

# Coordinated monitoring of the eccentric O-star binary Iota Orionis: optical spectroscopy and photometry

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## ABSTRACT

With the objective of investigating the wind–wind collision phenomenon and supporting contemporaneous X-ray observations, we have organized a large-scale, coordinated optical monitoring campaign of the massive, highly eccentric O9 III+B1 III binary Iota Orionis. Successfully separating the spectra of the components, we refine the orbital elements and confirm the rapid apsidal motion in the system. We also see strong interaction between the components during periastron passage and detect phase-locked variability in the spectrum of the secondary star. However, we find no unambiguous signs of the bow shock crashing on the surface of the secondary, despite the predictions of hydrodynamic simulations. Combining all available photometric data, we find rapid, phase-locked variations and model them numerically, thus restricting the orbital inclination to  $50^\circ \lesssim i \lesssim 70^\circ$ .

**Key words:** binaries: general – stars: early-type – stars: general – stars: individual: Iota Orionis.

## 1 INTRODUCTION

A significant fraction of massive stars can be found in binary systems, with high probability that both stars are nearly equally massive. Many of these early-type binaries have orbits with high eccentricity, especially those systems with longer periods. Systematic study of early-type, highly eccentric binaries provides insight into a variety of phenomena, including various aspects of the physics of colliding winds, tidal interactions and tidally induced pulsations.

The study of colliding wind phenomena has greatly advanced in recent years. Following the pioneering study of Shore & Brown (1988), there have been major advances in several directions, both observational and theoretical. In the X-ray range, among the best-studied massive colliding wind WR+O systems are HD 193793 (Williams et al. 1990; Pollock et al. 1995) and  $\gamma^2$  Velorum (Willis, Schild & Stevens 1995; Stevens et al. 1996). In the case of the strongly interacting O+O binaries, Corcoran (1996) presented a collection of X-ray observations, including the system that is the focus of this study,  $\iota$  Ori. Another system worth mentioning is Eta Carinae, now also believed to be a colliding wind binary with high eccentricity (Corcoran et al. 1998; Pittard et al. 1998).

Practically all recent optical/UV studies of colliding wind

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WR+O binaries are mentioned by Stevens & Howarth (1999). Examples of O+O star systems include 15 Mon (Gies et al. 1997b); 29 UW Canis Majoris (Wiggs & Gies 1993; Bagnuolo et al. 1994); Plaskett's star (Wiggs & Gies 1992); and V729 Cyg (Rauw, Vreux & Bohannan 1999).

A wide variety of phenomena is predicted to occur in a massive binary system. They generally stem from a presumably strong interaction between the components, including enhanced mass loss, mass transfer between the components, various wind–wind collision effects, and stellar oscillations. Some of these phenomena lead to eruptive events and others gradually modify the orbital parameters. For longer periods, the system is more likely to have a large orbital eccentricity, accentuating many of the effects and making new ones appear. The interaction effects can be identified through the use of high-resolution spectroscopy by searching for peculiar absorption or emission line profiles and their phase-locked variability. However, it is very likely that several interaction effects operate simultaneously in a given system, thus making it rather difficult to isolate the physical agent responsible for any particular component of variability. The problem is further complicated in relatively close binaries with high mass-loss rates. Hence, it is also desirable to undertake a study of binaries where the mass-loss rates are not so extreme, and where there is a good chance of isolating the effects arising from the tidal interactions from the effects of the interacting winds.

To this end we have undertaken a multiwavelength campaign on the O-star binary Iota Orionis (HR 1889; HD 37043: O9 III+ B1 III), which is a highly eccentric ( $e = 0.764$ ), relatively long-period ( $P = 29.133\,76$  d) binary. General properties of the system were reviewed by Stevens (1988) and by Stickland et al. (1987, hereafter SPLH87), who studied the system in order to determine whether the mass loss from either of the components is enhanced at periastron with respect to other orbital phases, finding no clear evidence for this effect. Gies et al. (1996) studied the system to determine whether non-radial oscillations driven by tidal forces at periastron are present, failing to detect the expected line-profile variability. Moreno & Koenigsberger (1999) developed a two-dimensional tidal interaction model including stellar rotation. The model was applied to  $\iota$  Ori by following the vertical displacements and velocity amplitudes of each equatorial stellar surface element throughout the orbital cycle. It appears that with a radius for the primary star of  $15 R_{\odot}$ , the tidal oscillation amplitudes might not exceed the limit of detectability, which could explain the non-detection of oscillation by Gies et al. (1996). The model enables one to predict the behaviour of the stellar surface of both stars in the system. The growth and recession of the tidal bulges translate into photospheric line-profile and intensity variations, with the strongest variability in  $\iota$  Ori expected to occur within a time-span of  $\pm 0.75$  d around periastron passage (i.e.,  $\pm 0.02$  in phase).

$\iota$  Ori has been recently observed by the ASCA satellite (cf. Pittard et al. 2000) in order to investigate wind–wind collision phenomena. Here we present the results of the complementary optical campaign, organized in support of the X-ray observations.

## 2 OBSERVATIONS

### 2.1 Spectroscopy

New, high-quality spectroscopic observations of  $\iota$  Ori were obtained at several locations.

- (1) Three high signal-to-noise ratio ( $S/N \sim 250$ – $350$  in the

blue,  $\sim 350$ – $450$  in the red) spectra were obtained on 1996 and 1997 September using the 2.15-m telescope of Complejo Astronómico El Leoncito (CASLEO), San Juan, Argentina. The Cassegrain REOSC spectrograph, with a  $400 \text{ line mm}^{-1}$  grating and  $1024 \times 1024$  CCD ( $24 \times 24 \mu\text{m}^2$  pixel size) provided spectral coverage of  $\lambda\lambda$  3670–6120 and 5830–8130 Å, and spectral resolution of  $\sim 0.5$ – $0.8$  Å (blue–red regions, based on the FWHM of the interstellar lines). The initial processing was performed using the data reduction package written by J. Chauville.

- (2) Six spectra of  $\iota$  Ori were collected around periastron passage on 1997 February using the Aurélie spectrograph fed by the 1.52-m telescope at Observatoire de Haute Provence (OHP), France. The spectrograph was equipped with a  $1200 \text{ line mm}^{-1}$  grating blazed at 5000 Å, providing a reciprocal dispersion of  $7.6 \text{ Å mm}^{-1}$  over a wavelength range from 6510 to 6710 Å. The detector was a Thomson TH7832 linear array with pixel size of  $13 \mu\text{m}$ . The spectral resolution as derived from the FWHM of the calibration lines was  $\sim 0.3$  Å. A few additional observations were gathered with the same equipment in 1998 November, but using a  $1800 \text{ line mm}^{-1}$  grating providing reciprocal dispersion of  $5 \text{ Å mm}^{-1}$ . All the OHP data were reduced in a standard way using the MIDAS software developed at the European Southern Observatory. To achieve a first-order correction of the telluric absorption lines between 6500 and 6600 Å, a template of the telluric spectrum was built from observations of HD 36512 ( $\nu$  Ori, B0 V) at very different airmasses. The mean  $S/N$  in the continuum after correction for the telluric absorptions reaches  $\sim 200$  in the spectral range 6510–6600 Å and  $\sim 500$  elsewhere.

