

From perfect liquid to viscous fluid: hydrodynamics at RHIC

Pasi Huovinen

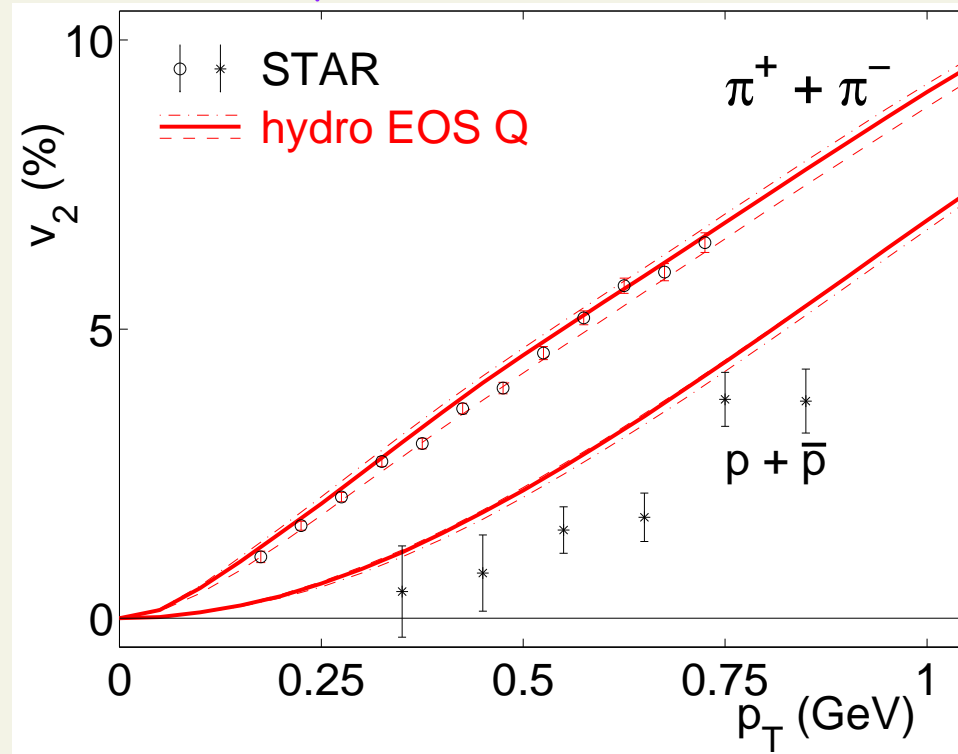
Purdue University

Heavy Ion Tea

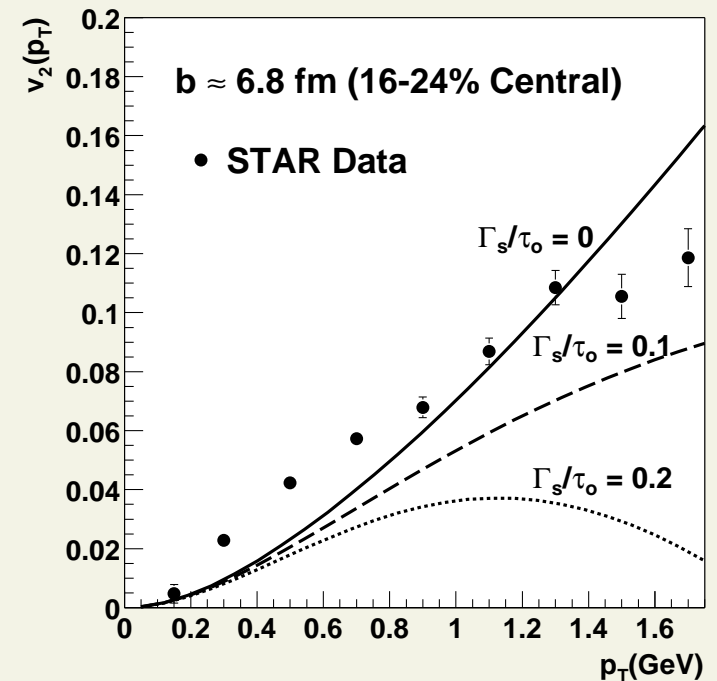
March 4, 2008, [LBNL](#), Berkeley, CA

Success of ideal hydrodynamics

Kolb, Heinz, Huovinen et al ('01)
 minbias Au+Au at RHIC



Teaney ('04): $\eta/s \propto \Gamma_s/\tau_0$



dissipation \rightarrow reduction of v_2

\Rightarrow the idea of plasma as “perfect fluid”

- but how perfect? - what if $\eta/s > 0$, e.g. conjectured $\eta/s \geq 1/(4\pi)$
- tune the calculation to allow large viscosity?

Ideal hydrodynamics

space-time evolution given by: $\partial_\mu T^{\mu\nu}(x) = 0$
 $\partial_\mu N^\mu(x) = 0$

- local conservation of energy, momentum and baryon number

local equilibrium, no dissipation: $T^{\mu\nu} = (e + p)u^\mu u^\nu - P g^{\mu\nu}$
 $N^\mu = n u^\mu$

local, macroscopic variables: energy density $e(x)$
pressure $P(x)$
flow velocity $u^\mu(x)$

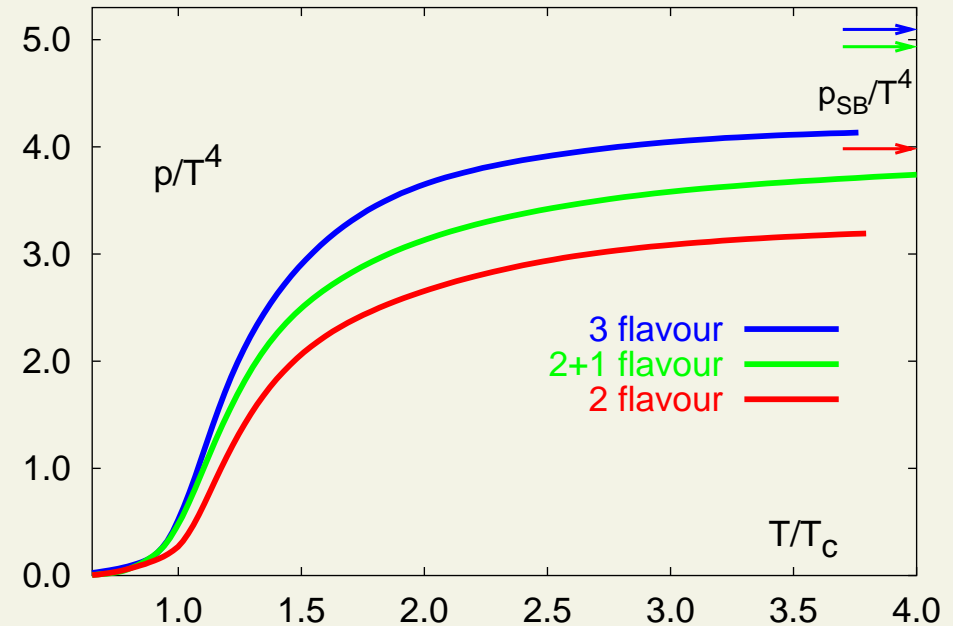
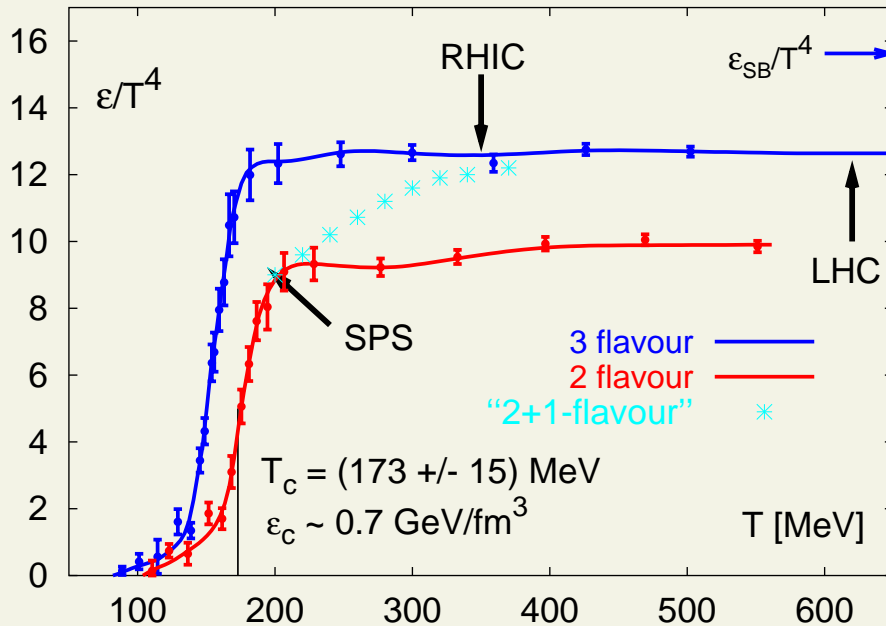
NO dissipation \rightarrow conserves entropy $\partial_\mu (s u^\mu) = 0$

