

The optimal $\omega_{J}$ can now be determined by maximizing

$$
\psi(\omega)= \begin{cases}\frac{\omega\left(2-\omega \rho\left(\tilde{Q}^{-1} A\right)\right)}{\left(1-\omega \lambda_{\min }\left(\tilde{Q}^{-1} A\right)\right)^{2}} & \text { for } \quad 0<\omega_{J}<\frac{2}{\rho\left(\tilde{Q}^{-1} A\right)+\lambda_{\min }\left(\tilde{Q}^{-1} A\right)} \\ \frac{\omega\left(2-\omega \rho\left(\tilde{Q}^{-1} A\right)\right)}{\left(1-\omega \rho\left(\tilde{Q}^{-1} A\right)\right)^{2}} & \text { for } \quad \omega_{J} \geq \frac{2}{\rho\left(\tilde{Q}^{-1} A\right)+\lambda_{\min }\left(\tilde{Q}^{-1} A\right)} .\end{cases}
$$

with respect to $\omega$.
The solution to maximizing the first term is

$$
\begin{equation*}
\omega_{J}=\frac{1}{\rho\left(\tilde{Q}^{-1} A\right)-\lambda_{\min }\left(\tilde{Q}^{-1} A\right)} . \tag{18}
\end{equation*}
$$

The second term is decreasing in the considered range, consequently it is maximal for

$$
\begin{equation*}
\omega=\frac{2}{\lambda_{\min }\left(\tilde{Q}^{-1} A\right)+\rho\left(\tilde{Q}^{-1} A\right)} . \tag{19}
\end{equation*}
$$

For $\rho\left(\tilde{Q}^{-1} A\right)>3 \lambda_{\min }\left(\tilde{Q}^{-1} A\right),(18)$ is larger than (19) and maximizes $\psi(\omega)$ in the considered range. It leads to

$$
\sigma_{2} \leq \frac{1}{\rho\left(D^{-1} \tilde{Q}\right)\left(\rho\left(\tilde{Q}^{-1} A\right)-2 \lambda_{\min }\left(\tilde{Q}^{-1} A\right)\right)}
$$

Since in general for finite element problems $\lambda_{\min }\left(\tilde{Q}^{-1} A\right)$ is small compared to the largest eigenvalue, the best choice for the first smoothing property (17) is also a good choice for the second smoothing property with

$$
\sigma_{2} \leq \frac{\rho\left(\tilde{Q}^{-1} A\right)}{\rho\left(D^{-1} \tilde{Q}\right)\left(\rho\left(\tilde{Q}^{-1} A\right)-\lambda_{\min }\left(\tilde{Q}^{-1} A\right)\right)^{2}}
$$

If $\rho\left(\tilde{Q}^{-1} A\right)<3 \lambda_{\min }\left(\tilde{Q}^{-1} A\right), \psi(\omega)$ is maximal for (19) with

$$
\sigma_{2} \leq \frac{4 \lambda_{\min }\left(\tilde{Q}^{-1} A\right)}{\rho\left(D^{-1} \tilde{Q}\right)\left(\rho\left(\tilde{Q}^{-1} A\right)-\lambda_{\min }\left(\tilde{Q}^{-1} A\right)\right)^{2}}
$$

Clearly if

$$
4 \lambda_{\min }\left(\tilde{Q}^{-1} A\right)<\rho\left(\tilde{Q}^{-1} A\right)
$$

(17) is a better choice for $\omega_{J}$ than (19). Note that (19) minimizes $\rho\left(I-\omega \tilde{Q}^{-1} A\right)$ and leads therefore to the fastest convergent splitting, but does not necessarily lead to the best smoother!

## 5 Analysis of the Inner Relaxation Parameter

Inner relaxation parameters occur when one uses SOR or SSOR locally on each processor. Determining a best choice for the inner relaxation parameter in SSOR is a very difficult task. Even in the case of only one processor, one needs to analyze the following complicated functions

$$
\begin{aligned}
\phi_{1}(\omega) & =\frac{\omega(2-\omega)\left(2-\omega(2-\omega) \rho\left(\tilde{Q}(\omega)^{-1} A\right)\right)}{\rho\left(D^{-1} \tilde{Q}(\omega)\right)} \\
\phi_{2}(\omega) & =\frac{\omega(2-\omega)\left(2-\omega(2-\omega) \rho\left(\tilde{Q}(\omega)^{-1} A\right)\right)}{\rho\left(D^{-1} \tilde{Q}(\omega)\right)\left(\rho\left(I-\omega(2-\omega) \tilde{Q}(\omega)^{-1} A\right)\right)^{2}}
\end{aligned}
$$

Although there are some interesting results on the choice of relaxation parameters for SSOR as an iterative solver in [10], these results do not transfer to smoothers, and applying the same techniques to analyze the parameter gives nonconclusive results. However, one can derive some interesting results for SOR and hybrid SOR.