- (3) Multiple observations were obtained in 1997 September–October with a fiber-fed échelle spectrograph attached to the 1-m telescope of the Ritter Observatory (USA) and equipped with a liquid nitrogen-cooled Wright Instruments Ltd. CCD camera, which incorporates an  $800 \times 1200$  EEV CCD05-20-0-202 thick sensor with pixel dimensions of  $22.5 \times 22.5 \mu\text{m}^2$ . The data reduction used Interactive Data Language (IDL) programs described in Gordon & Mulliss (1997). The spectral resolution element corresponds to an instrumental function of width about 4.2 pixel FWHM, i.e.  $0.21$ – $0.25$  Å. About 70 Å of each of nine échelle orders was covered; in addition to H $\alpha$ , features of interest in the spectra included O III  $\lambda 5592$ , C III  $\lambda 5696$  and He I  $\lambda 5876$ . Under highly variable weather conditions,  $S/N$  in the range 100–200 were achieved in exposure times ranging from 5 to 15 min. They were raised to  $S/N \sim 200$ – $350$  after triple-pixel rebinning of the original spectra. Experience with repeated high- $S/N$  observations of K-giant radial velocity standards indicates that the velocity stability of this instrumentation is better than  $1 \text{ km s}^{-1}$ . Removal of the telluric absorption lines was performed using the spectrum of  $\zeta$  Aql as a template.

- (4) Spectroscopic observations were carried out at the Observatorio Astronómico Nacional San Pedro Mártir (SPM), Mexico, on 1997 September 22 and 23, using an échelle spectrograph mounted on the 2.11-m telescope. The  $2048 \times 2048$  Tektronics CCD detector provided  $0.4$ – $0.9$  Å spectral resolution and  $\lambda\lambda$  3200–7350 Å spectral coverage. A total of 27 spectra were obtained, eight on September 22 and 19 on September 23, with exposure times ranging between 40 and 180 s. The spectra on each night were obtained within a time-span of less than 2 h. An average spectrum was constructed for each night. The  $S/N$  ratio for the average spectrum ranges from 1000 in the blue to 600–800 in the red. The spectra were processed with IRAF.

- (5) One spectrum ( $S/N \sim 300$ ) was obtained at Observatoire du Mont Mégantic (OMM), Québec, Canada, using the 1.6-m

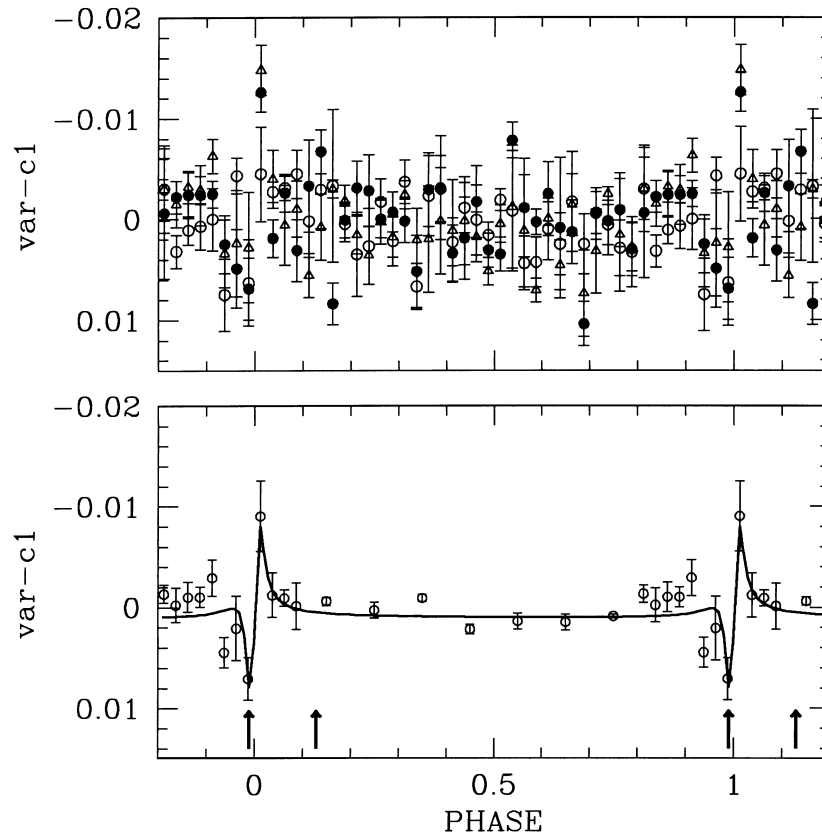
telescope with the attached Cassegrain spectrograph and Loral 2048 × 2048 CCD. The 600 line mm<sup>-1</sup> grating provided  $\lambda\lambda 5770$ –7120 Å spectral coverage and  $\Delta\lambda = 2.0$  Å (3 pixel) resolution.

## 2.2 Photometry

Iota Ori was monitored in Johnson *UBV* optical filters at one of the automatic 25-cm telescopes in Arizona (cf. Young et al. 1991 for a general description). Between 1996 December and 1998 April we obtained 724 *UBV* observations of Iota Ori (about 240 data points in each filter; 1–3 *UBV* cycles per night) with a typical accuracy for a single observation of 0.005–0.007 mag. Iota Ori was observed along with HR 1848 and 1840 as comparison and check stars, respectively, using a sequence of 10-s integrations: check (ch, *UBV*) – sky (s, *UBV*) – comparison (c, *UBV*) – Iota (I, *UBV*), where the sequence of *UBV* observations was ch-s-c-l-c-l-c-l-c-s-ch.

In the absence of apparent (with amplitude  $A \geq 0.01$  mag) long-term variations, we folded the individual *UBV* data with  $P_{\text{orb}} = 29.133\,76$  d and  $E_0 = \text{HJD } 245\,1121.658$  (SPLH87 and below), binned them to 0.1 phase bins and combined the binned *U*, *B* and *V* light curves by simple averaging after removal of the mean in each filter, as the differences between the binned *U*, *B* and *V* light curves did not exceed the observational errors. We immediately noticed flux variations taking place around periastron passage (phase 0.0). To provide better time resolution and improve the accuracy even further, we again phased the individual *U*, *B* and *V* data and binned them to 0.025 phase bins, then ordered and trimmed the data within each bin for each filter, taking out the

upper and lower 10 per cent (*B* and *V*) or 15 per cent (*U*) of the data (as judged by a median value of the given bin) and calculating the average values for the binned and trimmed data in each filter. Then we averaged the individual *U*, *B* and *V* light curves, thus reaching an accuracy of  $\sigma \leq 0.001$ –0.002 mag for a binned data point. Detecting no significant variability at phases 0.1–0.8, we increased the bin size to 0.1 while combining the individual *UBV* light curves at phases 0.1–0.8, keeping the sufficiently high 0.025 phase resolution around periastron (Fig. 1). The 0.025 phase resolution was chosen after some experimentation as a reasonable compromise between the simultaneously required high accuracy ( $\leq 0.002$  mag per bin) and high time resolution. Inspecting the average light curve in Fig. 1 one may notice the somewhat larger error bars for the points closest to phase zero. One may also note the deviation of the *B* lightcurve from the *U* and *V* data around  $\phi = 0.9625$  and 0.0125 respectively (Fig. 1, upper panel). It is tempting to assume that enhancement of the intrinsic variability of the system around periastron could be a possible source of the deviations. This is supported by the fact that the individual light curves comprising the data were taken at different orbital cycles. Indeed, the time sampling of the 0.025 phase bins at  $\phi = 0.9625$  and 0.0125 is slightly different: the  $\phi = 0.9625$  bin contains five observations in *U* and five in *V* of the same time sequence, but only four *B* points; the  $\phi = 0.0125$  bin includes five *U*, four *V* (all times are overlapping with *U*) and six *B* points. With the current data in hand, the suggestion of the enhanced intrinsic variability during periastron passage may be regarded only as tentative, pending acquisition of data of much higher quality.