matter characterized by: equation of state $P = P(e, n)$

Unknowns: initial state, final state, equation of state

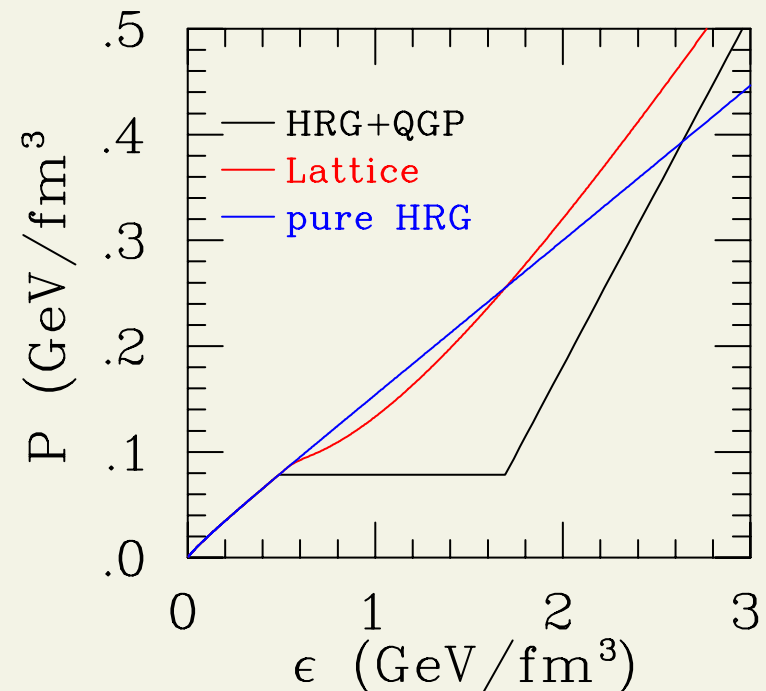
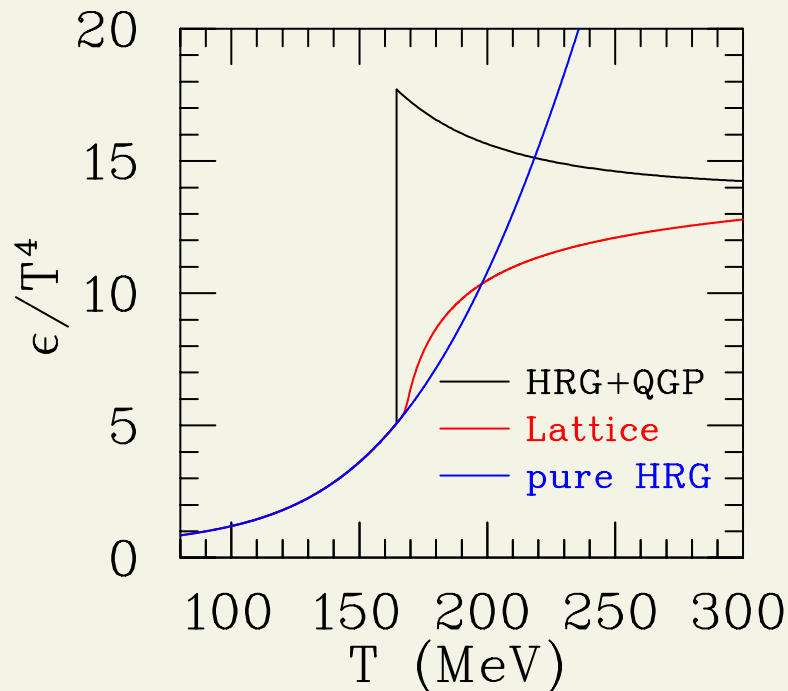
QCD equation of state

lattice QCD (Karsch & Laermann, hep-lat/0305025)



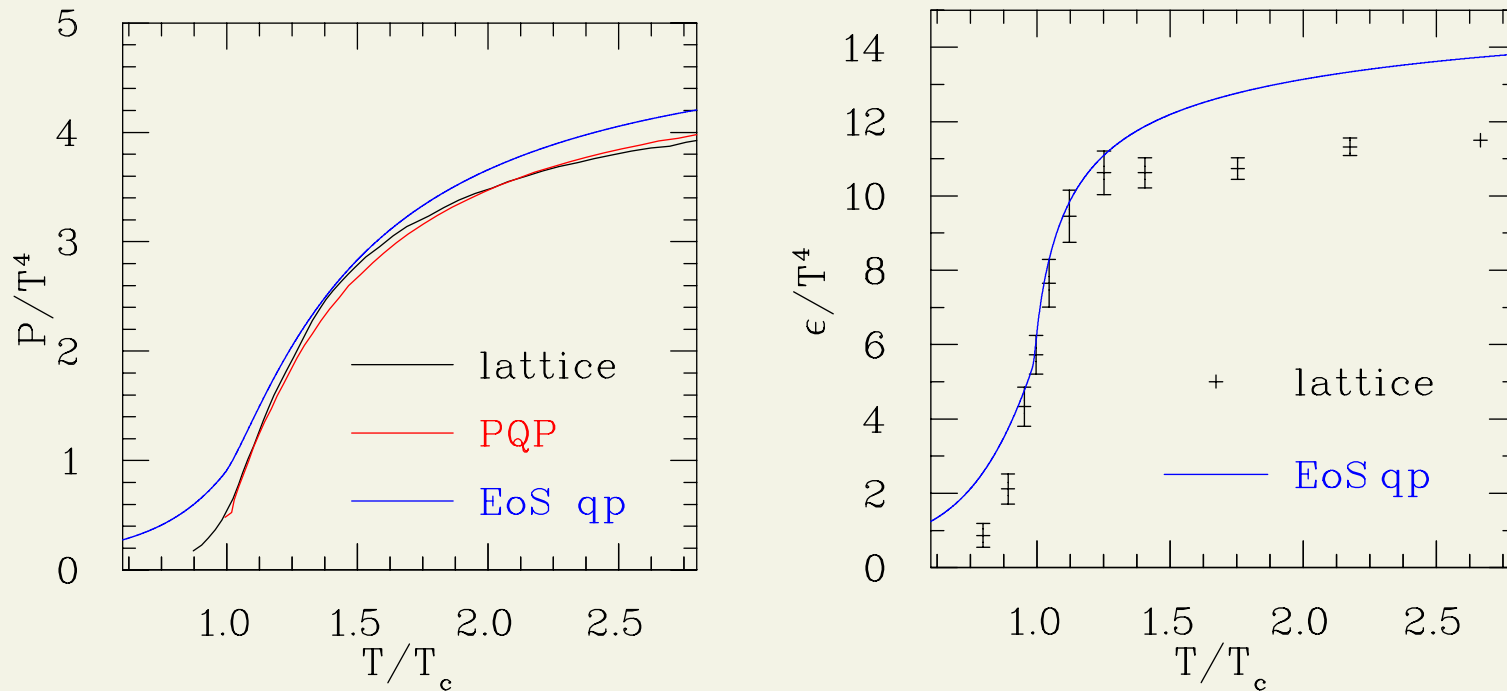
- jump in effective degrees of freedom at $T_c \approx 170 - 190 \text{ MeV} \sim 2 \cdot 10^{12} \text{ K}$
- phase transition is a smooth crossover
- hydro calculations: usually first order phase transition

Equation of state



- **HRG+QGP**: ideal gas of partons and hadron resonance gas, first order PT
- **Lattice**: parametrized lattice EoS + hadron resonance gas, crossover
- **pure HRG**: hadron resonance gas only, no PT

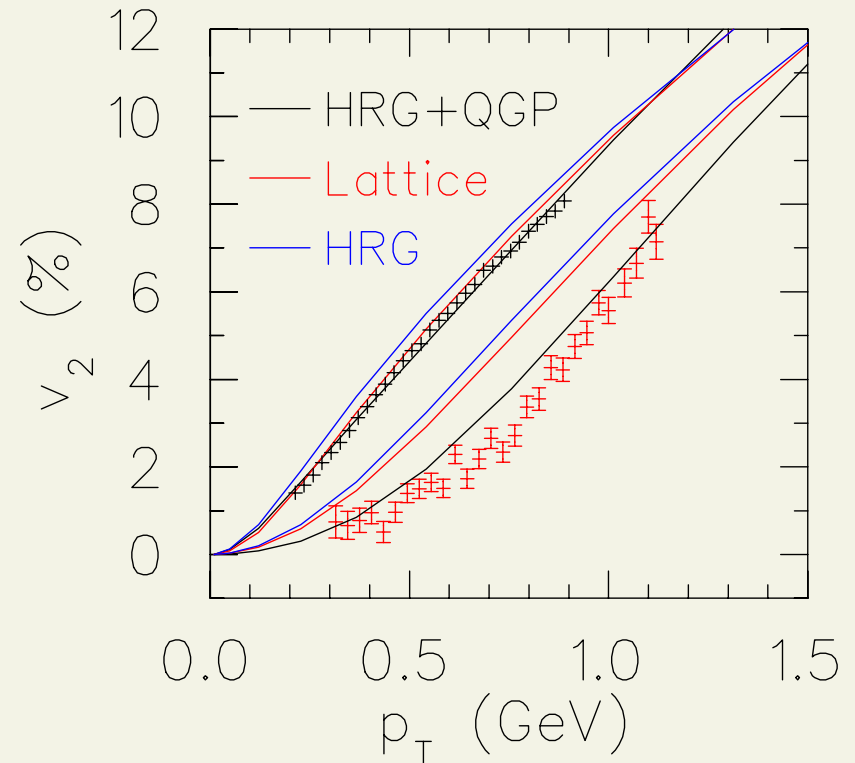
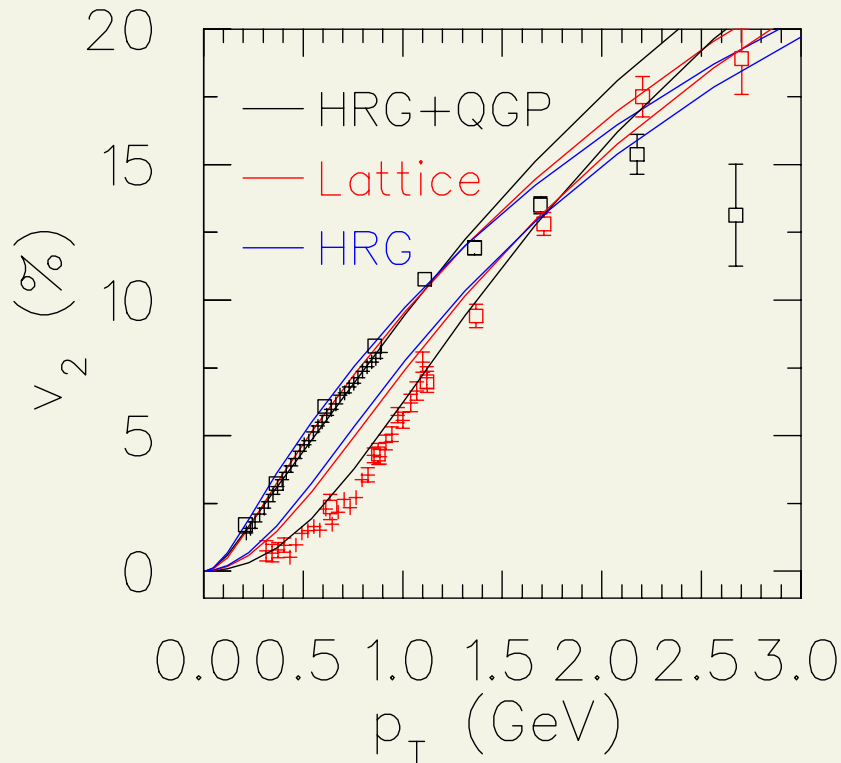
Parametrized lattice EoS



- **Extrapolation to finite quark masses?**
- **Schneider & Weise:** quasiparticle model for plasma with known masses
 - Tune the model to fit lattice results with lattice masses
 - Change masses to more physical ones