Lemma 4 Assume $A$ is symmetric positive definite, $Q=\frac{1}{\omega} D-L$, where $D$ is the diagonal and $-L$ the lower triangular part of $A$, and $0<\omega<2$. Assume also that

$$
\begin{equation*}
\sigma_{1} \leq \frac{(2-\omega) \omega}{\left(1+\omega \gamma^{-}\right)\left(1+\omega \gamma^{+}\right)} \tag{20}
\end{equation*}
$$

where

$$
\begin{aligned}
& \gamma^{-}=\left\|D^{-1} L\right\| \\
& \gamma^{+}=\left\|D^{-1} L^{T}\right\|
\end{aligned}
$$

and $\|$.$\| denotes any vector induced matrix norm. Then$

$$
\sigma_{1} e^{T} Q^{T} D^{-1} Q e \leq e^{T}\left(Q^{T}+Q-A\right) e
$$

The upper bound for $\sigma_{1}$ is maximal, if

$$
\omega=\frac{\sqrt{\left(2 \gamma^{+}+1\right)\left(2 \gamma^{-}+1\right)}-1}{\gamma^{+}+\gamma^{-}+2 \gamma^{-} \gamma^{+}} \leq 1
$$

Proof:
Since $Q=\frac{1}{\omega} D-L$,

$$
Q+Q^{T}-A=\frac{2}{\omega} D-L-L^{T}-D+L+L^{T}=\left(\frac{2}{\omega}-1\right) D
$$

Now

$$
e^{T} Q^{T} D^{-1} Q e \leq \rho\left(D^{-1} Q^{T} D^{-1} Q\right) e^{T} D e
$$

and

$$
\begin{aligned}
\rho\left(D^{-1} Q^{T} D^{-1} Q\right) & \leq\left\|D^{-1} Q^{T}\right\|\left\|D^{-1} Q\right\| \\
& =\left\|\frac{1}{\omega} I-D^{-1} L^{T}\right\|\left\|\frac{1}{\omega} I-D^{-1} L\right\| \\
& \leq \frac{1}{\omega^{2}}\left(1+\omega \gamma^{+}\right)\left(1+\omega \gamma^{-}\right)
\end{aligned}
$$

This leads to the following inequality for $\sigma_{1}$

$$
\sigma_{1} \leq \psi_{1}(\omega)=\frac{(2-\omega) \omega}{\left(1+\omega \gamma^{-}\right)\left(1+\omega \gamma^{+}\right)}
$$

Now,

$$
\frac{d \psi_{1}(\omega)}{d \omega}=\frac{2-2 \omega-\left(\gamma^{-}+\gamma^{+}+2 \gamma^{-} \gamma^{+}\right) \omega^{2}}{\left(1+\omega \gamma^{-}\right)^{2}\left(1+\omega \gamma^{+}\right)^{2}}
$$

which vanishes for

$$
\omega_{1}=\frac{-\sqrt{\left(2 \gamma^{+}+1\right)\left(2 \gamma^{-}+1\right)}-1}{\gamma^{+}+\gamma^{-}+2 \gamma^{-} \gamma^{+}}<0
$$

and

$$
\omega_{2}=\frac{\sqrt{\left(2 \gamma^{+}+1\right)\left(2 \gamma^{-}+1\right)}-1}{\gamma^{+}+\gamma^{-}+2 \gamma^{-} \gamma^{+}}
$$

Since

$$
\frac{d^{2} \psi_{1}\left(\omega_{2}\right)}{d \omega^{2}}=-\frac{2 \sqrt{\left(2 \gamma^{+}+1\right)\left(2 \gamma^{-}+1\right)}}{\left(1+\omega_{2} \gamma^{-}\right)^{2}\left(1+\omega_{2} \gamma^{+}\right)^{2}}<0
$$

and the relative minimum $\omega_{1}$ is outside of $(0,2), \omega_{2}$ is the maximum of $\psi_{1}(\omega)$ in $(0,2)$. Since

$$
\gamma^{+}+\gamma^{-}+2 \gamma^{+} \gamma^{-}=\frac{1}{2}\left(\left(2 \gamma^{+}+1\right)\left(2 \gamma^{-}+1\right)-1\right)
$$

it is easy to show that $\omega_{2} \leq 1$.
q.e.d.