**Figure 1.** Upper panel: the individual differential *U* (triangles), *B* (open circles) and *V* (filled circles) observations of  $\iota$  Ori, each with the mean subtracted, and grouped in 0.025 phase bins (see text). Lower panel: the *UBV*-averaged light curve (open circles with  $2\sigma$  error bars) of  $\iota$  Ori folded with  $P_{\text{orb}} = 29.133\,76$  d and binned to 0.1 and 0.025 phase bins with the modelled light curve (solid line) corresponding to  $i = 60^\circ$  orbital inclination. The arrows point to the moments of potential eclipses.

### 3 RADIAL VELOCITY MEASUREMENTS

In order to obtain the most precise radial velocity (RV) measurements possible, we first constructed a template for the primary star. Gaussian (O III 5592, C IV 5801, 5812) or Voigt (He I lines, H $\alpha$ ) profiles were fitted to the lines of the primary at phases when the primary and secondary spectra are well separated (i.e. around periastron passage). We created a zero-approximation artificial spectrum of the primary by assigning to its components the averaged half-widths and intensities of the observed lines and placing all the artificial components at the laboratory wavelengths. This artificial spectrum was properly Doppler-shifted and compared with the composite (primary + secondary, originally observed) spectra at different phases, with emphasis on the parts of the composite profile the least affected by the presence of the secondary. We quickly realized that this artificial spectrum of the primary generally underestimates the line depths of the primary, if based solely on the data taken around periastron. Measuring line depths at out-of-periastron phases where the lines of the secondary, nevertheless, are sufficiently displaced from the profiles of the primary and accordingly adjusting the fitted parameters (mainly the line intensities), we finally constructed a compromise version of the artificial spectrum of the primary. We found the H $\alpha$  profile of the primary to be blueshifted relative to the O III and He I lines. This phenomenon (sometimes referred to as Balmer progression; cf. the discussion in Hutchings 1968) is probably indicative of an expanding atmosphere. Hence, we applied a  $\Delta RV = -28.4 \text{ km s}^{-1}$  correction to the laboratory wavelength of H $\alpha$  while creating the final version of the artificial spectrum of the primary.

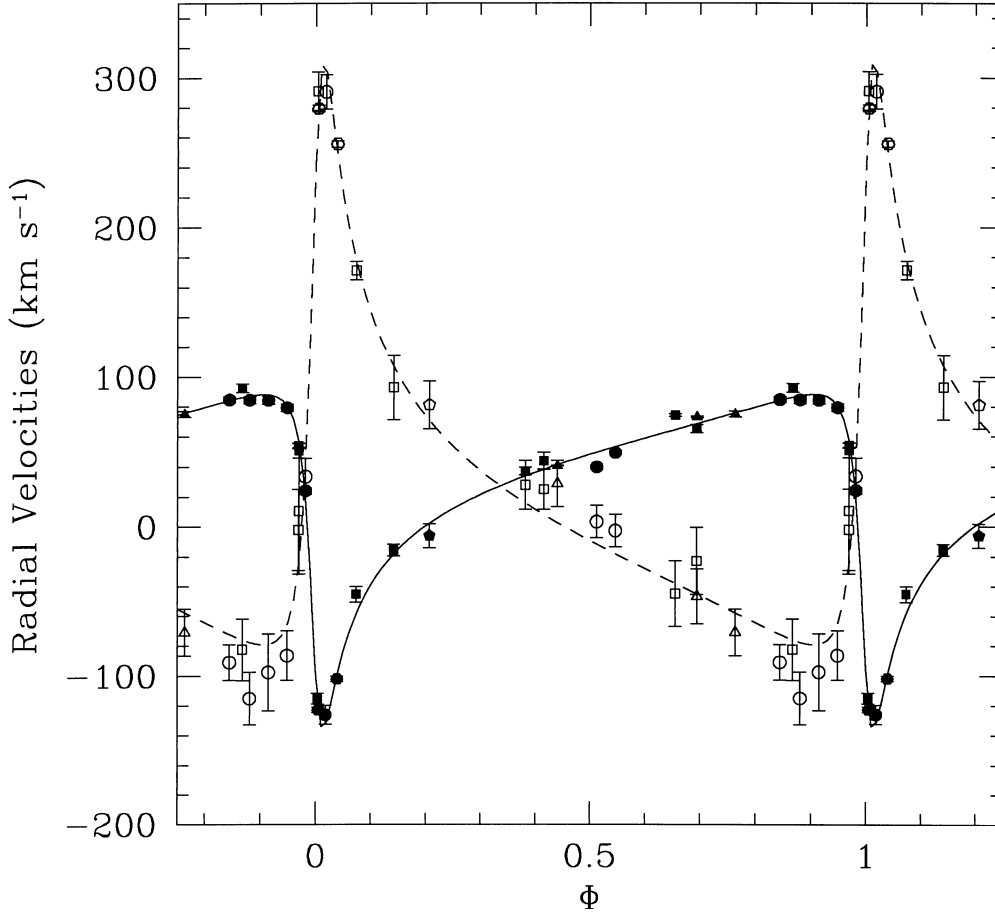
This artificial spectrum was then cross-correlated (the FXCOR task in IRAF) with the individual observed composite spectra. The RVs of the primary component measured in this way were

subsequently adjusted (usually by  $\Delta RV < 20 \text{ km s}^{-1}$ ) to provide the best match for the parts of the composite profile where the presence of the secondary was assumed to be negligible (line wings of the composite profiles for the phases 0.3–0.7; red side for the phases 0.7–0.9, etc.). Subtracting the appropriately shifted artificial spectrum of the primary from the composite spectra, we thus obtained a set of secondary spectra. These were then mutually cross-correlated, appropriately shifted and combined into a mean spectrum of the secondary. While constructing this secondary template, we omitted all the spectra taken around periastron, in the expectation that they could be affected by phase-locked spectral variability at those phases. Then we cross-correlated the secondary template with the individual difference spectra (composite spectrum minus appropriately Doppler-shifted artificial spectrum of the primary) to obtain more robust velocity estimations for the secondary component. Subtracting the primary (secondary) template from the appropriately (orbital motion) shifted composite spectra, we obtained a collection of phase-resolved secondary (primary) spectra to look for any phase-related profile variability. In Table 1 we list: the HJD dates calculated for the mid-exposure times; the source of the data; the orbital phases; the heliocentric RVs and corresponding  $\sigma(RV)$  as measured via cross-correlation, along with the calculated RVs. We note that the restored secondary and primary profiles are in fair agreement with the profiles obtained by Gies et al. (1993; 1996). Some of the quoted RV errors may seem unrealistically small, as they are based exclusively on the results of the cross-correlation procedure, thus neglecting the additional source of errors stemming from the uncertainty in the ‘absolute’, zero-point wavelength calibration. Unfortunately, the latter error cannot be uniformly assessed from the available data, as for a given observatory they sometimes comprise only 1–3 spectra covering different wavelengths. As we already mentioned, the velocity stability was better than  $1 \text{ km s}^{-1}$

