



# Simulations of dynamical friction including spatially-varying magnetic fields

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**COOL 05 Workshop**  
**Galena, IL**  
**September 21, 2005**

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Work supported by US DOE,  
Office of NP, SBIR program,  
contract # **DE-FG03-01ER83313**



## Motivation

- Parameters for RHIC II cooler are unprecedented
  - see I. Ben-Zvi et al., Proc. COOL 03 Workshop (2003).
  - friction forces must be understood to within a factor of  $\sim 2$
- There's a need for high-fidelity simulations
  - We are using the VORPAL code
    - C. Nieter and J.R. Cary, Journal of Computational Physics (2004)
    - <http://www-beams.colorado.edu/vorpal/>
- Goals of the simulations
  - Resolve differences in theory, asymptotics, parametric models
    - Understand magnetization in limit of small Coulomb logarithm
  - Quantify the effect of magnetic field errors
- Numerical approach:
  - use  $O(N^2)$  algorithm from astrophysical dynamics community
    - 4<sup>th</sup>-order predictor-corrector with aggressive variation of time step
    - accurately resolves close binary collisions



## 4<sup>th</sup>-Order Predictor/Corrector Hermite Algorithm

- Algorithm developed and used extensively by galactic dynamics community
  - J. Makino, The Astrophysical Journal **369**, 200 (1991)
  - J. Makino & S. Aarseth, Publ. Astron. Soc. Japan **44**, 141 (1992)

- Predictor step:

$$\mathbf{v}_{p,j} = \frac{1}{2} (t - t_j) \dot{\mathbf{a}}_j + (t - t_j) \mathbf{a}_j + \mathbf{v}_j$$

$$\mathbf{x}_{p,j} = \frac{1}{6} (t - t_j) \dot{\mathbf{a}}_j + \frac{1}{2} (t - t_j) \mathbf{a}_j + (t - t_j) \mathbf{v}_j + \mathbf{x}_j$$

- where

$$\mathbf{a}_i = \frac{q_i}{m_i} \mathbf{v}_i \times \mathbf{B} + \frac{q_i}{4\pi\epsilon_0 m_i} \sum_j \frac{q_j \mathbf{r}_{ij}}{(r_{ij}^2 + r_c^2)^{3/2}} \quad \mathbf{r}_{ij} = \mathbf{x}_{p,j} - \mathbf{x}_{p,i}$$

$$\dot{\mathbf{a}}_i = \frac{q_i}{m_i} \mathbf{a}_i \times \mathbf{B} + \frac{q_i}{4\pi\epsilon_0 m_i} \sum_j q_j \left[ \frac{\mathbf{v}_{ij}}{(r_{ij}^2 + r_c^2)^{3/2}} + \frac{3(\mathbf{v}_{ij} \cdot \mathbf{r}_{ij}) \mathbf{r}_{ij}}{(r_{ij}^2 + r_c^2)^{5/2}} \right] \quad \mathbf{v}_{ij} = \mathbf{v}_{p,j} - \mathbf{v}_{p,i}$$

$r_c \gg 0$  “cloud” radius



## Hermite Algorithm – including a Magnetic Field

- The corrector step:
 
$$\mathbf{x}_i(t_i + \Delta t_i) = \mathbf{x}_{p,i} + \frac{1}{24} \Delta t_i^4 \mathbf{a}_{0,i}^{(2)} + \frac{1}{120} \Delta t_i^5 \mathbf{a}_{0,i}^{(3)}$$

$$\mathbf{v}_i(t_i + \Delta t_i) = \mathbf{v}_{p,i} + \frac{1}{6} \Delta t_i^3 \mathbf{a}_{0,i}^{(2)} + \frac{1}{24} \Delta t_i^4 \mathbf{a}_{0,i}^{(3)}$$
  - where  $\mathbf{a}_{0,i}^{(2)}$  and  $\mathbf{a}_{0,i}^{(3)}$  are linear functions of  $\mathbf{a}(t)$  and  $\dot{\mathbf{a}}(t)$  evaluated at times  $t_i$  and  $t_i + \Delta t_i$
  