This result is interesting, since it shows that the best $\omega$ in the context of (20) does not lead to overrelaxation, as is the case when SOR is used as an iterative solver, but to underrelaxation. In the special case $\gamma^{+}=\gamma^{-}=\gamma$, one obtains $\omega=\frac{1}{1+\gamma}<1$.

Note that Ruge and Stüben [7] suggest the norm

$$
\|A\|_{v}=\max _{1 \leq i \leq n}\left\{\frac{1}{v_{i}} \sum_{j=1}^{n} v_{j}\left|a_{i j}\right|\right\}
$$

where $v$ is a vector with positive elements $v_{i}$. This choice leads to

$$
\begin{aligned}
& \gamma^{-}=\max _{k}\left\{\frac{1}{v_{k} a_{k k}} \sum_{j<k} v_{k}\left|a_{k j}\right|\right\} \\
& \gamma^{+}=\max _{k}\left\{\frac{1}{v_{k} a_{k k}} \sum_{j>k} v_{k}\left|a_{k j}\right|\right\}
\end{aligned}
$$

which can be easily computed in practice.
Using a similar argument, one can show that the second smoothing property (12) is also fulfilled for SOR.

Lemma 5 Assume $A$ is symmetric positive definite, $Q=\frac{1}{\omega} D-L$, where $D$ is the diagonal and $-L$ the lower triangular part of $A$, and $0<\omega<2$. Assume also that

$$
\sigma_{2} \leq \frac{(2-\omega) \omega}{\left(|1-\omega|+\omega \gamma^{-}\right)\left(|1-\omega|+\omega \gamma^{+}\right)}
$$

where $\gamma^{+}$and $\gamma^{-}$are defined as in Lemma 4. Then

$$
\sigma_{2} e^{T}(Q-A)^{T} D^{-1}(Q-A) e \leq e^{T}\left(Q+Q^{T}-A\right) e
$$

For

$$
\begin{equation*}
\gamma^{-}>\frac{1}{2}, \quad \gamma^{+}>\frac{\gamma^{-}}{2 \gamma^{-}-1} \tag{21}
\end{equation*}
$$

the upper bound for $\sigma_{2}$ is maximal if

$$
\begin{equation*}
\omega=\frac{\sqrt{\left(2 \gamma^{+}-1\right)\left(2 \gamma^{-}-1\right)}-1}{2 \gamma^{+} \gamma^{-}-\gamma^{-}-\gamma^{+}} \leq 1 . \tag{22}
\end{equation*}
$$

For all other cases the maximal upper bound for $\sigma_{2}$ is obtained by setting $\omega=1$.

The proof is similar to the proof of Lemma 4. The inequality for $\sigma_{2}$ is determined in a similar way as was done for $\sigma_{1}$ in Lemma 4. In order to determine the best $\omega$, one needs to examine the continuous function

$$
\psi_{2}(\omega)= \begin{cases}\frac{(2-\omega) \omega}{\left(1+\omega\left(\gamma^{-}-1\right)\right)\left(1+\omega\left(\gamma^{+}-1\right)\right)} & \text { for } \quad 0<\omega \leq 1 \\ \frac{(2-\omega) \omega}{\left(\omega\left(\gamma^{-}+1\right)-1\right)\left(\omega\left(\gamma^{+}+1\right)-1\right)} & \text { for } \quad 1 \leq \omega<2\end{cases}
$$

If $2 \gamma^{+} \gamma^{-} \leq \gamma^{+}+\gamma^{-}$, it turns out that this function is increasing in $(0,1)$ and decreasing in $(1,2)$, which shows that the optimal $\omega$ is 1 . Only in the special case $2 \gamma^{+} \gamma^{-}>\gamma^{+}+\gamma^{-}$, which is equivalent to (21), is the absolute maximum in $(0,1)$ as given by (22) and can be obtained by straightforward differentiation of $\psi_{2}(\omega)$.

