

KILLING THE ELECTRON CLOUD EFFECT IN THE LHC ARCS*

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Abstract

A getter/electrode assembly has been devised to suppress the regeneration mechanism of the electron cloud effect in the arc dipoles of LHC. The assembly consists of a copper foil electrode, supported through an insulating layer on a stainless steel skid, which would rest upon the flat bottom of the beam screen. The electrode is coated with NEG to provide effective pumping of all non-inert gases from the vacuum. When the electrode is biased $\sim +100$ V, both electrons from beam ionizations and secondary electrons produced on the beam screen surface would be cleared, killing the regeneration mechanism. The NEG surface can be regenerated by passing a current through the electrode to heat it to ~ 240 C. The heat transfer (radiant + conductive) to the beam screen during regeneration is estimated ~ 10 W/m, within limits to maintain the beam screen at nominal 20 K temperature during regeneration.

ELECTRON CLOUD EFFECT

The electron cloud effect (ECE) arises from the multipactoring of electrons within the vacuum chamber. In the Large Hadron Collider (LHC) electrons are produced when gas molecules are ionized by protons in a high-intensity bunch, and also when synchrotron light strikes the side walls. The critical energy of this light is 44 eV, for which photoemission yield is maximum for many materials.

Such electrons typically have \sim eV energies and so are still inside the beam screen when the next proton bunch passes. The Coulomb field on that bunch is sufficient to accelerate ambient electrons to energies of $\sim 100 - 1000$ eV. After the bunch passes each such energetic electron travels to the side wall before arrival of the next bunch, strikes the wall, may emit one or more secondary electrons. If the secondary electron yield $\delta > 1$, the process is regenerative and the ambient electron density will grow exponentially. For copper surfaces such as the beam screen, δ typically has values ranging from a low of ~ 1.1 to a high of ~ 1.7 . The value of this parameter, notoriously difficult to control, is thus critical to the multipactoring process of the electron cloud effect.

ECE has several bad effects on collider operation. First, the electron distribution responds dynamically to the proton bunch as it passes and can excite transverse mode-coupling instability (TMCI), coupled-bunch instabilities, head-tail motion within the proton bunch, tune spread, beam loss and incoherent emittance growth. There is an extensive literature on the effect of electron cloud effect on beam dynamics [1]. Zimmermann and Benedetto give

a recent summary of the understanding [2]. Simulations using the HEAD-TAIL code indicate the possibility of long-term emittance growth that could be detrimental with a storage time of hours. Also transverse mode-coupling instabilities will be excited by the electron cloud. While it has been possible to control the fast instabilities in the SPS, control has required a degree of chromaticity that could be problematic for LHC. Presently there is no cure for long-term incoherent emittance growth due to the electron cloud, if it occurs, other than reducing the electron density.

Second, the electron cloud desorbs gases from the walls of the beam screen, and could push the limit for beam lifetime and pressure-bump instabilities. It has been proposed that the surfaces of the beam screen could be conditioned. Conditioning appeared to improve the pressure in the beam screen in the COLDEX experiment [3], but the heat load attributed to electron cloud did not significantly reduce with time. Furthermore cryopumping onto the cold bore relies upon there being a limited coverage of helium on the cold bore surface. If there are any cold leaks in the arcs the helium coverage reach levels that would inhibit cryopumping of the cold bore.

Lastly the energetic electrons heat the surfaces that they impact. This heat has been studied both by models and by experiment [4], and is expected to be $\sim 1-2$ W/m for LHC bunch intensity, depending upon the value of secondary electron yield for the copper surfaces of the beam screen. The heat load from ECE with bunches of 10^{11} protons could approach installed refrigeration capacity for 25 ns bunch spacing (the mode most favorable for high-luminosity physics), and might limit operation to the 75 ns bunch spacing foreseen for initial operation. A shield has been installed behind the slots that couple the beam

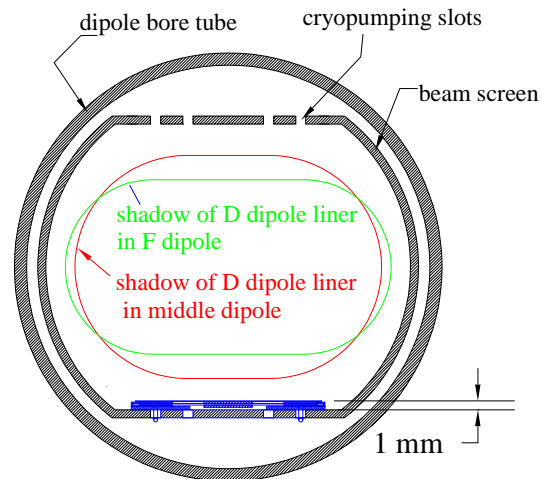


Figure 1. ECE electrode assembly (blue) on base of beam screen. Shadows of the beam screen from the D dipole are shown at the F dipole (green) and middle dipole (red).