- Adding B-field breaks 4<sup>th</sup>-order scaling, unless
  - Lab-frame B is purely longitudinal, constant in time
  - $\mathbf{v} \times \mathbf{B}$  force is evaluated again at the predicted positions
  - magnetic term in velocity correction (far right term above):
    - is split into self-field & magnetic terms
    - the coefficient in front of  $\mathbf{a}_{self-field}^{(3)}$  is changed from  $\frac{1}{24}$  to  $\frac{5}{72}$   $\mathbf{a}_{magnetic,i}^{(3)}$



## Initial Study of Magnetic Field Errors – Motivation

- The effect of magnetic field errors in a solenoid on the dynamical velocity drag (i.e. friction) of an ion in an electron cooler is not well understood
  - The parametric model of Parkhomchuk treats field errors as an effective transverse rms velocity of the electron Larmor circles
    - Contribution appears in same place as  $V_{e,rms,\parallel}$
    - In the absence of an explicit model, field errors have been treated as an effective increase in  $V_{e,rms,\parallel}$
- Our primary interest is the cooler for RHIC II
  - We consider the CELSIUS ring here, to take advantage of recent experiments
  - We consider two very different models for the errors



## Magnetic field errors – “Model 1”

- A sum of sinusoidal terms (lab frame)

$$B_x = \sum_i b_i \frac{k_{x,i}}{k_{z,i}} \exp(k_{x,i}x) \exp(k_{y,i}y) \sin(k_{z,i}z + \varphi_{z,i})$$

$$B_y = \sum_i b_i \frac{k_{y,i}}{k_{z,i}} \exp(k_{x,i}x) \exp(k_{y,i}y) \sin(k_{z,i}z + \varphi_{z,i})$$

$$B_z = B_0 + \sum_i b_i \exp(k_{x,i}x) \exp(k_{y,i}y) \cos(k_{z,i}z + \varphi_{z,i})$$

$$k_{z,i}^2 = k_{x,i}^2 + k_{y,i}^2 \quad \lambda_i = 2\pi / k_{z,i}$$

- a more general form of the equations is allowed
- we assume  $b_i \ll B_0$  for all  $i$
- appropriate choices for  $b_i, \lambda_i$ , etc. are not yet clear
- here, we consider a single component



## Magnetic field errors – “Model 2”

- A sum of piece-wise constant “tilts” (lab frame)

$$B_x = \sum_i b_{x,i} H(z - z_{x,i}) H(z_{x,i+1} - z)$$

$$B_y = \sum_i b_{y,i} H(z - z_{y,i}) H(z_{y,i+1} - z)$$

$$B_z = B_0$$

- $H(x)$  is the unit Heaviside function
- we assume  $b_{ix}, b_{iy} \ll B_0$  for all  $i$
- small abuse of Maxwell’s eqn.’s at discontinuities
- parameters taken from design report
  - M. Sedlacek et al., “Design and Construction of the CELSIUS Electron Cooler,” [http://preprints.cern.ch/cgi-bin/setlink?base=cernrep&categ=Yellow\\_Report&id=94-03](http://preprints.cern.ch/cgi-bin/setlink?base=cernrep&categ=Yellow_Report&id=94-03)
  - amplitude of tilts (highly variable) is  $\sim 1.e-03$
  - length of segments (highly variable) is  $\sim 20$  cm



## Fields are Lorentz-transformed to beam frame

- VORPAL cooling sim.'s are in the beam frame

$$B_z' = B_z(x', y', \gamma\beta ct')$$

$$B_x' = \gamma B_x(x', y', \gamma\beta ct') \quad B_y' = \gamma B_y(x', y', \gamma\beta ct')$$

$$E_x' = -\beta c B_y' \quad E_y' = \beta c B_x' \quad E_z' = 0$$

- **E** fields are dominant for “relativistic” coolers
  - because electrons are non-relativistic in the beam frame
  - not strictly true for CELSIUS, for which  $\beta \sim 0.3$