In this case, it appears that the use of a relaxation parameter is, in general, not beneficial. Only if condition (21) is fulfilled, which implies a matrix that is not diagonally dominant, should underrelaxation lead to better convergence.

In the following theorem, the SOR hybrid method is investigated.
Theorem 6 Assume that $A$ is symmetric positive definite, $D_{B}$ the block diagonal matrix with diagonal blocks $A_{k k}, Q=\frac{1}{\omega_{J}} \tilde{Q}$ as defined in (4), $\tilde{Q}_{k}=\frac{1}{\omega_{k}} D_{k}-L_{k}$ with $0<\omega_{k}<2$ for $k=1, \ldots, p$ and $0<\omega_{J} \leq \frac{1}{\rho\left(D_{B}^{-1} A\right)}$. Assume also that

$$
\sigma_{1} \leq \omega_{J} \min _{1 \leq k \leq p} \frac{\left(2-\omega_{k}\right) \omega_{k}}{\left(1+\omega_{k} \gamma_{k}^{+}\right)\left(1+\omega_{k} \gamma_{k}^{-}\right)}
$$

with

$$
\begin{aligned}
\gamma_{k}^{+} & =\left\|D_{k}^{-1} L_{k}^{T}\right\| \\
\gamma_{k}^{-} & =\left\|D_{k}^{-1} L_{k}\right\|
\end{aligned}
$$

where $\|$.$\| denotes any matrix norm, induced by a vector norm. Then the smoothing$ property

$$
\sigma_{1} e^{T} Q^{T} D^{-1} Q e \leq e^{T}\left(Q+Q^{T}-A\right) e
$$

is fulfilled.

## Proof:

The proof requires the application of Lemma 4 to the diagonal blocks of $A$, which leads to

$$
e_{k}^{T}\left(\tilde{Q}_{k}^{T}+\tilde{Q}_{k}-A_{k k}\right) e_{k} \geq \frac{\left(2-\omega_{k}\right) \omega_{i}}{\left(1+\omega_{k} \gamma_{k}^{-}\right)\left(1+\omega_{k} \gamma_{k}^{+}\right)} e_{k}^{T} \tilde{Q}_{k}^{T} D_{k}^{-1} \tilde{Q}_{k} e_{k}
$$

Now

$$
e^{T} A e \leq \rho\left(D_{B}^{-1} A\right) e^{T} D_{B} e \leq \frac{1}{\omega_{J}} e^{T} D_{B} e
$$

using the assumption $\omega_{J} \leq \frac{1}{\rho\left(D_{B}^{-1} A\right)}$. We thus obtain

$$
\begin{aligned}
e^{T}\left(Q^{T}+Q-A\right) e & =\frac{1}{\omega_{J}} \sum_{k=1}^{p} e_{i}^{T}\left(\frac{2}{\omega_{k}} D_{k}-L_{k}-L_{k}^{T}\right) e_{k}-e^{T} A e \\
& =\frac{1}{\omega_{J}} \sum_{k=1}^{p}\left(\frac{2}{\omega_{k}}-1\right) e_{k}^{T} D_{k} e_{k}+\frac{1}{\omega_{J}} e^{T} D_{B} e-e^{T} A e
\end{aligned}
$$

$$
\begin{aligned}
& \geq \frac{1}{\omega_{J}} \sum_{i=1}^{p} e_{k}^{T}\left(\tilde{Q}_{k}^{T}+\tilde{Q}_{k}-A_{k k}\right) e_{k}+\left(\frac{1}{\omega_{J}}-\rho\left(D_{B}^{-1} A\right)\right) e^{T} D_{B} e \\
& \geq \frac{1}{\omega_{J}} \sum_{k=1}^{p} \frac{\left(2-\omega_{k}\right) \omega_{k}}{\left(1+\omega_{k} \gamma_{k}^{-}\right)\left(1+\omega_{k} \gamma_{k}^{+}\right)} e_{i}^{T} \tilde{Q}_{k}^{T} D_{k}^{-1} \tilde{Q}_{k} e_{k} \\
& \geq \omega_{J} \min _{1 \leq k \leq p} \frac{\left(2-\omega_{k}\right) \omega_{k}}{\left(1+\omega_{k} \gamma_{k}^{-}\right)\left(1+\omega_{k} \gamma_{k}^{+}\right)} e^{T} Q^{T} D^{-1} Q e \\
& \geq \sigma_{1} e^{T} Q^{T} D^{-1} Q e
\end{aligned}
$$

