APPENDIX E

TARGET AND REFERENCE WAVE PACKETS

Here we derive analytic expressions for the target (reference) wave packet defined by Eq. (5.25) (Eq. (5.26)). Using the short-pulse limit of the B_x pulse propagator (5.11) and the free-evolution operator of first-order in J given in footnote 52 of Chapter V, we find

$$\left| \begin{array}{c} 1 \\ 1 \end{array} \right\rangle = \frac{J}{2} area_{B_{0}}^{t_{w}} d e^{iH_{1}(t_{w})} e^{iH_{1}} \left| 0,0 \right\rangle.$$
 (E.1)

The target wave packet (E.1) is assumed to originate from the vibrational ground-state of the electronically unexcited complex; $area_I = dt A_I(t)$ is the integrated pulse envelope.

We can use harmonic-oscillator creation and annihilation operators to go further with Eq. (E.1). Adopting the usual definitions $q_a = (2m^-)^{-1/2}(a^\dagger + a)$, $p_a = i(m^-/2)^{1/2}(a^\dagger - a)$, $q_b = (2m^-)^{-1/2}(b^\dagger + b)$, $p_b = i(m^-/2)^{1/2}(b^\dagger - b)$, and introducing the corresponding phase-space translation operators $T_a(-) = \exp(-a^\dagger - a^*)$ and $T_b(-) = \exp(-b^\dagger - a^*)$, we have

$$H_1 = T_a(\)(H_0 + \ _1)T_a^{\dagger}(\)$$
 (E.2)

$$H_1 = T_b(\)(H_0 + \ _1)T_b^{\dagger}(\)$$
 (E.3)

$$H_2 = T_a(\)T_b(\)(H_0 + \ _2)T_b^{\dagger}(\)T_a^{\dagger}(\) .$$
 (E.4)

From the relations (E.1) through (E.4) and some operator algebra it follows that

$$\left| \begin{array}{c} 1 \end{array} \right\rangle = \frac{J}{2} area_B e^{-2^{-2} i_{1-vib}/2} \tag{E.5}$$

$$\int_{0}^{vib/2} d \exp\{i(a_{1} a_{1}) + 2i^{2}\sin a + (1 e^{i} a)a^{\dagger} + (1 + e^{i} a)b^{\dagger}\} |0,0\rangle.$$

The dimensionless displacement is $=(m/2)^{1/2}d$, and the waiting time has been set to $t_w = \frac{1}{v_{ib}}/2$.

In the short-pulse limit, the reference wave packet (5.26) prepared from the vibronic ground state remains a quasi-classical coherent state of the form

$$\left| \begin{array}{ccc} \\ \\ \\ \end{array} \right\rangle = i \frac{3}{2} area_{A} area_{C} area_{D} e^{i p^{t_{p}} i d^{t_{d}}} e^{iH_{2}t_{d}} e^{iH_{1}(t_{p}+t_{w}+t_{d})} e^{iH_{0}t_{p}} \left| 0,0 \right\rangle. \tag{E.6}$$

Relations (E.3) and (E.4) plus some operator manipulations (and the choice $t_w = v_{ib}/2$) lead to the more explicit expression

$$\left| \begin{array}{c} 1 \\ 1 \end{array} \right\rangle = i \frac{3}{2} area_A area_C area_D$$
 (E.7)

exp
$$i \frac{1 - vib}{2} + i \binom{p}{p-1} t_p = i \binom{d+1}{d+1-2} t_d$$

 $+ \frac{2}{3} e^{i t_d} - \frac{2}{3} e^{-i t_p} - 2^{-2} + (1 - e^{i t_d}) a^{\dagger} + (1 + e^{-i t_p}) b^{\dagger} \Big\} |0,0\rangle.$

The inner product of (E.5) and (E.7) leads to the experimentally isolable complex overlap, Eq. (5.29), that is plotted in Chapter V.