## Basic Parameter set with 5 variations

Symbol	Meaning	Value	Units
$B_0$	solenoid field	0.1	T
$L_{\text{sol}}$	solenoid length	2.5	m
$\beta$	proton bunch velocity / c	0.308	
$\tau_{\text{lab}}$	interaction time (lab frame)	$2.7 \times 10^{-8}$	s
$\tau_{\text{beam}}$	interact. time (beam frame)	$2.6 \times 10^{-8}$	s
$\Delta t$	largest time step	$2.6 \times 10^{-12}$	s
$dt_{\text{min}}$	smallest time step	$8.0 \times 10^{-14}$	s
$\omega_{\text{pe}}$	e- plasma frequency	$4.1 \times 10^8$	rad/s
$\Omega_L$	e- Larmor frequency	$1.8 \times 10^{10}$	rad/s
$r_L$	e- Larmor (gyro-) radius	$7.9 \times 10^{-6}$	m
$L_{x,y,z}$	sim. domain dimensions	$6.0 \times 10^{-4}$	m
$n_e$	e- number density	$5.4 \times 10^{13}$	$\text{m}^{-3}$
$N_e$	# of simulated e-'s	$1.2 \times 10^3$	
$\Delta_{e,\perp}$	transverse rms e- velocity	$1.4 \times 10^5$	m/s
$\Delta_{e,\parallel}$	long. rms e- velocity	$3.0 \times 10^3$	m/s
$\Delta_{\text{eff},\parallel}$	effective long. rms e- vel.	$9.0 \times 10^3$	m/s

We consider 5 separate cases –  
2 with field errors & 3 without

$$V_{\text{ion},\parallel} = 10,000 \text{ m/s}$$

$$V_{\text{ion},\perp} = 10,000 \text{ m/s}$$

“cld” –  $\Delta_{e,\parallel} = 3,000$  (no errors)

“wrm” –  $\Delta_{e,\parallel} = 9,000$  (no errors)

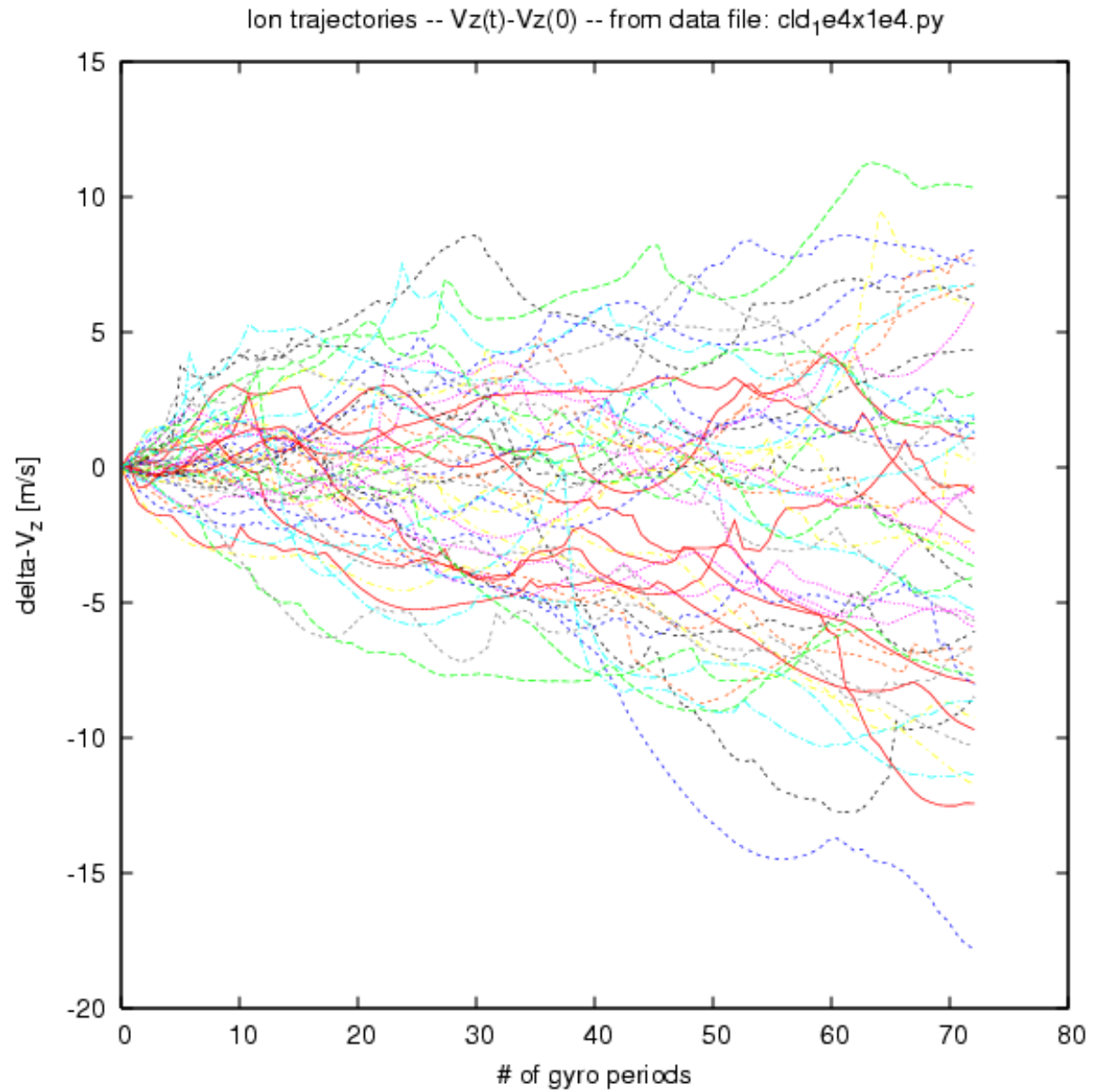
“hot” –  $\Delta_{e,\parallel} = 18,000$  (no errors)