**Table 1.** Measured radial velocities.

HJD–245 0000	Observatory <sup>a</sup>	Phase	RV (km s <sup>−1</sup> )	Primary $\sigma_{RV}$ (km s <sup>−1</sup> )	RV <sub>calc</sub> (km s <sup>−1</sup> )	RV (km s <sup>−1</sup> )	Secondary $\sigma_{RV}$ (km s <sup>−1</sup> )	RV <sub>calc</sub> (km s <sup>−1</sup> )
347.883	CASLEO	0.441	40.6	0.2	41.4	28.8	15.4	2.7
505.314	OHP	0.844	84.8	2.0	84.2	−90.8	11.9	−72.1
506.375	OHP	0.881	84.6	2.5	87.3	−115.0	17.7	−77.5
507.369	OHP	0.915	84.4	2.5	88.1	−97.4	25.8	−78.9
508.353	OHP	0.949	79.4	2.5	80.4	−86.2	16.6	−65.5
509.318	OHP	0.982	24.3	1.8	19.9	33.6	11.9	40.3
510.379	OHP	0.018	−125.9	6.4	−132.8	290.8	11.6	307.4
704.876	Ritter	0.695	65.3	2.8	68.7	−23.0	22.4	−45.0
704.893	CASLEO	0.695	73.1	0.7	68.7	−46.6	18.4	−45.0
706.899	CASLEO	0.764	75.0	2.0	79.9	−70.7	15.7	−64.6
709.928	Ritter	0.868	92.6	2.8	86.3	−82.3	20.7	−75.8
712.895	Ritter	0.969	54.1	1.9	57.3	−2.1	27.1	−25.1
712.903	Ritter	0.970	50.8	4.7	56.8	10.6	42.2	−24.2
713.901	Ritter	0.004	−115.1	3.6	−116.8	291.1	12.9	279.4
713.924	SPM	0.005	−122.7	1.9	−120.1	279.5	2.2	285.2
714.932	SPM	0.039	−101.9	1.8	−99.3	255.7	2.1	248.8
715.928	Ritter	0.074	−45.2	5.3	−58.6	171.4	6.2	177.6
717.915	Ritter	0.142	−15.6	3.8	−18.4	92.9	21.6	107.3
719.810	OMM	0.207	−6.0	8.0	2.3	81.2	16.0	71.1
724.914	Ritter	0.382	37.0	2.6	33.9	27.8	16.3	15.8
725.901	Ritter	0.416	43.9	5.7	34.3	25.0	13.4	15.1
732.877	Ritter	0.655	74.3	1.0	64.7	−44.8	22.1	−38.0
1136.580	OHP	0.512	39.8	1.1	49.6	3.4	10.9	−11.6
1137.574	OHP	0.546	49.4	1.1	53.3	−2.7	10.9	−18.1

<sup>a</sup>CASLEO – Complejo Astronómico El Leoncito (Argentina); OHP – Observatoire de Haute-Provence (France); SPM – the San Pedro Mártir Observatory (Mexico); OMM – Observatoire du Mont Mégantic (Québec); Ritter – the Ritter Observatory (USA).



**Figure 2.** Radial velocity curve of  $\iota$  Ori as determined from our RV measurements. The filled (open) symbols indicate the RVs of the primary (secondary) component: triangles – CASLEO; circles – OHP; pentagons – OMM; squares – Ritter; hexagons – SPM.

in the spectra obtained at Ritter Observatory. As for the rest of the spectra, the measurements of the interstellar lines (whenever possible) showed random deviations not exceeding  $10 \text{ km s}^{-1}$ . The overall uniformity of the data set can be evaluated by comparing the data plotted in Fig. 2, where different symbols correspond to different observatories. In the following RV analysis we refer to the cross-correlation errors only.

#### 4 ORBITAL PARAMETERS

A revised orbital solution for  $\iota$  Ori was derived from our RV measurements using an improved version of the algorithm of Wolfe, Horak & Storer (1967) and using  $1/\sigma$  RV weights (cf. Table 1). The problem of relatively slow convergence of the algorithm for orbits with  $e \gtrsim 0.8$  was circumvented by choosing an appropriate number of terms in the asymptotic expansion, four terms in our case. Since the radial velocities of the secondary are probably affected by line profile variations (see below), a simultaneous determination of all the orbital parameters from the data of both components seems inappropriate. We have thus obtained the orbital solution of the primary independently of the RVs of the secondary. The RV curve of the secondary was then derived after fixing the values of  $e$ ,  $\omega$  and  $T_0$  as determined for the primary star alone, with amplitude and apparent systemic velocity as the only free parameters.

Since there is strong evidence for apsidal motion in  $\iota$  Ori

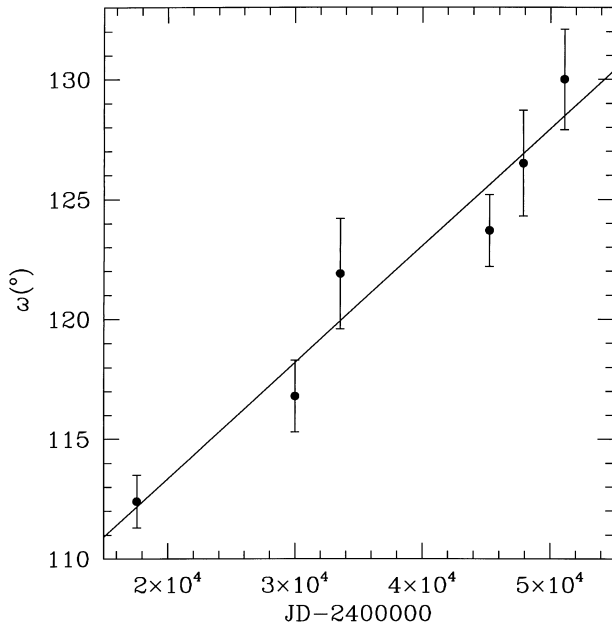
**Table 2.** New orbital solution derived from our RV measurements. The quoted uncertainties are  $1\sigma$  errors.

	Primary	Secondary
$P$ (days)	29.133 76(fixed)	
$\gamma$ ( $\text{km s}^{-1}$ )	$31.3 \pm 1.2$	$20.4 \pm 2.1$
$K$ ( $\text{km s}^{-1}$ )	$111.9 \pm 2.5$	$195.7 \pm 4.3$
$q = M_1/M_2$	1.749	
$e$	$0.764 \pm 0.007$	
$\omega(^{\circ})$	$130.0 \pm 2.1$	
$T_0$ (HJD–245 0000)	$1121.658 \pm 0.046$	
$M \sin^3 i$ ( $M_{\odot}$ )	15.0	8.5
$a \sin i$ ( $R_{\odot}$ )	41.5	72.7

(SPLH87), it is problematic to combine data sets from very different epochs to derive an improved orbital period. Therefore, we have restricted ourselves to the combination of our data with the most recent RV data set available in the literature (Hilditch et al. 1991), thus obtaining the most probable period of  $29.134\,34 \pm 0.000\,20 \text{ d}$ . Since the longitude of periastron increases with time, this value should provide an upper limit to the actual orbital period. We note, however, that at any orbital phase the RVs of the primary from the Hilditch et al. (1991) data set are systematically closer to the apparent systemic velocity than our data. Hilditch et al. adopted an empirical correction to account for the effects of spectral-line blending on the RVs of the primary.