Anisotropy as function of p_T

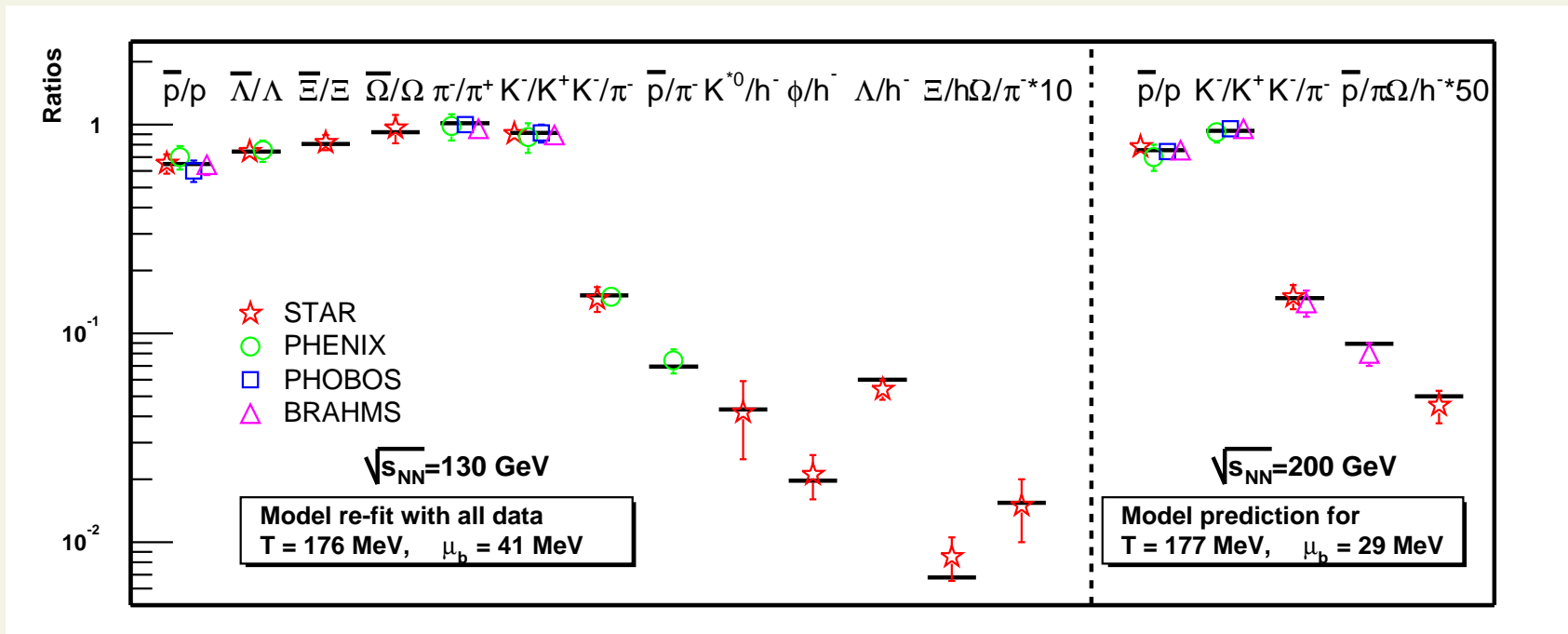
minbias Au+Au at RHIC ($\sqrt{s_{NN}} = 200$ GeV)



- All EoSs lead to **similar pion** $v_2(p_T)$
- **Proton** $v_2(p_T)$ is **sensitive** to EoS:
 - Lattice EoS gives as bad fit than EoS without any phase transition!
 - An EoS with a first order phase transition is closest to the data

Thermal models

- Hadronic phase: ideal gas of massive hadrons and resonances
- in chemical equilibrium

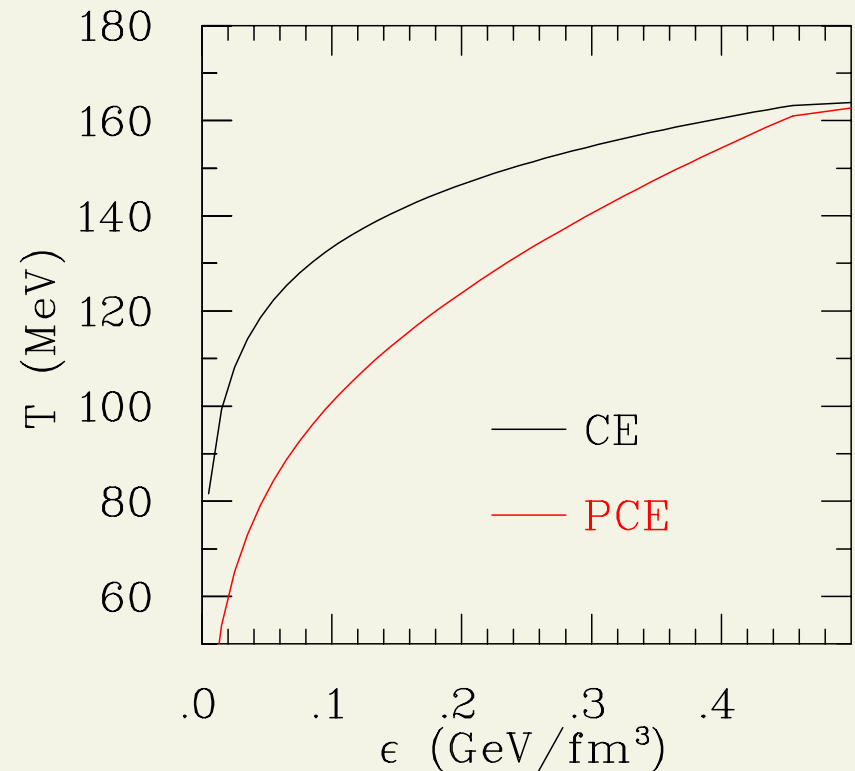
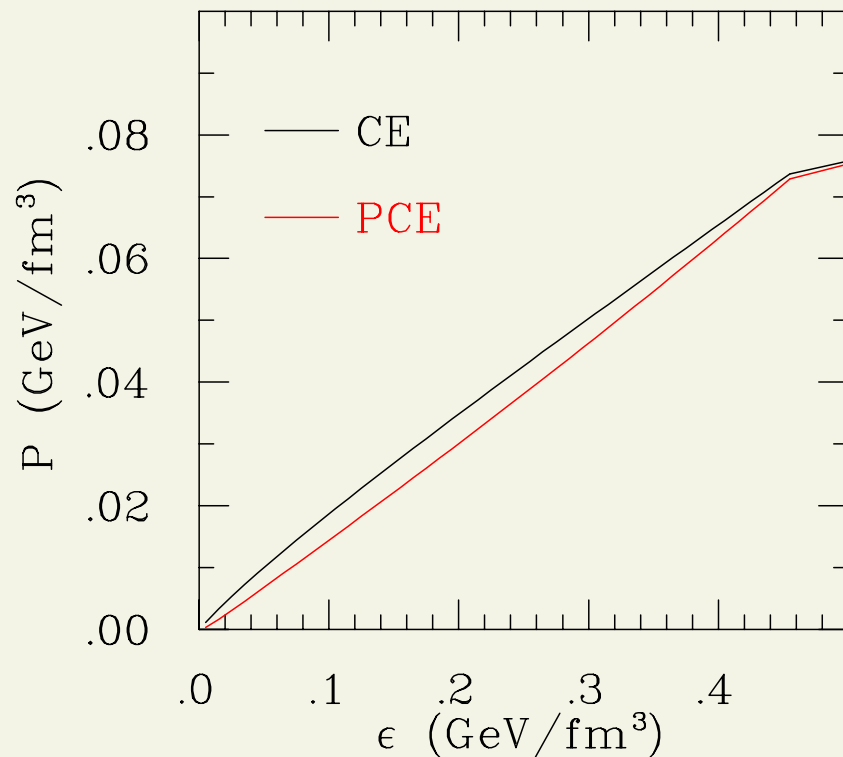


- Particle ratios (approximately) correspond to a system in $T \approx 160\text{--}170$ MeV temperature
- Evolution to $T \approx 100\text{--}120$ MeV temperature

⇒ In hydro particle ratios become wrong

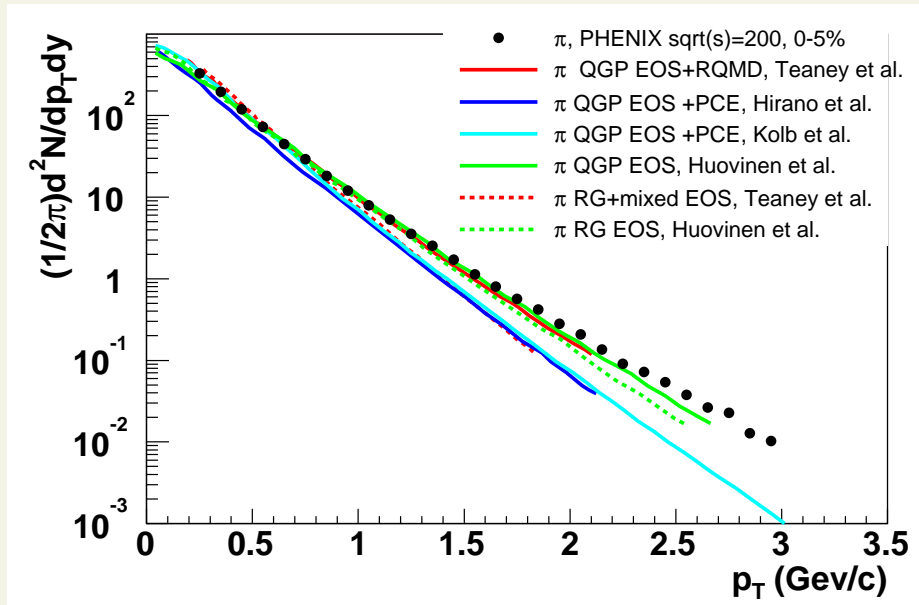
Chemical non-equilibrium

- Treat number of pions, kaons etc. as conserved quantum numbers below T_{ch} (Bebie et al, Nucl.Phys.B378:95-130,1992)
- $P = P(\epsilon, n_b)$ changes very little, but $T = T(\epsilon, n_b)$ changes. . . .