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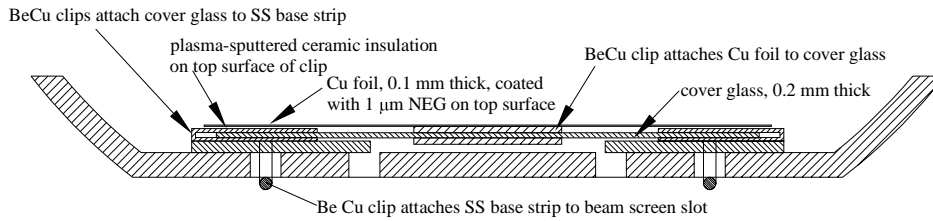


Figure 2. Detail of ECE electrode assembly.

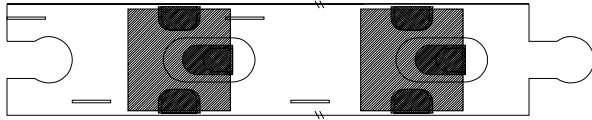


Figure 3. Plan view of a segment of the ECE electrode, showing attachment of Cu electrode to SS skid.

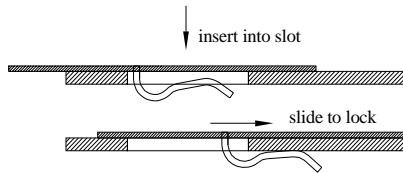


Figure 4. Detail showing attachment of SS skid through slots in beam screen.

screen vacuum to the cold walls in order to reduce the transport of these energetic electrons onto the cold bore. Unfortunately it also reduces cryopumping by a factor 3.

THE ECE ELECTRODE

The ECE electrode is a getter/electrode structure designed to integrate into the beam screen as shown in Figure 3. It contains a copper electrode, mounted on a 316LN stainless steel skid which is in turn anchored to the floor of the beam screen. The electrode is electrically and thermally isolated to operate at +100 V with respect to the beam screen.

The electrode is coated with Ti/Zr/V getter following the process developed by Benvenuti *et al.* [5]. The coating of a $\sim 2 \mu\text{m}$ thick NEG layer on the electrode will be performed in the magnetron sputtering facility of Chigiato at CERN. The NEG layer will provide distributed high-throughput pumping to assist the cryopumping to the bore tube in maintaining excellent vacuum in the arcs.

The electrode is insulated from the skid by means of a thin glass plate (a standard 22 mm square microscope cover glass). The borosilicate glass has large but finite resistivity, so that any charge developed on the insulator during operation will be cleared. The copper electrode is attached to the glass by a BeCu tab clip; the glass is attached to the skid by two side tab clips as shown in Figure 3.

It is necessary to provide sufficient thermal isolation to limit heat transfer to the beam screen so that the electrode

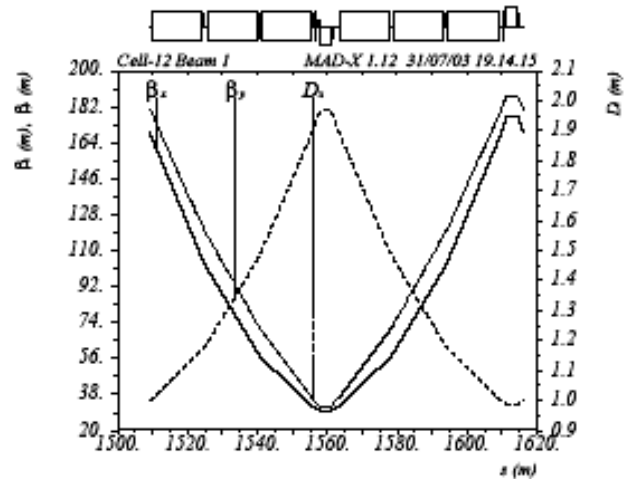


Figure 5. Arc lattice, showing b_v in the three dipoles of a half-cell.

can be heated to 240 C to activate the NEG layer without increasing the beam screen temperature. The arrangement shown does this by providing a $\sim 1 \text{ cm}$ long conduction path through the 0.2 mm thick glass from the skid clip to the copper clip.

The copper foil electrode is chosen to be $100 \mu\text{m}$ thick, sufficient to provide a low impedance for image currents of the proton beam but thin enough to provide a favorable characteristic for heating the strip during NEG regeneration. The top surface of the side tab clips is coated with a plasma-sprayed Al_2O_3 insulating layer to provide reliable electrical isolation from ground.