It could be that this correction is not large enough to account for the entire effect of the blending. In the following, we will therefore focus on the analysis of our new data set alone.

We have tested several assumptions on the orbital period, adopting either the value 29.13376 d of SPLH87 or using the period of 29.13434 d, as derived above. Both options yield orbital elements that overlap within their errors. The solution obtained adopting the period proposed by SPLH87 is given in Table 2 and displayed in Fig. 2.

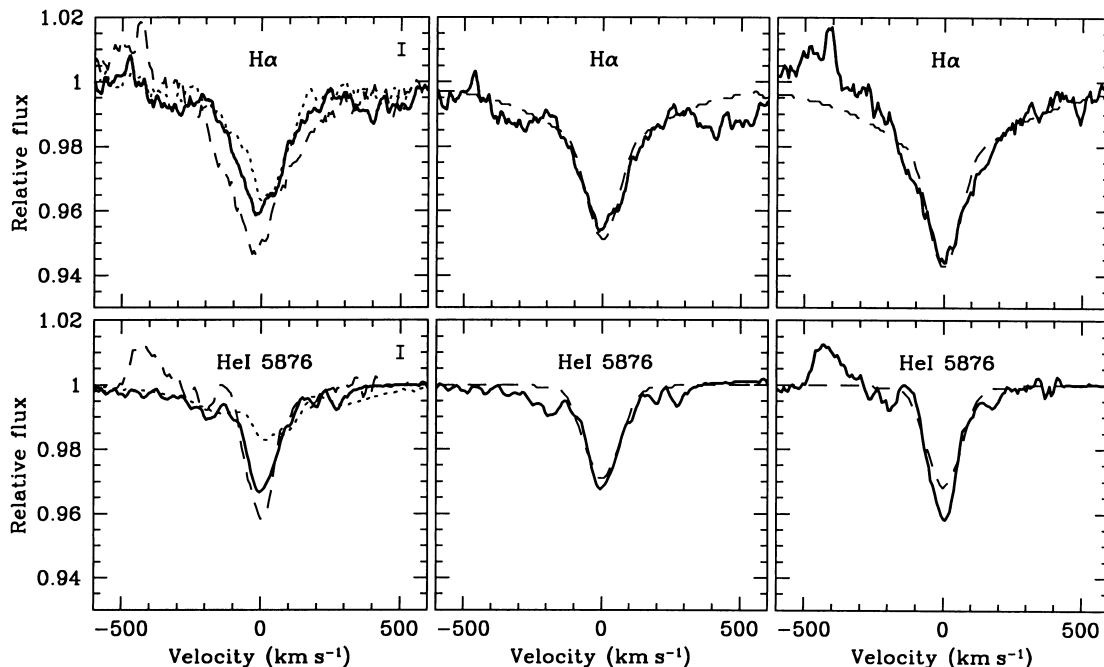


**Figure 3.** Least-squares fit to the observed variations of the longitude of periastron in  $\iota$  Ori (see text).

The eccentricity of our solution is in excellent agreement with the value from SPLH87. A rather unexpected finding is the difference in the apparent systemic velocities of the two components ( $\gamma_1 - \gamma_2 = 10.9 \text{ km s}^{-1} \approx 5 \sigma$ ). Fitting the RVs of the secondary as quoted by SPLH87, we find that the same discrepancy also exists in their data set. This phenomenon is quite common in evolved O+O binaries, and is usually attributed to the influence of a stellar wind on the formation of absorption lines. In this interpretation, the component with the strongest wind should display the most blue-shifted absorptions. However, in the case of  $\iota$  Ori, it is the secondary component with the presumably weaker wind which has the bluer apparent systemic velocity. We therefore believe that the  $\gamma$ -velocity difference in  $\iota$  Ori could be an artefact resulting from the phase-locked line profile variability of the secondary star (see Section 5).

Considering apsidal motion, we notice that the longitude of periastron in our solution ( $\omega = 130^\circ 0 \pm 2^\circ 1$ ) is larger than the value ( $126^\circ 2$ ) obtained from an extrapolation of the empirical linear relation given by SPLH87. We have therefore performed a least-squares linear fit (see Fig. 3) in time to all the available measurements of  $\omega$  (Plaskett & Harper 1909; Pearce 1953; Miczaika 1951; SPLH87; Hilditch et al. 1991; this paper). The best fit corresponds to a rate of apsidal motion  $\dot{\omega} = 0^\circ 00049 \pm 0^\circ 00003 \text{ d}^{-1}$ , slightly larger than the value derived by SPLH87 ( $0^\circ 00041 \text{ d}^{-1}$ ).

The mass ratio  $q = m_1/m_2 = 1.75$  in our solution is lower than the value found by SPLH87,  $q = 2.05$ . However, since this difference arises from an increase of  $K_1$  and a decrease of  $K_2$  (the RV amplitudes of the primary and secondary, respectively) in our data compared with SPLH87, the total semimajor axis of the system  $a \sin i = 114.2 R_\odot$  is nearly unchanged with respect to the results of SPLH87 ( $112.8 R_\odot$ ) and therefore our revision of the orbital elements of  $\iota$  Ori should have little impact on the modelled colliding wind effects (Pittard 1998). Concerning the predicted



**Figure 4.** Phase-averaged line profiles of the secondary component. Left panel: full line – phases 0.1–0.9 (each profile: mean of 10 He I profiles and 11 H $\alpha$  profiles); dotted line – phases 0.9–1.0 (mean of 2 He I profiles and 5 H $\alpha$  profiles, respectively); long-dashed line – phases 0.0–0.1 (mean of 4 He I profiles and 3 H $\alpha$  profiles, respectively). Typical  $2\sigma$  error bars are shown in upper right corners of the panel. Middle panel: full line – phases 0.1–0.9, observed profiles; long-dashed line – modelled profiles. Right panel – the same as in middle panel, but for phases 0.0–0.1.

proximity effects (Stevens 1988; Moreno & Koenigsberger 1999), we notice that the downward-revised mass ratio will lead to a smaller critical volume of the primary near periastron than expected from the SPLH87 solution. Hence, the predicted deformations of the primary could be somewhat larger than the model estimates (Moreno & Koenigsberger 1999).

## 5 DISCUSSION: OBSERVED PROXIMITY EFFECTS

### 5.1 Spectroscopy

In  $\iota$  Ori the phase-locked profile variations induced by the rapid change of separation between the components during periastron passage must be disentangled from the stochastic line variability affecting the primary spectrum. The latter type of variability was not detected in spectra acquired within the same night, but was clearly present as a night-to-night and cycle-to-cycle component. However, in  $\iota$  Ori these stochastic variations are generally of small intensity and usually confined to the cores of the profiles (cf. Gies et al. 1993).