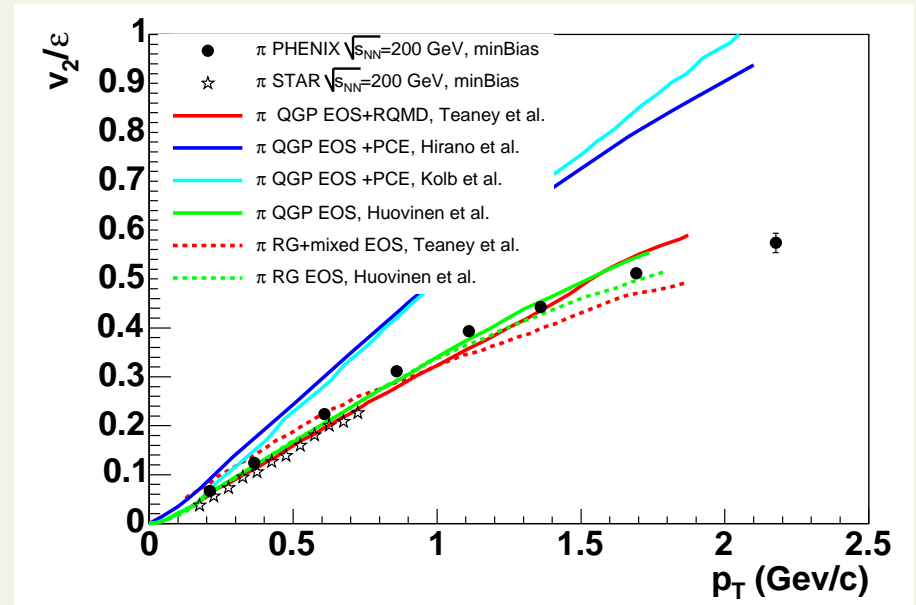


PHENIX whitepaper, NPA 757, 184 (2005)

π p_T -spectrum:



π $v_2(p_T)$:



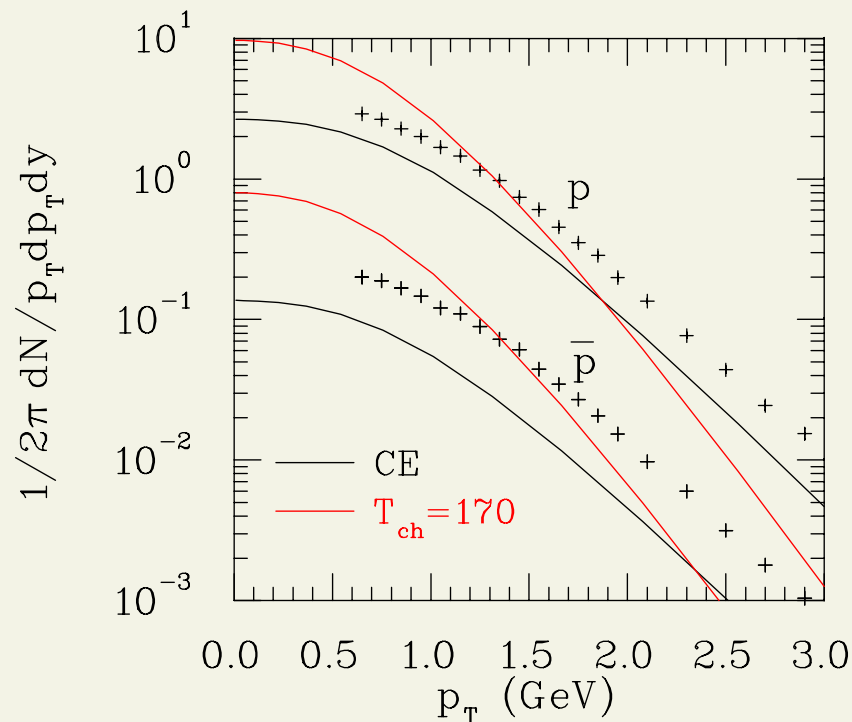
- **Blue curves:** Hirano & Tsuda and Kolb & Rapp, ideal hydro, no chemical equilibrium

⇒ **IT DOES NOT WORK!**

- If ideal hydro reproduces particle yields, it does not reproduce their distributions nor elliptic flow

Closer look at (anti)protons

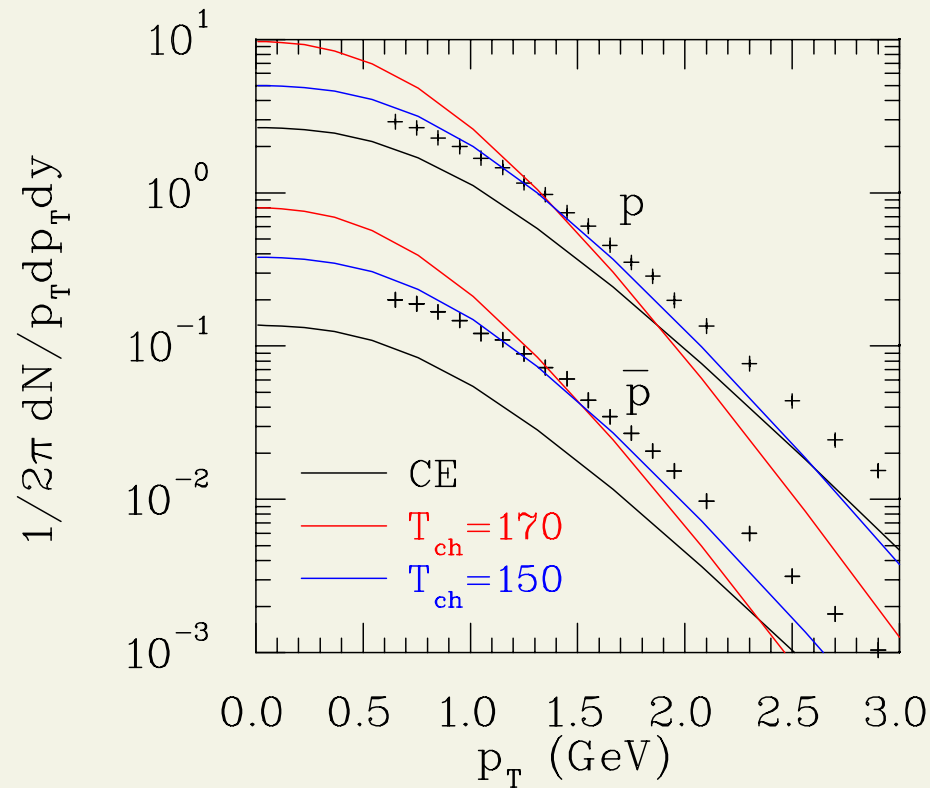
- **Too many protons and antiprotons:**



dN/dy	π^+	p	\bar{p}
PHENIX	286.4 ± 24.2	18.4 ± 2.6	13.5 ± 1.8
$T_{ch} = 170$	268	25	20

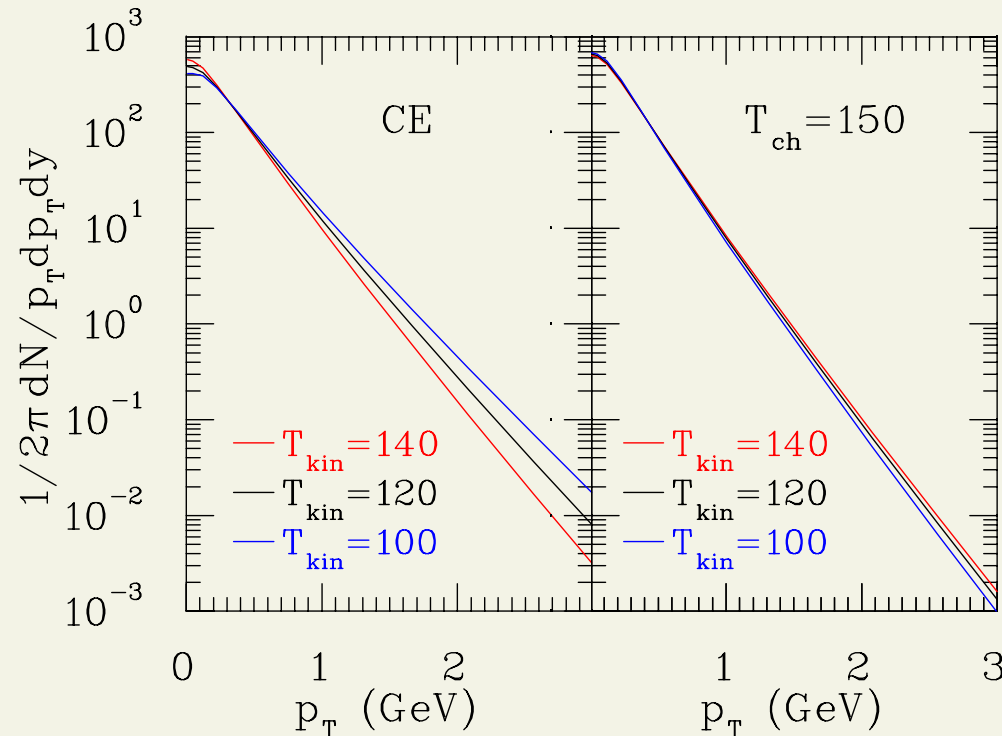
- Thermal models slightly overestimate p/π ratio
- Fit of π , K , p and \bar{p} ratios $\Rightarrow T \approx 150$ MeV
(Cleymans et al., Phys.Rev.C71:054901,2005)

Second look at (anti)protons



- Cooler chemical freeze-out temperature $T_{ch} = 150 \text{ MeV}$
⇒ yields OK.
- Slope too steep
 - Lower freeze-out temperature does not work

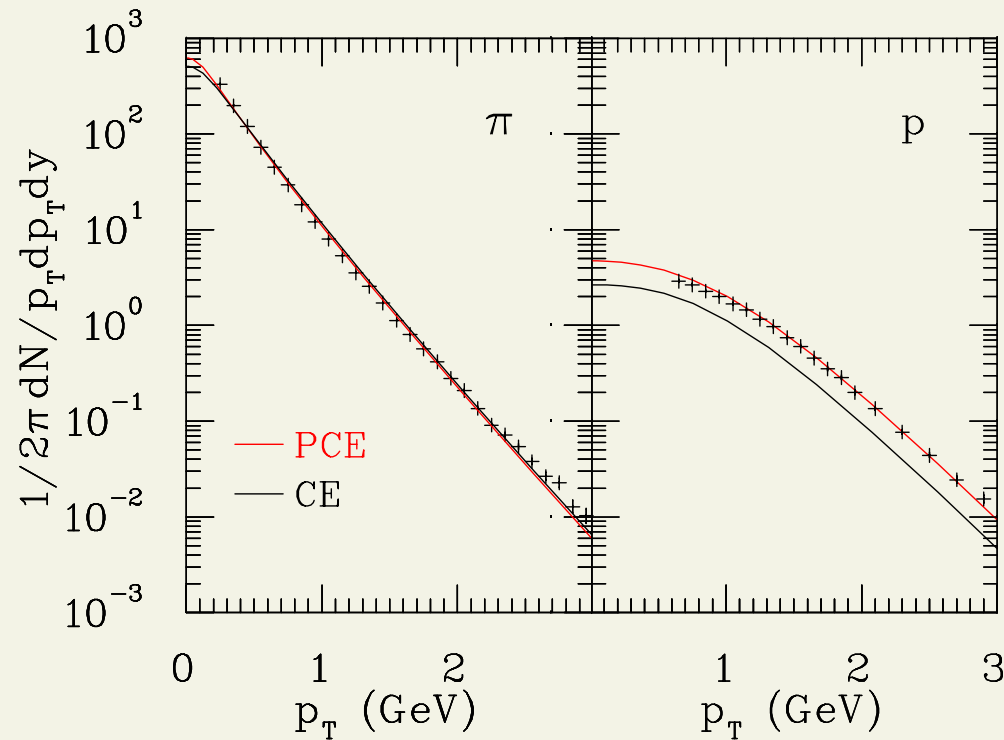
Effect of T_{kin} on pions



- Longitudinal expansion does work ($p dV$) $\Rightarrow \frac{dE_T}{dy}$ **decreases**
- If particle # is conserved, $\langle p_T \rangle$ **decreases**
- In chemical equilibrium mass energy of baryon-antibaryon pairs is converted to kinetic energy $\Rightarrow \langle p_T \rangle$ **increases!**