This shows that a good choice of inner relaxation parameters should improve the hybrid SOR method, but even more crucial is a good choice of $\omega_{J}$ for (11) to be fulfilled. The best choice under the assumptions of this theorem is $\omega_{J}=1 / \rho\left(D_{B}^{-1} A\right)$. It is possible to obtain a good estimate of $\rho\left(D_{B}^{-1} A\right)$ using the procedure described in the next section. However, this approach would be very expensive and is therefore not practical, since it requires solving a linear system on each processor in each CG iteration step.

## 6 Practical Determination of the Outer Relaxation Parameter

The result on the optimal outer relaxation parameter obtained in Section 4 is only useful if it can be applied in practice. It is very important to get good estimates for $\rho\left(\tilde{Q}^{-1} A\right)$. This can be achieved by applying $k$ steps of preconditioned conjugate gradient to $A x=b$ with the preconditioner $\tilde{Q}$. Note that $\tilde{Q}$ needs to be symmetric positive definite. The preconditioning step is here just the application of one sweep of the smoother, which is fairly inexpensive. Due to the relationship of the conjugate gradient method and the Lanczos method [5], one can derive the tridiagonal Lanczos matrix $T_{k}$ from the parameters obtained within CG as can be found e.g. in [8]. Since the eigenvalues of $T_{k}$ approach the eigenvalues of $\tilde{Q}^{-1} A$ with increasing $k$, one can estimate the eigenvalues of $T_{k}$. This can be done using the Gershgorin estimate or, since $T_{k}$ is very small, using an eigenvalue solver for tridiagonal systems (such as the QR algorithm or bisection [5]). It is possible to get good estimates with this procedure with a fairly small number of CG iterations, e.g. $k=10$ or $k=15$. The use of this procedure increases the setup time of AMG. However, for problems that require a good smoothing parameter, the resulting decrease in number of iterations and solve time far outweighs this increase in setup time, as can be seen in the next section. Note that in some cases conjugate gradient has found to be unstable and it might be better to use a stable implementation of the preconditioned Lanczos algorithm. We used conjugate gradient, since it was immediately available to us, and we observed no instabilities in our experiments.

## 7 Numerical Results

The methods described in the previous sections are applied to various very large 3 -dimensional elasticity problems composed of 3 concentric spherical shells. An octant of this domain is shown in Figure 1. The outer shells are composed of steel, the inner shell is composed of lucite. We consider problems without and with slide surface boundary conditions. In the case of slide surface boundary conditions the steel and lucite spheres are allowed to slide tangentially relative to each other. Adding the slide surface boundary conditions leads to an indefinite problem. It is,


Figure 1: Finite element discretization of a sphere using quadrilateral elements.
however, possible to reduce the system to a positive definite system through the elimination of a subset of equations [4]. In our experiments we use the reduced system. The problems were run on the ASCI White Computer at LLNL. Two different problem sizes are considered: the smaller problem is 497,664 elements and is run on 32 processors; the larger problem consists of almost 4 million elements and is run using 256 processors. Since the considered problem has multiple degrees of freedom per grid point, we use AMG for problems derived from systems of partial differential equations, employing the function, or "unknown", approach [7]. This approach coarsens each physical variable separately and interpolates only within variables of the same type. The smoothers are symmetric Gauß-Seidel with $\omega_{J}=1$ and the $\omega_{J}$ as given in (17), which is obtained using at most 10 CG-iterations. We are able to use a nodal hybrid Gauß-Seidel, i.e. a block Gauß-Seidel method with $3 \times 3$ blocks, due to the structure of the problem. Unfortunately, the nodal structure is destroyed after the first level, so use of a nodal smoother beyond the finest level does not make sense. $\omega_{J}$ was here estimated with 15 CG-iterations, since 10 CG-iterations turned out to be not good enough. We also consider hybrid SSOR. Since we have no procedure to determine the best inner relaxation parameters, we present results for uniform $\omega=0.75,0.5$ and 0.25 for the moderate size problem and use $\omega=0.5$ for the large problem. In all these experiments, AMG is used as a preconditioner for CG. Therefore, in order to not destroy the symmetry of the problem, only symmetric smoothers are used.