Calculating the standard deviations for all but two (phases 0.0–0.1 – see below) H $\alpha$  profiles, we found that the stochastic component can cause line profile variations with full amplitude up to 3 per cent (referring to the continuum). When possible, we minimized the influence of this stochastic variability by: (a) averaging the spectra taken at different cycles; and (b) grouping the spectra within wide phase bins. Following this simple recipe, we immediately find that all the line profiles of the secondary undergo dramatic changes during periastron passage (Fig. 4). To demonstrate this, all the individual profiles were Doppler shifted to the frame of the secondary in accordance with the orbit of the secondary and grouped within 3 phase bins: 0.0–0.1, 0.1–0.9, 0.9–1.0. We found the profiles to be somewhat depleted before periastron ( $\phi = 0.9$ –1.0) passage and greatly enhanced right after that ( $\phi = 1.0$ –1.1) compared with the ‘neutral’ mean profiles taken at  $\phi = 0.1$ –0.9. The additional emission at  $v = -(500 - 300)\text{ km s}^{-1}$  is an artefact created by the phase-related variability of the profiles of the primary (see below). Note that at phases  $\phi = 0.9$ –1.0 (and *only* around these phases) we are able to look at the rear side of the secondary star, which is not affected by the flux coming from the primary.

The deepening of the profiles of the secondary after periastron has been noted by Gies et al. (1996) and could be related to (i) illumination of the hemisphere of the secondary component by the nearby and far more luminous primary star and, additionally, to (ii) the Struve–Sahade effect, which, among other possibilities, might be caused by heating of the secondary by the radiation coming from the wind–fixed solar collision zone (either back-scattered photospheric flux or direct X-ray irradiation: Gies et al. 1997a). To verify this hypothesis, we calculated grids of the He I 5876 and H $\alpha$  line profiles using the non-local thermodynamic equilibrium code TLUSTY (Hubeny & Lanz 1995) for an H (15 levels) and He (14 levels) atmosphere, turbulent velocity of  $5$ – $25\text{ km s}^{-1}$  (Hubeny & Lanz 1992), gravitational acceleration of  $\log g = 4.1$  [centimetre-gram-second (cgs) units; in correspondence with the values derived from the light curve fit; (see below), rotational velocity  $v_e \sin i = 50$ – $100\text{ km s}^{-1}$ , primary/secondary flux ratio in the optical  $I_1/I_2 = 7.9$  (see Section 5.2)] and variable effective temperatures, keeping fixed solar chemical composition. Apparently, the outcome of the modelling should depend on the initial choice of  $\log g$  and chemical composition. However, our intention

will be to follow the differential, phase-related changes in the surface temperature. We are aware that the absolute values might be uncertain at a level exceeding the quoted errors. However, we argue that the differences in the temperatures are less sensitive to the choice of He/H ratio,  $v_{\text{turb}}$ ,  $\log g$ , etc., as those parameters, once chosen, remain fixed for all the spectra. Fine tuning of the model providing accurate absolute temperatures could be done only via incorporation of wide-range spectral data with an adequate coverage of orbital phases. Unfortunately, we lack such data.

The modelling provides very good fits (Fig. 4) for the averaged profiles corresponding to the 0.1–0.9 phases, with  $T_{\text{eff}} = 25\,000 \pm 1000\text{ K}$ , in coincidence with  $T_{\text{eff}} = 24\,000$ – $25\,400\text{ K}$  for B1 III–V stars [Schmidt-Kaler 1982; note that the small radius of the B1 component as derived from the light curve fitting (below) has a better correspondence with luminosity class V than the previously accepted III],  $v_{\text{turb}} = 25\text{ km s}^{-1}$  and  $v_e \sin i = 75 \pm 10\text{ km s}^{-1}$  (in the range of the previously published  $v_e \sin i$  estimates for the secondary,  $v_e \sin i = 68$ – $80\text{ km s}^{-1}$ : Gies et al. 1993, 1996; Howarth et al. 1997). Despite the relatively low quality of the fits (Fig. 4) to the phase 0.0–0.1 profiles (fixed  $v_{\text{turb}} = 25\text{ km s}^{-1}$  and  $v_e \sin i = 75\text{ km s}^{-1}$ ), we are able to conclude that the H $\alpha$  profile can be fitted only assuming much lower  $T_{\text{eff}} = 19\,000 \pm 1000\text{ K}$ . At these phases, there is no possibility of matching the observed He I profile, with its central part being far too strong for any reasonable choice of  $T_{\text{eff}}$ . Note also the broadening of the observed H $\alpha$  profile.

It is clearly futile to attempt to find precise  $T_{\text{eff}}$  values for the asymmetric, peculiar shaped (especially He I 5876 line) profiles corresponding to 0.9–1.0 phases. They are indicative of  $T_{\text{eff}} \gtrsim 28\,000\text{ K}$ . Although the absolute values of the temperatures might be somewhat biased by the initial choice of parameters, there is little doubt that we observe fairly large, phase-dependent line profile variations, which we attempted to model as photospheric temperature variations. We may conclude that they cannot be caused by the changing separation between the components, as the more luminous primary has practically negligible influence on the profiles of the secondary, being of comparable  $T_{\text{eff}}$ ; we have been able to check this by introducing an external illumination factor in TLUSTY.

Could the variations be the result of the crashing of the bow shock on to the secondary? Note that from general considerations one expects to see significant imbalance in the momenta of the wind (up to 100 times: SPLH87). In a simplistic approach, at  $\phi = 0.0$ –0.1 the shock zone is seemingly detached from the atmosphere of the secondary; otherwise, the profiles of the secondary would be significantly depleted, as the crashing bow shock would completely engulf the H $\alpha$  formation zone, reaching the region of formation of He I lines as well, and raising the inside-zone temperature by  $\sim$ two orders of magnitude. However, the real situation might be more complicated, with everything related to the relaxation time-scale for the zone of the shocked gas. This may range from hours to days, critically depending on the poorly known velocity of the wind of the primary in the vicinity of the surface of the secondary: any extrapolation of the widely accepted  $\beta$ -velocity law (cf. Pittard et al. 2000) is very risky, owing to the

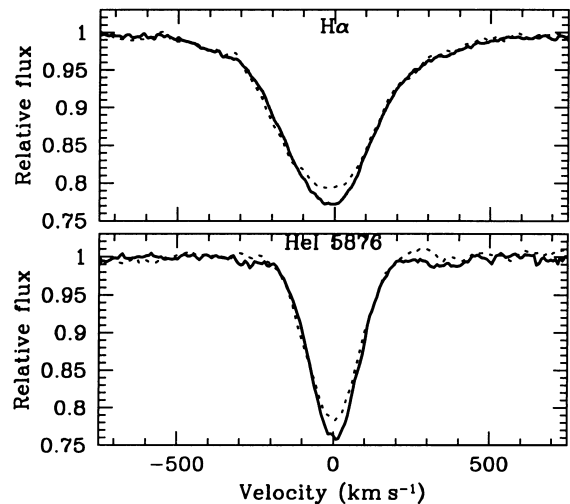
<sup>1</sup> The chosen value of  $v_{\text{turb}}$  seems rather high. However, our choice was motivated not only by the quality of the model fits, but also by the matter of convergence of the ‘low-temperature’ models,  $T_{\text{eff}} = 15\,000$ – $20\,000\text{ K}$ . This particular choice should not much affect the conclusions, as long as we are interested only in differential temperature changes.

fact that the wind braking effect (a sudden deceleration of the wind of the primary by the radiation field of the secondary: Gayley, Owocki & Cranmer 1997; see also the discussion in Pittard et al. 2000) is presently only schematically accounted for. Assuming that the zone of shocked gas can settle to some equilibrium within hours (the assumption being supported by the extremely efficient radiative cooling in the relatively high-density region of the secondary's atmosphere), we proceed by suggesting that the observed 'low-temperature' H $\alpha$  profile at  $\phi = 0.0$ – $0.1$ , as well as the anomalously strong He I 5876 line, can be formed *inside* the wind–wind collision zone – more specifically, in the decelerated and cooled part of the wind of the primary. In a rather speculative vein, this assumption may be counteracted by two arguments. (i) The dynamical equilibrium reached by the shocked gas could be continuously violated during periastron passage, when the rapidly changing separation can significantly affect the ram pressure balance. (ii) The rapidly cooled wind–wind collision zone is prone to large-scale, turbulent gas motions (Stevens, Blondin & Pollock 1992; Walder & Folini 1998). It remains to be seen whether the apparently smooth *absorption* profile can be formed in such a violent, fractalized medium; however, note the apparent broadening of the H $\alpha$  profile at phases 0.0– $0.1$  compared to the model fit (Fig. 4).