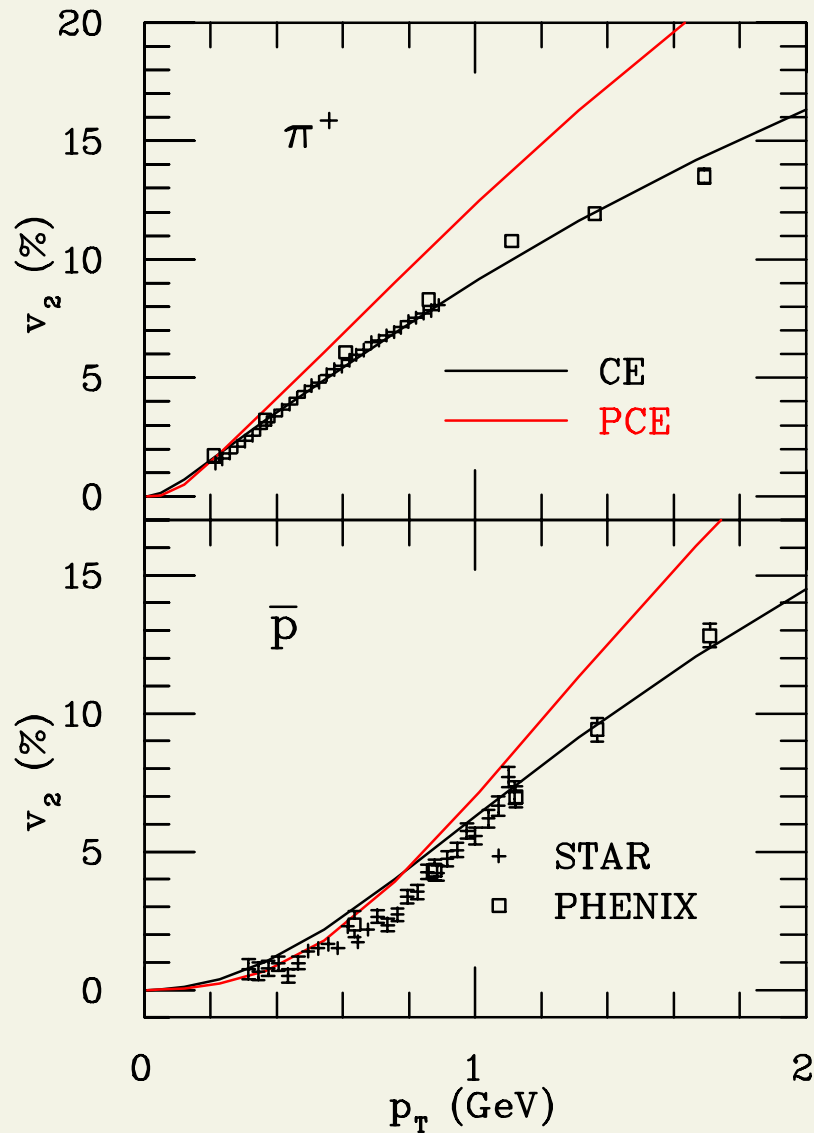
Initial state

- Need more transverse flow
 - Steeper initial density profile
 - Short initial time $\tau_0 = 0.2$ fm/c instead of $\tau_0 = 0.6-1.0$ fm/c



- If momentum distribution is isotropic, $\epsilon = 3P$ holds
- No need for exact thermalization

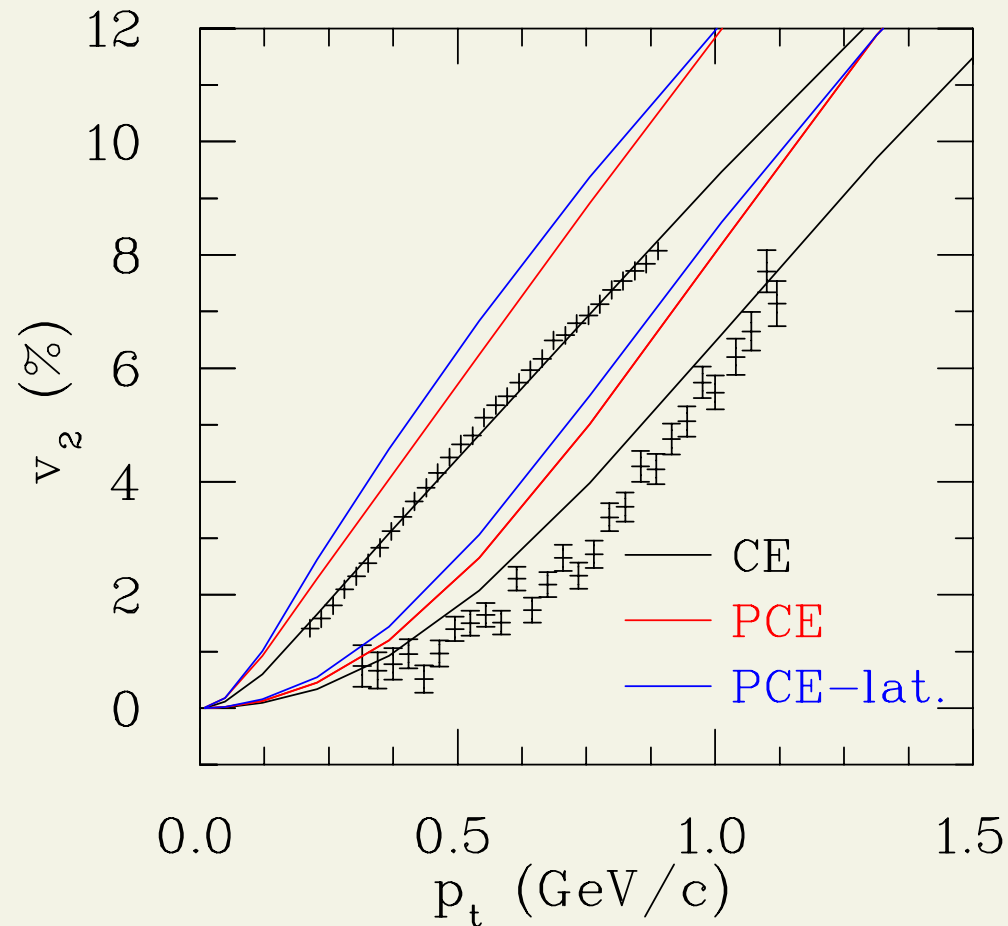
Elliptic flow $v_2(p_T)$



Failed!

- Dissipation required
- But where?
 - in hadronic phase?
 - in plasma?
 - or both?

And phase transition?



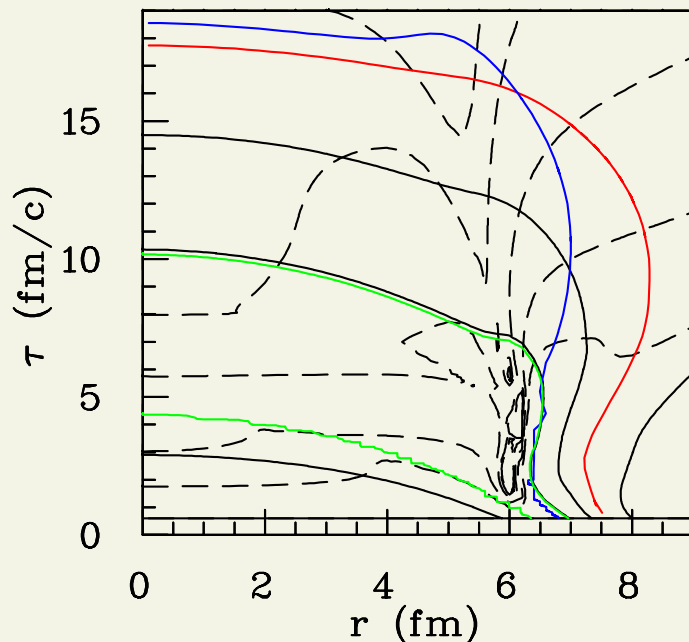
- smaller but still visible effect
- even worse fit than with a first order phase transition

Freeze-out

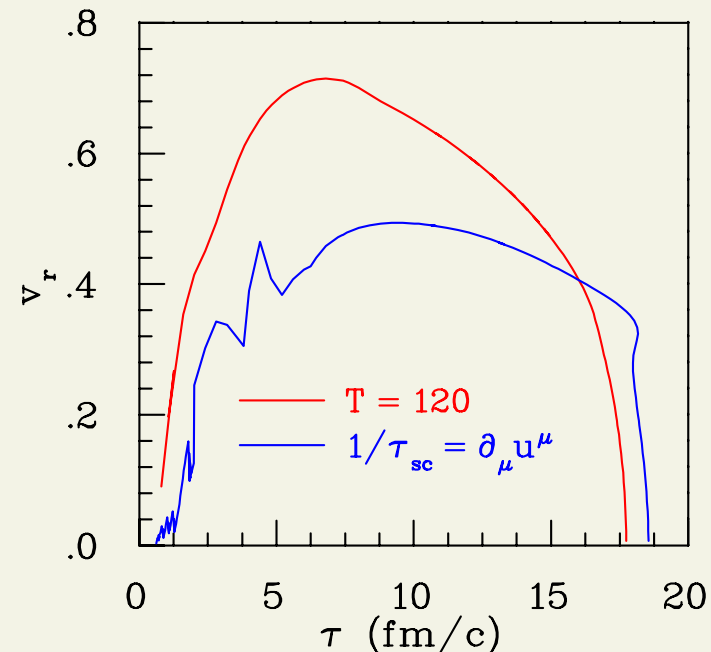
- when system no longer behaves as fluid but as free streaming particles
- momentum distributions cease to evolve \rightarrow they “freeze-out”
- **critereon:** expansion rate larger than scattering rate:

$$\tau_{\text{scat}}^{-1} \approx \partial_{\mu} u^{\mu}$$

- $\tau_{\text{scat}}^{-1} \propto T^4 \rightarrow$ freeze-out at $T = \text{const.}$
- **not so good approximation:**