The skid will be fabricated from six 2.5 m segments of straight 316LN material. The ends of each segment will be die-cut to form a key-in-lock joint to the next segment as shown in Figure 3. The curvature required for the 9 mm magnet sagitta of the dipole will be accommodated in the key-in-lock joints. The joints will be spot-welded to lock them together. The foil electrode is sufficiently flexible that it can be fabricated straight and mounted to the slightly curved skid. During cooldown the copper electrode will shrink by a strain $\delta L/L = 3 \times 10^{-4}$ relative to the SS skid. The differential shrinkage will be accommodated by bowing the electrode slightly as it is attached to the successive insulators. The additional length will be controlled by rod spacers placed between the skid and the foil electrode midway between the locations of successive mounting tabs. Once the tabs are all attached to the insulators the spacers will be removed.

APERTURE CONSIDERATIONS

The electrode assembly occupies 1 mm of vertical aperture, as shown in Figure 1. In the three dipoles of a given half-cell of the arc lattice, the vertical betatron function β_v has different maxima in each of the three dipoles, as shown in Figure 5. The dipole nearest the horizontally focusing quadrupole presents the most constraining vertical aperture. If one maps that constraining aperture with β_h , β_v , and dispersion D, the ‘shadow’ of the D dipole at the other two dipoles are shown in Figure 1. It is thus seen that if the electrode assembly were installed in only those two of the three dipoles in every half-cell, there would be *no reduction in beam aperture at all*. If the electrode assembly were installed in the third dipole, vertical aperture would be reduced by 1 mm. The various considerations of beam optics, placement errors, etc. will determine whether this loss of aperture is acceptable.

BEHAVIOR DURING MAGNET QUENCH

During magnet quench, the current in a dipole decreases from full excitation to zero in ~ 0.1 s. The time-changing flux through the beam screen region induces currents in the copper lining of the screen and also in the electrode proposed here. Rathgen [6] showed that while induced currents primarily cause tension within a conducting structure within a \sim uniform-field dipole, the small multipole fields produce some net forces as well. In order to evaluate this possibility for the ECE electrode assembly, we simulated the quench in the body of an LHC dipole, including the allowed field multipoles, using the 2-D electromagnetic code PE2D [7].

The outcome of these calculations is as follows:

- The electrode will dissipate ~ 550 J/m, heating it to ~ 160 K.
- The electrode will experience ~ 220 MPa of balanced tension, which is about half the yield strength of the cryogenic copper.
- The maximum instantaneous unbalanced force acting on the electrode consists of
 - 20 N/m horizontal (due to quadrupole);
 - 170 N/m vertical (due to sextupole).
- That force loading is well within the strength of the BeCu clips attaching the electrode, the insulator, and the skid.
- The vertical force would deflect the two outer edges of the electrode upwards. The deflection has been modeled using ALGOR [8]; the highest deflection is .03 mm.

ASSEMBLY IN DIPOLE

Because the ECE electrode has only recently been devised, it will be necessary to install the assembly into each beam screen after it is already installed in a dipole. For this purpose a trolley arrangement has been devised: a string of 5-wheel cars as shown in Figure 6. First the complete electrode assembly is inserted in the dipole on a carrier and the carrier is removed. The lead trolley car is

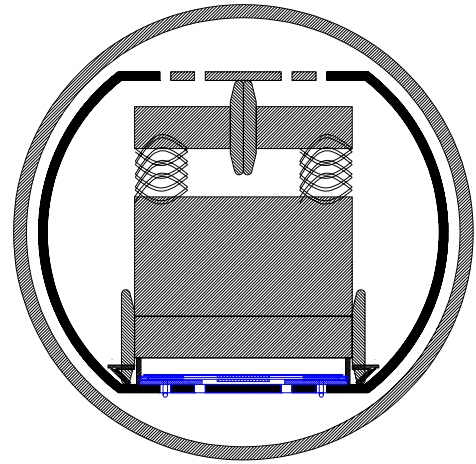


Figure 6. Installation trolley.

then positioned so that it sits on the near end of the assembly and locates itself in the beam screen from the bottom two corners and a spring-loaded idler against the top. As it is pulled through along the length of the dipole, its wheel assembly maintains local alignment and a set of spring fingers aligns the electrode assembly and presses it down so that the spring clips are positioned in the beam screen slots (Figure 4a). As each car proceeds, another is attached behind it so that finally the entire electrode assembly is in the proper position. Then the assembly is pulled towards the lead end to lock the clips to the screen (Figure 4b).

CONCLUSIONS

An electrode/getter assembly has been designed that is capable of being installed in the LHC arc dipoles in vertical aperture of 1 mm along the bottom of the beam screen. The electrode should clear all electrons and eliminate the electron cloud effect in the arcs. As a collateral benefit, the current to each electrode will give a direct measurement of the beam vacuum in each dipole, which information is not otherwise available and should be very useful during commissioning and operation of the collider.

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