More conventional solutions of the non-crashing paradox suggest either an extremely extended acceleration zone in the wind of the primary, or an unexpectedly strong wind braking effect. Another, though rather debatable, possibility that would change the wind momentum ratio in favour of the secondary, thus lifting the shock from the secondary's surface, is to assume that the wind velocity of the primary can be reduced by strong tidal interaction effects (a somewhat higher mass-loss rate with the related significant change in acceleration rate in the outer region of the primary's wind). However, the change in wind velocity would lead to a significant change in the hardness of the X-ray emission near periastron. The fact that we do not see this (Pittard et al. 2000) either implies that the X-ray emission does not come from the wind interaction region, or that there is something wrong with our expectation of a *considerably* changed wind structure at phase 0.0.

The phase-locked profile variations in the spectrum of the primary are more subtle and are evident only in the spectra taken around  $\phi = 0.00 \pm 0.01$  (Fig. 5). We confirm the finding of Gies et al. (1996): for a short time the central parts of the profiles of the primary are partially filled in by line emission probably arising in a tidal stream created by the gravitational pull of the secondary component. The variability in the red wing of the He I 5876 profile of the primary is caused by the phase-related changes in the spectrum of the secondary (as above). Simulations of the tidal stream emissivity will require incorporation of radiation transfer in a perturbed, three-dimensional stellar atmosphere, which is far beyond the scope of this paper. It is quite surprising that practically all interaction effects are restricted to such a short time interval: note also the optical light curve maximum (see below) and the 'spike' in the X-ray light curve (Pittard et al. 2000). However, it is consistent with the results of the tidal interaction model of Moreno & Koenigsberger (1999), who find that the time-span over which mass motions on the tidal bulge have velocities larger than  $1 \text{ km s}^{-1}$  is  $\sim 10 \text{ h}$ .

We found that a relatively strong ( $\sim 5$  per cent of the continuum) C III 5696 emission line follows the orbit of the primary, with some indication of phase-related variability. However, the emission was placed too close to the edge of the échelle order to achieve sufficient accuracy for quantifying those



**Figure 5.** Phase-averaged line profiles of the primary component: full line – all phases except  $\phi = 0.003$ – $0.005$  (mean of 12 He I profiles and 17 H $\alpha$  profiles); dotted line – phases 0.003– $0.005$  (mean of two spectra).

changes. Note that this line appears in emission in practically all luminous late-type O stars (Conti 1974; Walborn 1980).

## 5.2 Photometry

To interpret the rapid light curve variability taking place during periastron passage (Fig. 1), we applied a numerical code (Antokhina & Cherepashchuk 1994; Antokhina 1996) designed to calculate a light curve for an eccentric orbit under the standard assumptions (Wilson 1979) that (i) the shapes of the stars coincide with instantaneous equipotential Roche surfaces at all orbital phases, and (ii) the stars retain constant volumes during orbital revolution. The model takes into account changes in shapes of the components, mutual radiative heating, and possible geometrical eclipses. Note that at such a large eccentricity ( $e = 0.76$ ), the standard Roche model is far from being well justified. Near periastron, strong tidal interaction and resulting instabilities may distort the shapes of the components. In fact, we see some signs of such interaction: the profile variability in the primary component (Fig. 5) and rapid change of the longitude of periastron. Stevens (1988) used the equipotential Roche surface assumption and derived the variation in the (maximum) radius of  $\iota$  Ori's primary star between periastron and apastron, obtaining a difference of  $0.39 R_{\odot}$  in the radii at these two orbital phases. However, Moreno & Koenigsberger (1999) obtain a change in the radius of only  $0.02 R_{\odot}$  from the solution of the equations of motion of the surface elements in the equatorial region, assuming that the stellar radius is  $15.8 R_{\odot}$ . Furthermore, the shapes of the stars at periastron differ from those at other phases, having a prolate spheroidal appearance. Thus, the following results obtained with the standard Roche model should be treated only as indicative, pending the inclusion of the precise shapes of the stars into the light curve calculation method.

As model input we use the known spectroscopic values of  $M_1 \sin^3 i$ ,  $M_2 \sin^3 i$ ,  $P$ ,  $\omega$ ,  $e$ ,  $T_0$  (Table 2), as well as averages of the previously published data on  $v_e \sin i = 110 \text{ km s}^{-1}$  (primary) and  $70 \text{ km s}^{-1}$  (secondary), characteristic values of  $T_{\text{eff}} = 32\,500 \text{ K}$  (primary) and  $T_{\text{eff}} = 24\,000 \text{ K}$  (secondary). Gravity darkening coefficients  $\beta_1 = \beta_2 = 0.25$  and albedos  $A_1 = A_2 = 1$  are assumed as typical for early-type stars. The square-root limb



**Table 3.** Parameters of the system calculated for different inclinations.

$i$ deg	$\mu_1$	$\mu_2$	$M_1$ $M_\odot$	$M_2$ $M_\odot$	$a$ $R_\odot$	$R_1$ $R_\odot$	$R_2$ $R_\odot$	$F_1$	$F_2$	$L_1$ $\text{erg s}^{-1}$	$L_2$ $\text{erg s}^{-1}$	$\Delta(o - c)$
30	0.65	0.30	120.	68.	229.	14.3	3.8	8.8	20.9	$7.9 \times 10^{38}$	$1.7 \times 10^{37}$	3.86
40	0.75	0.30	56.5	32.1	178.	13.3	3.0	7.4	20.9	$6.8 \times 10^{38}$	$1.0 \times 10^{37}$	2.42
45	0.75	0.30	42.4	24.0	162.	12.7	2.7	7.4	20.9	$5.7 \times 10^{38}$	$8.5 \times 10^{37}$	2.00
50	0.70	0.45	33.4	18.9	149.	10.3	4.7	8.0	11.2	$4.1 \times 10^{38}$	$2.5 \times 10^{37}$	1.81
55	0.70	0.50	27.3	15.5	140.	9.6	5.0	8.0	9.8	$3.6 \times 10^{38}$	$2.9 \times 10^{37}$	1.74
60	0.65	0.55	23.1	13.1	132.	8.3	5.4	8.8	8.7	$2.6 \times 10^{38}$	$3.3 \times 10^{37}$	1.73
65	0.55	0.65	20.2	11.4	126.	6.3	6.2	11.0	7.1	$1.5 \times 10^{38}$	$4.5 \times 10^{37}$	1.73
70	0.60	0.55	18.1	10.3	122.	6.9	5.0	9.8	8.7	$1.8 \times 10^{38}$	$2.8 \times 10^{37}$	1.77
75	0.60	0.35	16.6	9.4	118.	6.9	2.6	9.8	16.1	$1.7 \times 10^{38}$	$7.7 \times 10^{36}$	2.19
80	0.50	0.30	15.7	8.9	116.	5.1	1.2	12.7	2.9	$9.9 \times 10^{37}$	$4.4 \times 10^{36}$	3.11

darkening law (Van Hamme 1993) is used for calculations of the limb darkening. We choose  $\lambda_0 = 4400 \text{ \AA}$  for the calculated monochromatic light curve, in correspondence with the central wavelength of the photometric  $B$  band.