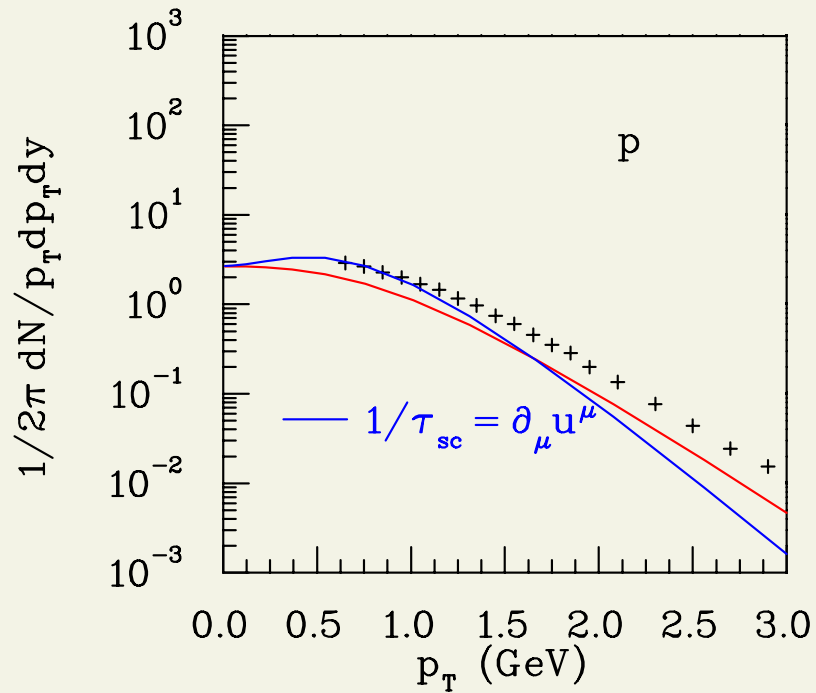
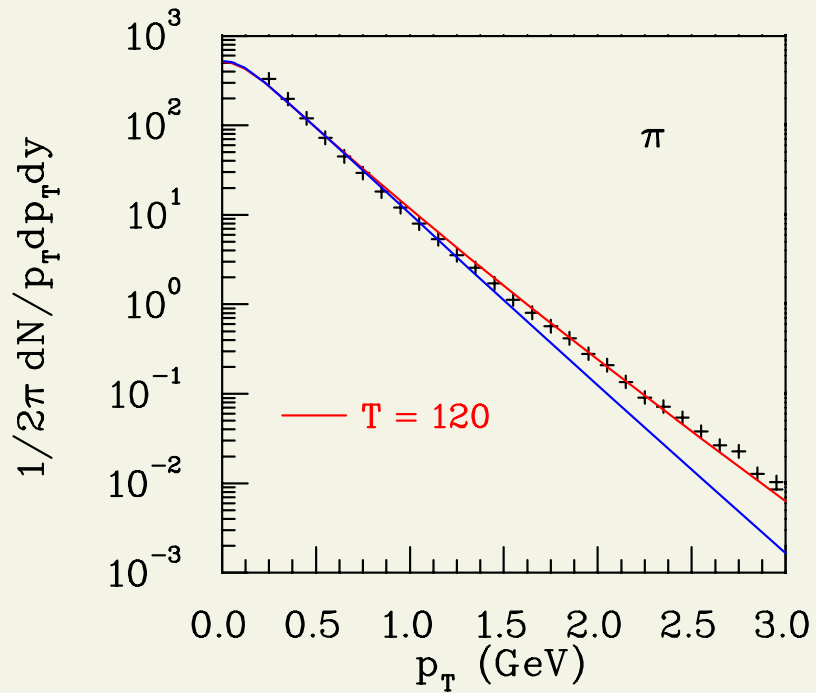


temperature and expansion rate contours



flow velocity on freeze-out surface

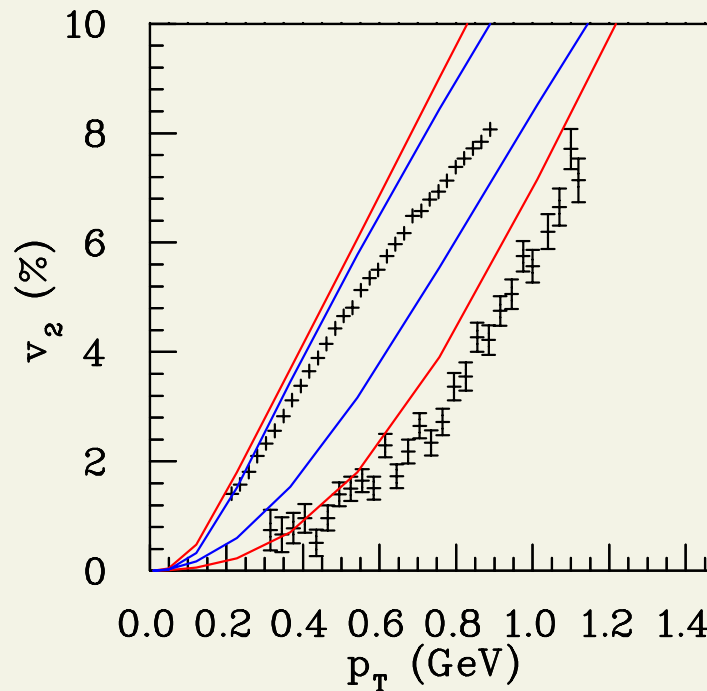
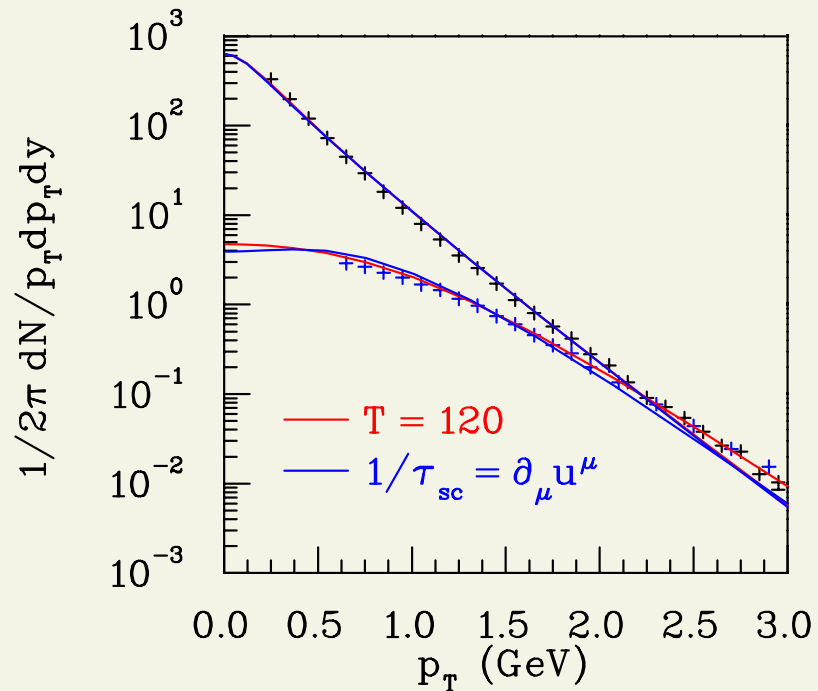
Effect on p_T spectra:



- Need more transverse flow

Early start?

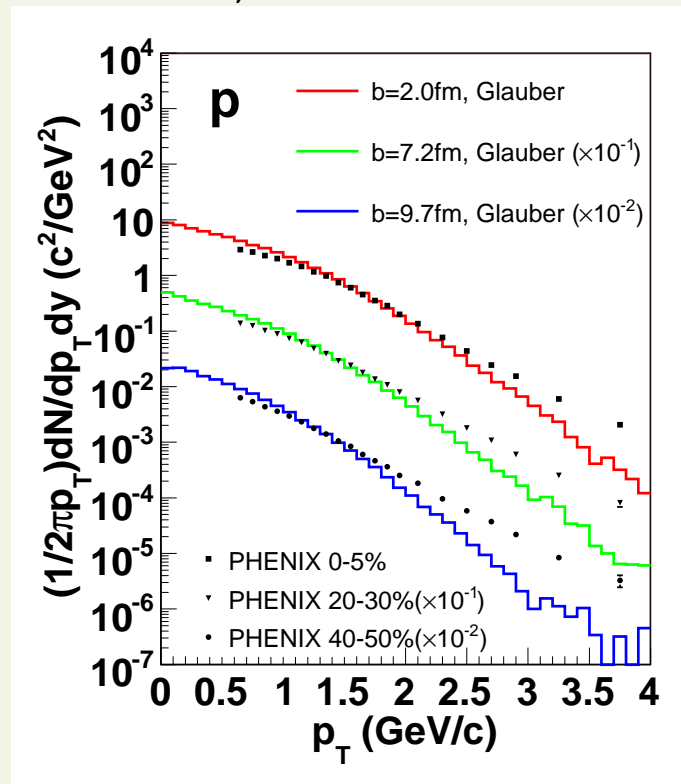
- freeze-out has only a small effect
- system freezes out soon after phase transition anyway
- results for chemical non-equilibrium and $\tau_0 = 0.2 \text{ fm}/c$:



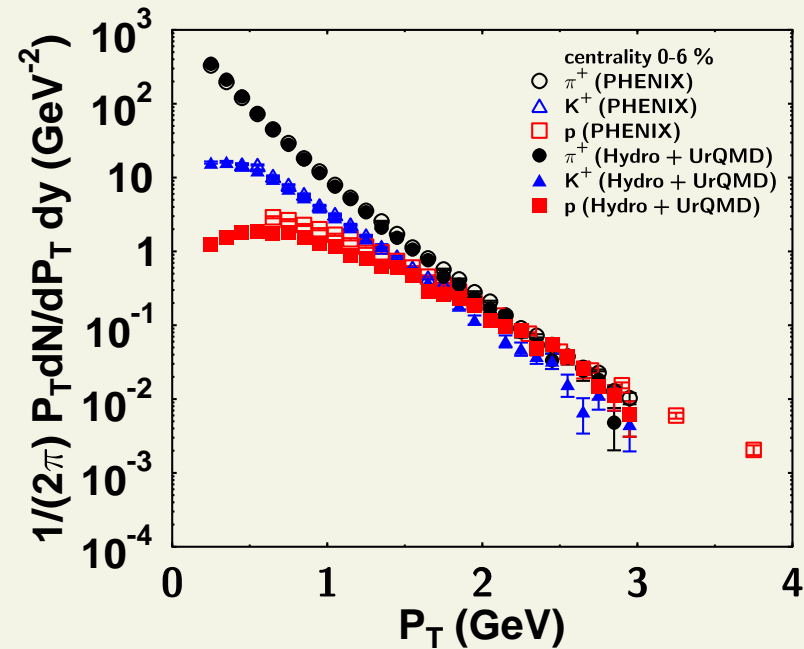
Hydro+cascade

- Cascade description describes chemistry and freeze-out
- Reproduces both spectra and v_2
- $n \rightarrow 2$ processes not included

Hirano et al., arXiv:0710.5795:



Nonaka & Bass, Phys.Rev.C75:014902,2007:



Viscous hydrodynamics

relativistic Navier-Stokes hydro: small corrections linear in gradients

$$T_{NS}^{\mu\nu} = T_{ideal}^{\mu\nu} + \eta(\nabla^\mu u^\nu + \nabla^\nu u^\mu - \frac{2}{3}\Delta^{\mu\nu}\partial^\alpha u_\alpha) + \zeta\Delta^{\mu\nu}\partial^\alpha u_\alpha$$
$$N_{NS}^\mu = N_{ideal}^\mu - \frac{n}{e+p}\kappa\nabla^\mu T$$

where $\Delta^{\mu\nu} \equiv u^\mu u^\nu - g^{\mu\nu}$, $\nabla^\mu = \Delta^{\mu\nu}\partial_\nu$

η, ζ shear and bulk viscosities, κ heat conductivity

two problems:

parabolic equations \rightarrow **acausal** Müller ('76), Israel & Stewart ('79) ...

instabilities Hiscock & Lindblom, PRD31, 725 (1985) ...

Causal viscous hydro

Müller, Israel & Stewart...

$$\Delta T^{\mu\nu} \equiv \pi^{\mu\nu} + \Pi \Delta^{\mu\nu} \quad , \quad \Delta N^\mu = -\frac{n}{e+p} q^\mu$$

bulk pressure Π , shear stress $\pi^{\mu\nu}$ heat flow q^μ treated as independent dynamical quantities that **relax** to their Navier-Stokes value on time scales $\tau_\Pi(e, n)$, $\tau_\pi(e, n)$, $\tau_q(e, n)$ - corresponds to keeping not only first but (certain) second derivatives.

Entropy four-flow including terms second order in dissipative fluxes:

$$S^\mu = su^\mu + \frac{\mu}{T} \frac{q^\mu}{h} - (\beta_0 \Pi^2 - \beta_1 q_\nu q^\nu + \beta_2 \pi_{\lambda\nu} \pi^{\lambda\nu}) \frac{u^\mu}{2T} - \frac{\alpha_0 q^\mu \Pi}{T} + \frac{\alpha_1 q_\nu \pi^{\nu\mu}}{T}$$

Require non-decrease of entropy:

$$0 \geq \partial_\mu S^\mu = \Pi X + q_\mu X^\mu + \pi_{\mu\nu} X^{\mu\nu}$$

Which terms to keep?

Muronga, Romatschke :

$$u^\alpha \partial_\alpha \pi^{\mu\nu} = -\frac{1}{\tau_\pi} \left(\pi^{\mu\nu} - 2\eta \nabla^{\langle\mu} u^{\nu\rangle} \right) \\ - (u^\mu \pi^{\nu\alpha} + u^\nu \pi^{\mu\alpha}) u^\lambda \partial_\lambda u_\alpha \\ - \frac{1}{2} \pi^{\mu\nu} \left(u^\lambda \partial_\lambda \ln \frac{\beta_2}{T} + \partial_\lambda u^\lambda \right)$$

Which terms to keep?

Muronga, Romatschke :

$$u^\alpha \partial_\alpha \pi^{\mu\nu} = -\frac{1}{\tau_\pi} \left(\pi^{\mu\nu} - 2\eta \nabla^{\langle\mu} u^{\nu\rangle} \right) - (u^\mu \pi^{\nu\alpha} + u^\nu \pi^{\mu\alpha}) u^\lambda \partial_\lambda u_\alpha - \frac{1}{2} \pi^{\mu\nu} \left(u^\lambda \partial_\lambda \ln \frac{\beta_2}{T} + \partial_\lambda u^\lambda \right)$$

Song & Heinz :

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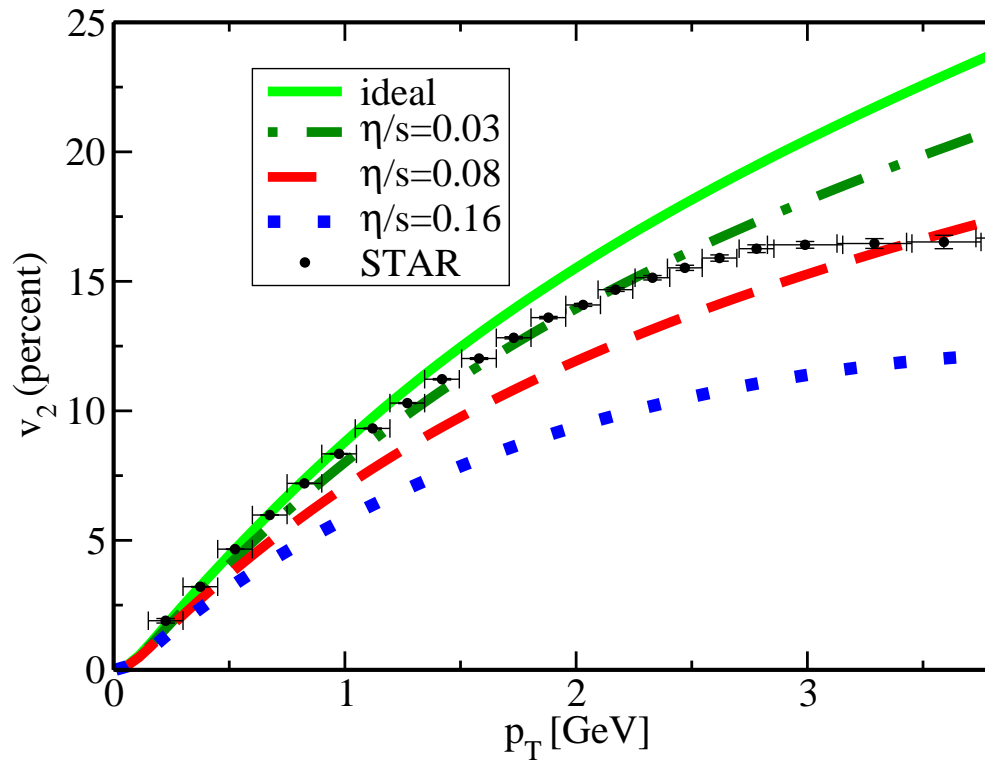
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Chaudhuri :
$$u^\alpha \partial_\alpha \pi^{\mu\nu} = -\frac{1}{\tau_\pi} \left(\pi^{\mu\nu} - 2\eta \nabla^{\langle\mu} u^{\nu\rangle} \right)$$

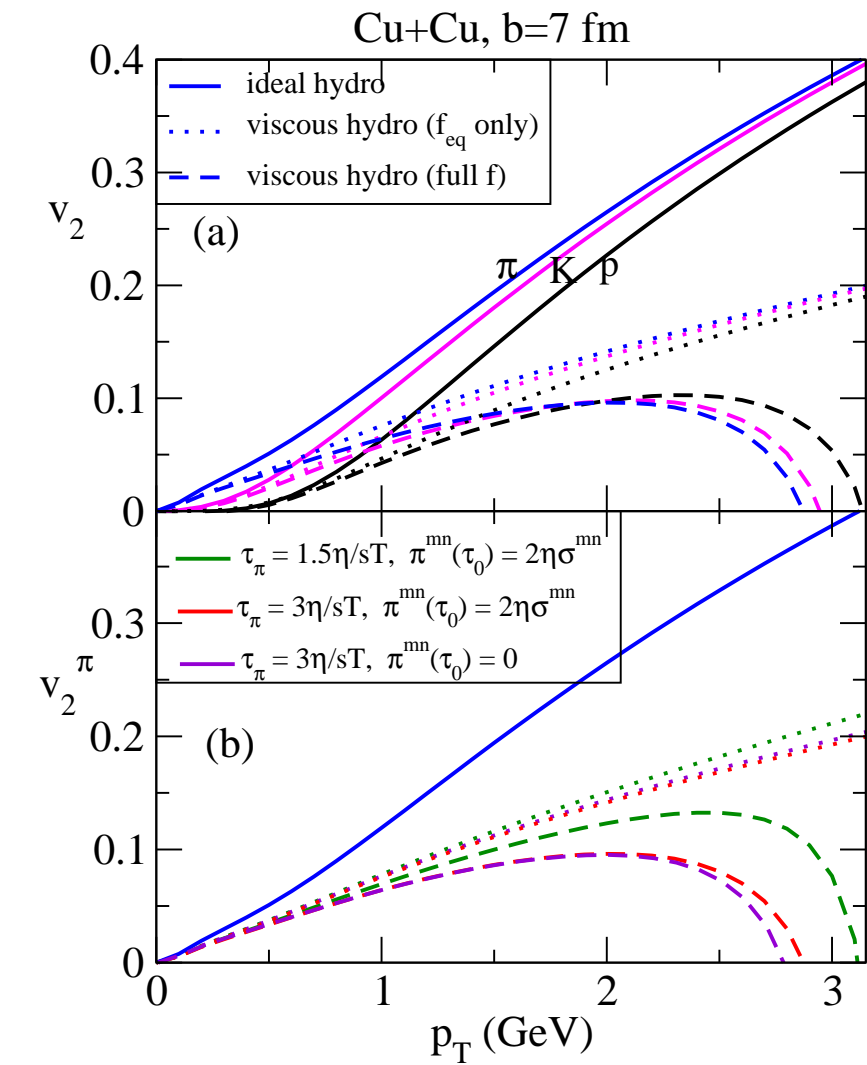
$\eta/s = 1/(4\pi)$ corrections from IS hydro



Romatschke & Romatschke, arxiv:0706.1522

10-20%

or



Song & Heinz, arxiv:0709.0742

50+% reduction??