Table 1 gives the estimates of the outer relaxation weights that have been used for the larger elasticity problems. The fact that the relaxation parameters on the finer levels are smaller than 0.5 , and thus $\rho\left(\tilde{Q}^{-1} A\right)>2$ indicates that hybrid (block) Gauß-Seidel is not a convergent iterative scheme for these problems. Table 2 contains the notations used for Tables 3 through 6. $\tilde{Q}$ denotes the symmetric GaußSeidel matrix, and $\tilde{Q}_{B}$ denotes the nodal or $3 \times 3$ block Gauß-Seidel matrix.

For moderate size problems, Table 3 shows a fairly small improvement of scaled smoothers over unscaled smoothers. However, for large problems, Table 4 shows

| without slide surfaces |  |  |  | with slide surfaces |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| level | $\omega_{J}$ | level | $\omega_{J}$ | level | $\omega_{J}$ | level | $\omega_{J}$ |
| 1 | 0.425 | 12 | 0.588 | 1 | 0.465 | 12 | 0.588 |
| 2 | 0.458 | 13 | 0.585 | 2 | 0.462 | 13 | 0.584 |
| 3 | 0.480 | 14 | 0.591 | 3 | 0.453 | 14 | 0.605 |
| 4 | 0.460 | 15 | 0.614 | 4 | 0.453 | 15 | 0.612 |
| 5 | 0.498 | 16 | 0.604 | 5 | 0.502 | 16 | 0.627 |
| 6 | 0.555 | 17 | 0.603 | 6 | 0.556 | 17 | 0.639 |
| 7 | 0.585 | 18 | 0.601 | 7 | 0.580 | 18 | 0.613 |
| 8 | 0.567 | 19 | 0.587 | 8 | 0.588 | 19 | 0.620 |
| 9 | 0.585 |  |  | 9 | 0.583 |  |  |
| 10 | 0.570 |  | al: | 10 | 0.585 |  | al: |
| 11 | 0.572 | 1 | 0.397 | 11 | 0.575 | 1 | 0.419 |

Table 1: Relaxation parameters for an elasticity problem, 3,981,312 elements, 256 procs
a significant improvement for scaled smoothers. Overall, the best time achieved is about 10 times as fast as the original test (which uses hybrid Gauß-Seidel without any relaxation parameter). Interestingly enough, it turns out that for this problem, the use of the nodal smoother does not improve convergence; apparently point smoothers are sufficient to smooth the error.

Table 5 shows this changes when we include slide surfaces. On the moderate size problem, using hybrid SGS without any smoothing parameter converges in 390 iterations, whereas scaled SGS converges about 3 times as fast. It is interesting that just applying the nodal smoother without any scaling parameter entails a similar number of iterations, showing that for this problem a nodal smoother is more suitable. Scaling the nodal smoother leads to a further improvement of another factor of about 2.3 . Scaling the point smoother improves this result only slightly. Reducing the number of sweeps increases the number of iterations, but decreases the time per iteration. The best result is about 7 times faster than the original solver.

The larger problem with slide surfaces diverges without any relaxation parameters, even when a nodal smoother is used on the finest level (Table 6). Scaled SGS converges within 192 iterations, but convergence is twice as fast when a scaled nodal smoother is used on the finest level. The overall fastest (with regard to time) combination solves this very large problem, which diverges when unscaled smoothers are employed, in about 5 minutes.

The results in Tables 3 through 6 show that a good inner relaxation parameter for hybrid SSOR is 0.5 . Overall, underrelaxation, i.e. choosing an inner relaxation parameter smaller than 1, beats hybrid SGS, which is equivalent to hybrid SSOR $\omega=1$. Overrelaxation $(\omega>1)$, which is not presented here, leads to a further decrease in performance.

## 8 Conclusions

The use of relaxation parameters for hybrid smoothers is analyzed. Both outer as well as inner parameters are considered. Analysis of the inner SOR relaxation parameter shows that, in most cases, underrelaxation (i.e. $\omega<1$ ) is preferred to overrelaxation. An outer relaxation parameter for symmetric positive definite matrices and splittings is determined that is optimal under the assumptions of the theory, and an automatic procedure to determine it is implemented. Numerical

| Smoother | Description |
| :--- | :--- |
| G | hybrid symmetric Gauß-Seidel (SGS) |
| $\mathrm{S}, \omega$ | hybrid SSOR, $\omega_{k}=\omega$ for $k=1, \ldots, p$ |
| SG | scaled hybrid SGS, $\omega_{J}=\frac{1}{\tilde{\rho}\left(\tilde{Q}^{-1} A\right)}$ |
| N | hybrid nodal $(3 x 3$ blocks) SGS |
| SN | scaled hybrid nodal SGS, $\omega_{J}=\frac{1}{\tilde{\rho}\left(\tilde{Q}_{B}^{-1} A\right)}$ <br> $<$ S1 $>/<$ S2 $>$ |
| $<$ S1 $>$ used only on finest level, <br> $<$ S2 $>$ used on coarser levels |  |