With the majority of the fundamental parameters fixed, we place restrictions on the orbital inclination. We start from calculating a grid of models with  $\Delta i = 10^\circ$  (afterwards changing to  $\Delta i = 5^\circ$  around the optimal  $i$  value), using  $\mu_1, \mu_2$  (the Roche lobe filling factors for the primary and secondary, respectively, calculated for the moment of periastron passage) as free parameters and attempting to find the best fit to the observed light curve using Fisher's statistical criterion (Balog, Goncharukii & Cherepashchuk 1981). As it turns out, the optimal solution is not very sensitive to the choice of  $i$ , which can be partially explained by the extremely small amplitude of the light curve and, as a consequence, by the *comparatively* large observational errors ( $\sim 0.001\text{--}0.002 \text{ mag point}^{-1}$ ). However, the primary reason for this is that the light variations are caused exclusively by tidal deformations of the components, not by eclipses. In the absence of eclipses, changes in inclination may be, to a large degree, compensated by appropriate changes in radii, so that the observations do not further constrain the radii and the inclination. In almost all cases one has to choose between two or more statistically equal solutions: usually, one with small  $\mu_1$  and large  $\mu_2$  and the other(s) with opposite  $\mu_1/\mu_2$  ratio. Our selection was based on the fact that the primary star is bigger, hotter, more massive and more luminous, with  $\Delta m = 2.0 \pm 0.2 \text{ mag}$  in the optical (SPLH87). We list the derived values in Table 3, where  $\mu_1, \mu_2$  are the Roche lobe filling factors at periastron;  $a$  is the semimajor axis;  $F_1$  and  $F_2$  are the ratios of surface rotation rate to synchronous rate for the primary and the secondary, respectively, calculated as  $F_j = \omega_j/\omega_{\text{kepl}} = (v_j/R_j)/(2\pi/P)$ , where  $v_j$  are taken from the known  $v_j \sin i$ ;  $L_1, L_2$  are the bolometric luminosities of the components and  $\Delta(o - c)$  is the normalized deviation for a given model of the light curve. At a confidence level of 1 per cent (referring to Fisher's statistics), the critical value is  $\Delta(o - c) = 1.90$ , and  $\Delta(o - c) = 1.59$  at 5 per cent. This provides a 99 per cent confidence interval for the inclination,  $50^\circ \leq i \leq 70^\circ$  with the optimal value at  $i = 60^\circ$  (the modelled light curve is shown in Fig. 1) and corresponding  $R_1 = 8.3 \pm 0.8 R_\odot, R_2 = 5.4 \pm 1.4 R_\odot$ , and  $M_1 = 23.1 M_\odot, M_2 = 13.1 M_\odot$ , with the masses being in reasonable agreement with the spectral types of the components. The derived radii are substantially smaller than the estimates from SPLH87, where the radii were calculated under the questionable assumption that the minimum in the light curve that is observed before periastron passage corresponds to a grazing eclipse. In our model, for the optimal value of  $i = 60^\circ$  the observed variations

could be interpreted as being induced by the proximity of the components alone, without eclipsing. However, smaller radii imply much smaller tidal deformations (Moreno & Koenigsberger 1999) and it is not clear whether such small deformations would produce the observed variations.

The new values of the radii point to luminosity class V rather than the traditionally assigned III. This immediately poses an additional problem: the system, judging by its total luminosity (case  $i = 60^\circ$ , Table 3), is a factor of 3 fainter than can be deduced from the mean distance modulus of the Ori OB1 association, considering  $\iota$  Ori as a member (cf. SPLH87). However, note that differences in the individual distance moduli of the association members may reach  $\pm 1 \text{ mag}$  (Humphreys 1978).

We reiterate our conclusion reached in the previous section, that heating by the *luminous* primary component cannot cause the observed phase-locked profile variability of the secondary star. Within our model, the calculated rise in the surface temperature of the secondary during periastron passage never exceeds 150 K.

The calculated asynchronous rotation rates deserve additional comment. As the values of  $F_i$  cannot be fixed initially, depending non-linearly on  $\mu_i$ , they have to be calculated iteratively. For a given  $\mu_i$  we set  $F_i = 1.0$ , then calculate the equipotential surface, its volume and mean radius, and, finally the new  $F_i$ . As these new  $F_i$  would deviate from 1, we re-solve the equation for the equipotential surface for the new  $F_i$ , repeating the process until convergence. This procedure ensures that values of  $v_e \sin i$  in our model match the observed values,  $\sim 110$  and  $\sim 70 \text{ km s}^{-1}$  for the primary and the secondary, respectively. One may question the apparently high values of  $F_1, F_2$  given in Table 3. However, note that they are defined relative to the average orbital angular rate,  $\omega_{\text{kepl}} = 2\pi/P$ . Comparison of the angular rotation rates of the components and the orbital angular rate *at periastron* provides significantly smaller values,  $\sim 0.5$  (Gies et al. 1996). Indeed, combining the newly derived values of  $R_1$  and  $R_2$  for  $i = 60^\circ$  with the observed values of  $v_e \sin i$  and assuming collinearity of the rotational and orbital axes, we obtain  $F_1 = F_2 = 0.28$  at periastron.

## 6 CONCLUSIONS

Combining new, high-quality spectroscopic data with abundant photometric measurements spanning over two years, we derive a complete set of orbital parameters for the massive, highly eccentric O9 III+B1 III binary  $\iota$  Orionis (Tables 2, 3). The newly derived mass ratio,  $q = m_1/m_2 = 1.75$ , is somewhat lower than the value found by SPLH87,  $q = 2.05$ . We confirm the high rate of apsidal motion, with our new estimation  $\dot{\omega} = (0^\circ.00049 \pm 0^\circ.00003) \text{ d}^{-1}$ .

Separating the heavily blended spectra of the components, we were able to see signs of tidal interaction taking place during the very short interval of periastron passage,  $\phi = 0.0 \pm 0.1$ . To our surprise, we find no clear evidence of the bow shock crashing on to the surface of the secondary even at periastron, i.e. presumably at the moment of maximum imbalance between the wind momenta.

The rapid, phase-locked photometric changes around periastron passage can be interpreted as arising from tidal deformations of the components (assuming that the shapes of the stars coincide with instantaneous equipotential Roche surfaces) and mutual illumination, without the necessity of invoking geometric eclipses. This leads to the relatively small values of the stellar radii (Table 3) and provides an estimate for the orbital inclination,  $50^\circ \lesssim i \lesssim 70^\circ$ .

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