Viscous hydro vs transport

We solve the full Israel-Stewart-Muronga equations, including vorticity terms from kinetic theory, in a 2+1D boost-invariant scenario. Shear stress only.

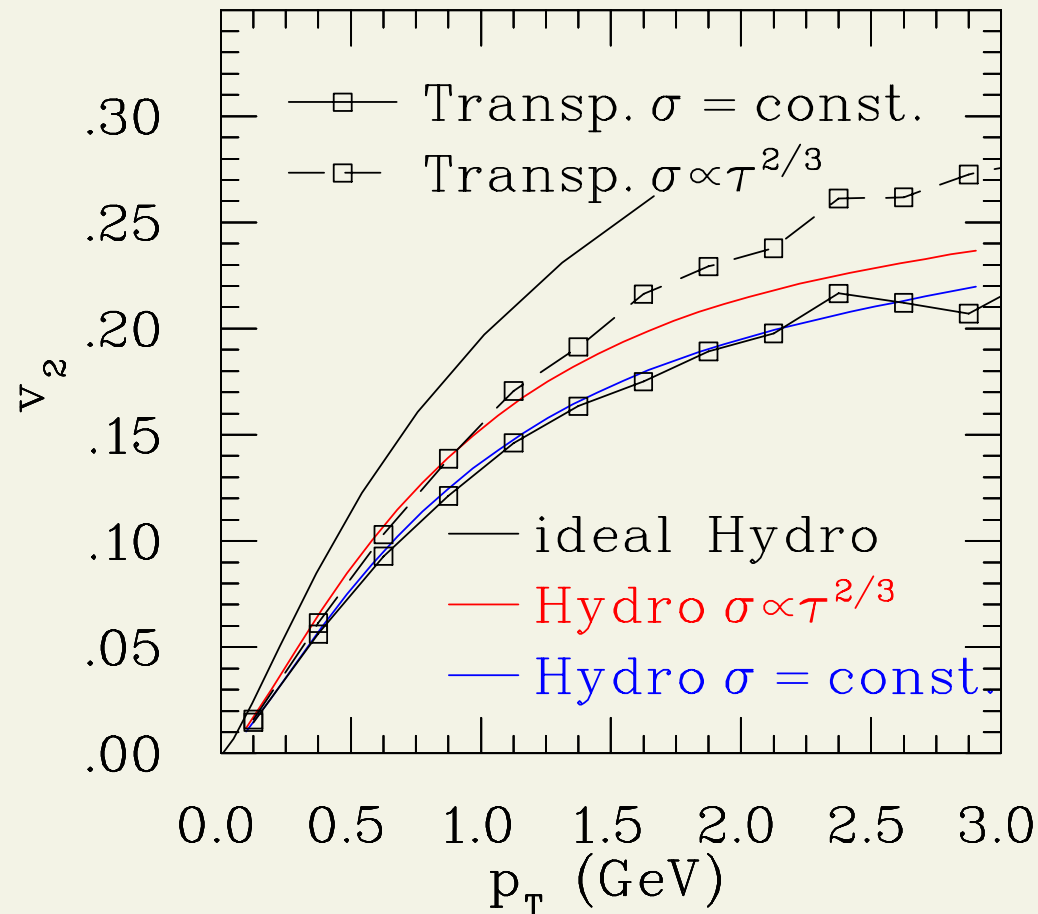
Mimic a known reliable transport model:

- massless Boltzmann particles $\Rightarrow \epsilon = 3P$
- only $2 \leftrightarrow 2$ processes, i.e. conserved particle number
- $\eta = 4T/(5\sigma_{\text{tr}})$
- either $\sigma_{\text{tr}} = \text{const.}$ the simplest in transport
or $\sigma_{\text{tr}} \propto \tau^{2/3}$ close to $\eta/s = 1/(4\pi)$

Our “RHIC-like” initialization:

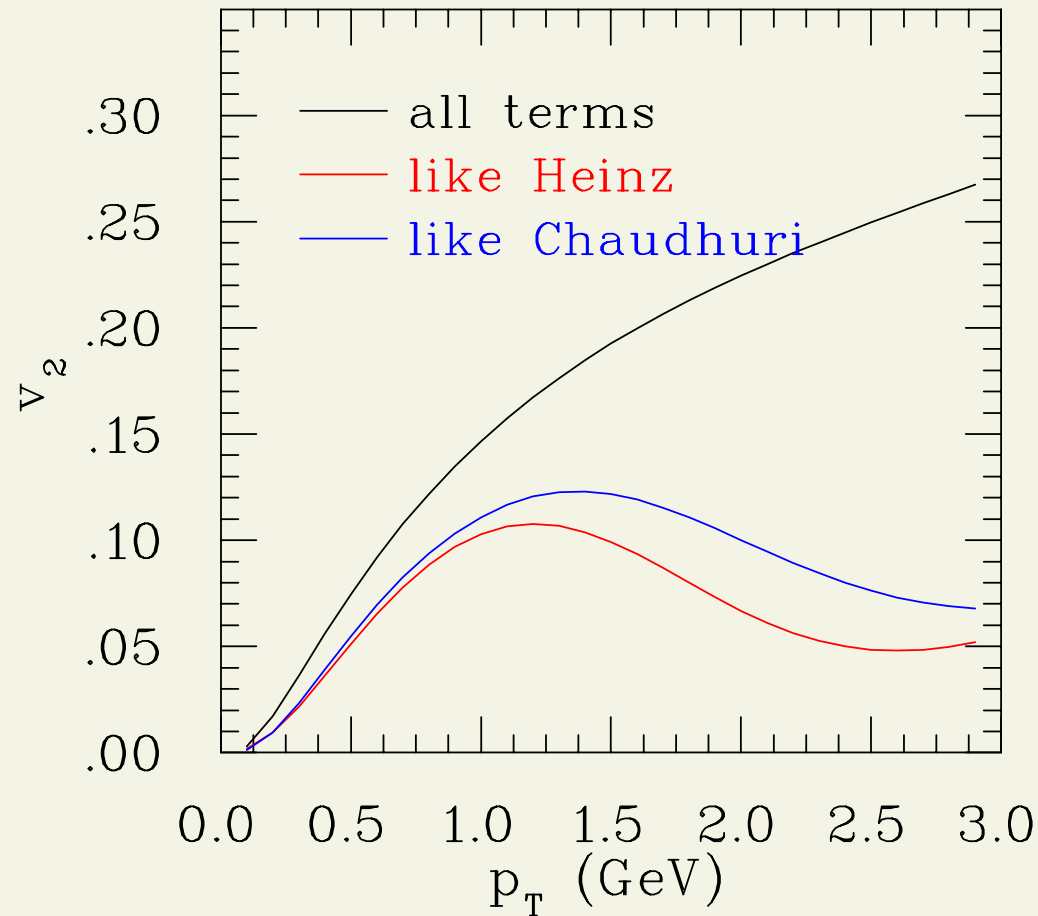
- $\tau_0 = 0.6 \text{ fm}/c$
- $b = 8 \text{ fm}$
- $T_0 = 385 \text{ MeV}$ and $dN/d\eta|_{b=0} = 1000$
- freeze-out at constant $n = 0.365 \text{ fm}^{-3}$

Viscous hydro vs transport v_2



- excellent agreement when $\sigma = \text{const} \sim 47 \text{ mb}$
- good agreement for $\eta/s \approx 1/(4\pi)$, i.e. $\sigma \propto \tau^{2/3}$
- BUT results sensitive to freeze-out criterion, especially at high p_T

Viscous hydro v_2



- **Important to keep all terms!**

Conclusions

improvements to ideal hydrodynamical models

- crossover phase transition
- correct particle ratios
- more physical freeze-out

increase elliptic flow

dissipative corrections required to fit the data

how large viscosity allowed in plasma unknown

prospects for applicability of Israel-Stewart causal hydrodynamics at RHIC look promising, based on comparisons with covariant transport in 2+1D Bjorken scenario

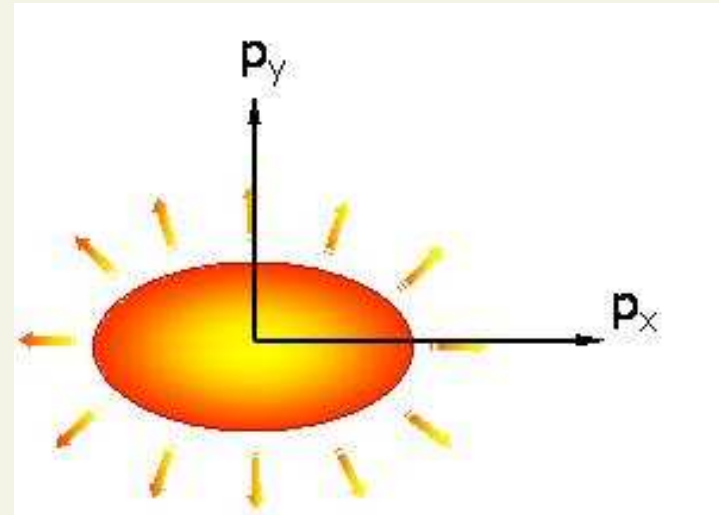
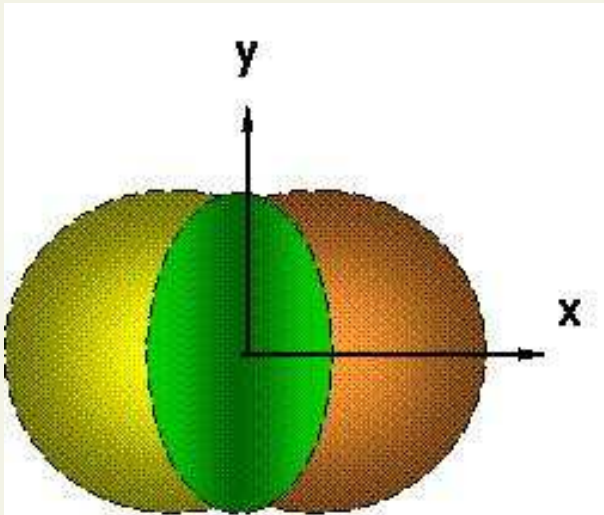
we expect 20 – 30% reduction to $v_2(p_T)$ due to dissipation

Elliptic flow v_2

spatial anisotropy



final azimuthal momentum anisotropy



$$\varepsilon \equiv \frac{\langle x^2 - y^2 \rangle}{\langle x^2 + y^2 \rangle}$$

$$v_2 \equiv \frac{\langle p_x^2 - p_y^2 \rangle}{\langle p_x^2 + p_y^2 \rangle}$$

sensitive to speed of sound $c_s^2 = \partial p / \partial e$ and shear viscosity η