Table 2: Smoother notations

| Smoother | no. of <br> sweeps | no. of <br> its | setup <br> time | solve <br> time | total <br> time |
| :---: | :---: | :---: | :---: | :---: | :---: |
| G | 2 | 74 | 15 | 266 | 281 |
| $\mathrm{~S}, 0.75$ | 2 | 63 | 15 | 233 | 248 |
| S, 0.5 | 2 | 45 | 15 | 164 | 179 |
| S, 0.25 | 2 | 54 | 15 | 197 | 212 |
| SG | 2 | 42 | 25 | 157 | 182 |
| N/G | 2 | 59 | 16 | 223 | 240 |
| SN/G | 2 | 39 | 22 | 148 | 170 |
| SN/G | 1 | 47 | 22 | 108 | 130 |
| SN/SG | 2 | 41 | 28 | 161 | 189 |
| SN/SG | 1 | 50 | 28 | 119 | 148 |

Table 3: Elasticity problem without slide surfaces, 497,664 elements, 32 procs, $\mathrm{n}=$ 1,545,483

| Smoother | no. of <br> sweeps | no. of <br> its | setup <br> time | solve <br> time | total <br> time |
| :---: | :---: | :---: | :---: | :---: | :---: |
| G | 2 | 484 | 25 | 2097 | 2122 |
| S, 0.5 | 2 | 67 | 25 | 301 | 326 |
| SG | 2 | 61 | 38 | 270 | 308 |
| N/G | 2 | 562 | 25 | 2477 | 2502 |
| SN/G | 2 | 57 | 32 | 258 | 290 |
| SN/G | 1 | 67 | 32 | 184 | 216 |
| SN/SG | 2 | 60 | 41 | 276 | 317 |
| SN/SG | 1 | 71 | 41 | 196 | 237 |

Table 4: Elasticity problem without slide surfaces, $3,981,312$ elements, 256 procs, $n$ $=12,152,595$

| Smoother | no. of <br> sweeps | no. of <br> its | setup <br> time | solve <br> time | total <br> time |
| :---: | :---: | :---: | :---: | :---: | :---: |
| G | 2 | 390 | 15 | 1427 | 1443 |
| S, 0.75 | 2 | 201 | 15 | 750 | 765 |
| S, 0.5 | 2 | 149 | 15 | 555 | 570 |
| S, 0.25 | 2 | 185 | 15 | 696 | 711 |
| SG | 2 | 135 | 26 | 517 | 543 |
| N/G | 2 | 142 | 17 | 547 | 564 |
| SN/G | 2 | 63 | 23 | 245 | 268 |
| SN/SG | 2 | 60 | 31 | 237 | 267 |
| SN/SG | 1 | 74 | 31 | 181 | 212 |

Table 5: Elasticity problem with slide surfaces, 497,664 elements, 32 procs, $\mathrm{n}=$ 1,587,825

| Smoother | no. of <br> sweeps | no. of <br> its | setup <br> time | solve <br> time | total <br> time |
| :---: | :---: | :---: | :---: | :---: | :---: |
| G | 2 | fail | 23 | - | - |
| S, 0.5 | 2 | 342 | 23 | 1375 | 1399 |
| SG | 2 | 192 | 37 | 767 | 804 |
| N/G | 2 | fail | 27 | - | - |
| SN/G | 2 | 146 | 29 | 600 | 630 |
| SN/SG | 2 | 86 | 41 | 365 | 406 |
| SN/SG | 1 | 106 | 41 | 272 | 313 |

Table 6: Elasticity problem with slide surfaces, $3,981,312$ elements, 256 procs, $\mathrm{n}=$ 12,320,217
experiments show that for certain elasticity problems significant improvements can be achieved using relaxation parameters.

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