“sin” –  $\Delta_{e,\parallel} = 3,000$  (Model 1)

“err” –  $\Delta_{e,\parallel} = 3,000$  (Model 2)



“cld” parameters –  $\Delta_{rms,\parallel} = 3000$  (no field errors)



Simulating dynamical

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p.

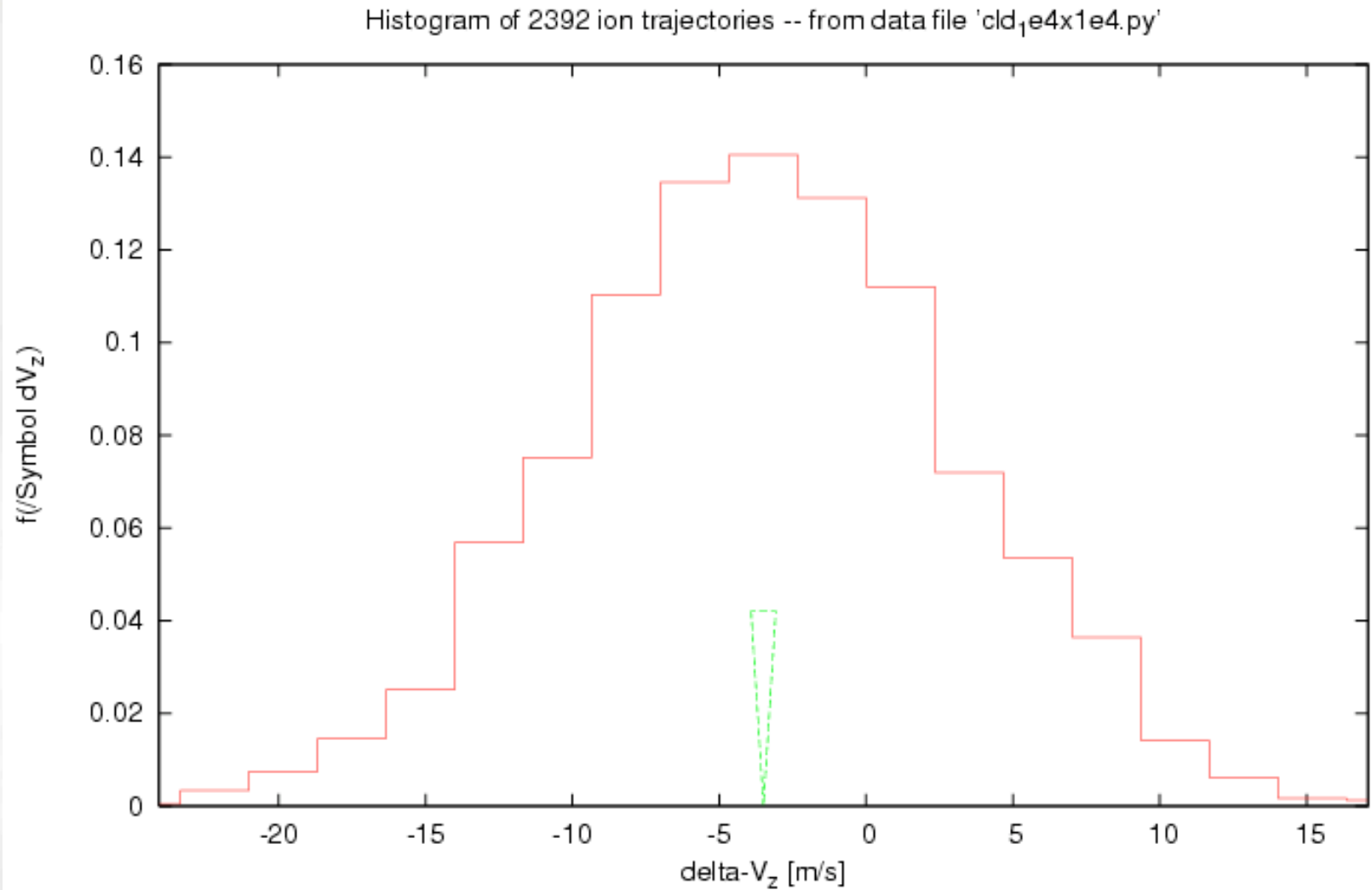


## Diffusive dynamics can obscure friction/drag

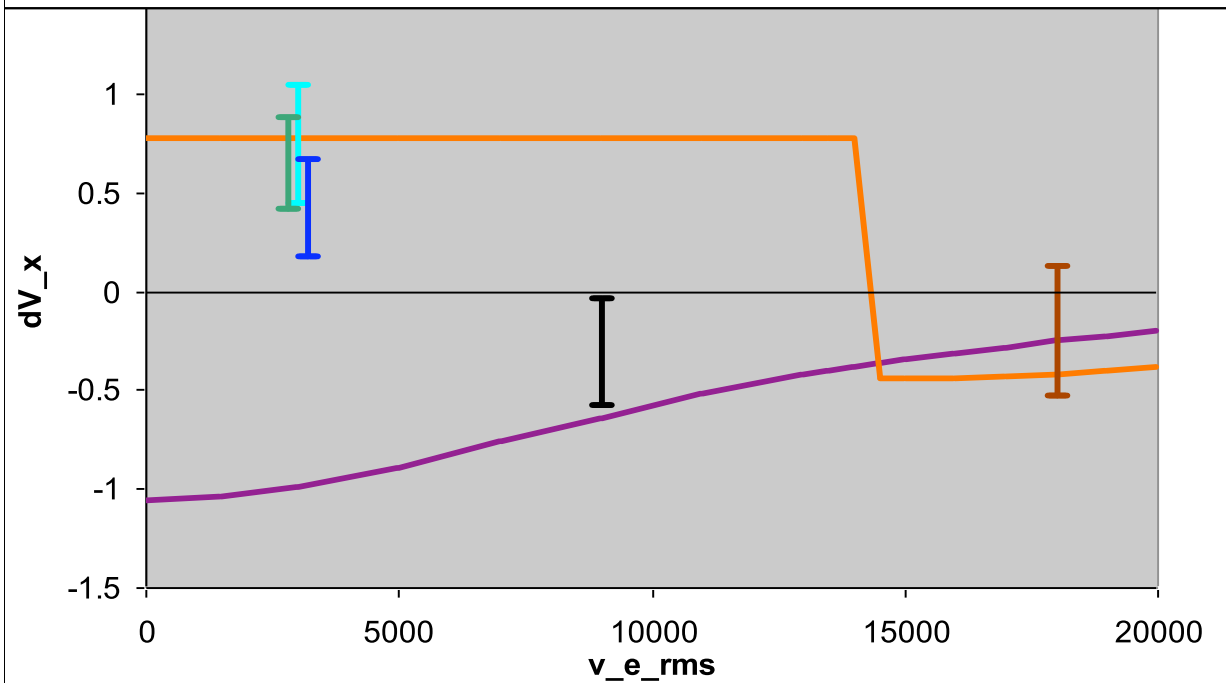
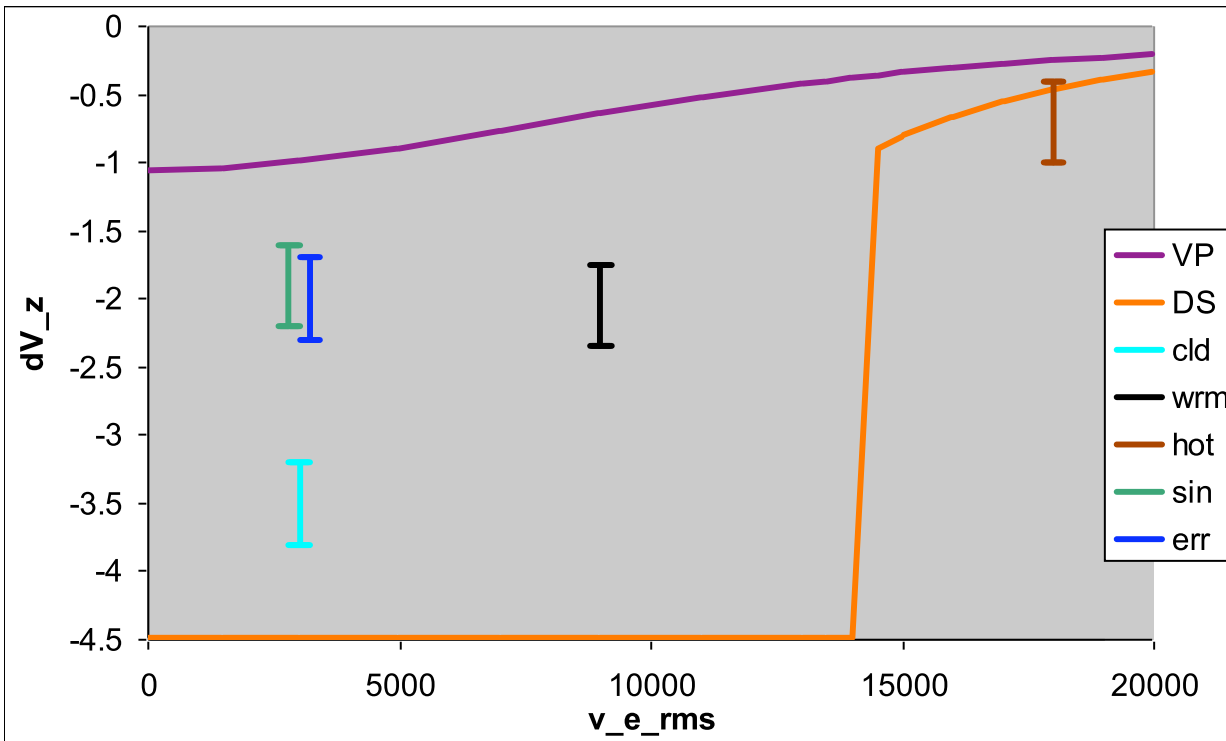
- For a single pass through the cooler
  - Diffusive velocity kicks are larger than velocity drag
  - Consistent with theory
- For sufficiently large  $\Delta_{e,\parallel}$ 
  - numerical trick of e-/e+ pairs can suppress diffusion
  - not valid for CELSIUS parameters
- Only remaining tactic is to generate 100's of trajectories
  - Central Limit Theorem states that mean velocity drag is drawn from a Gaussian distribution, with rms reduced by  $N_{\text{traj}}^{1/2}$  as compared to the rms spread of the original distribution
  - Hence, error bars are +/- 1 rms /  $N_{\text{traj}}^{1/2}$
- Not practical to routinely generate 100's or 1000's of trajectories "by hand"
  - run 8 trajectories simultaneously
  - use "task farming" approach to automate many runs



# “cld” parameters – $\Delta_{\text{rms},\parallel} = 3000$ (no field errors)



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## Error models yield similar results –

- Longitudinal velocity drag is significantly reduced
  - in agreement with parametric increase of  $\Delta_{rms,\parallel}$
- Transverse velocity drag is less affected
  - NOT consistent with parametric increase of  $\Delta_{rms,\perp}$



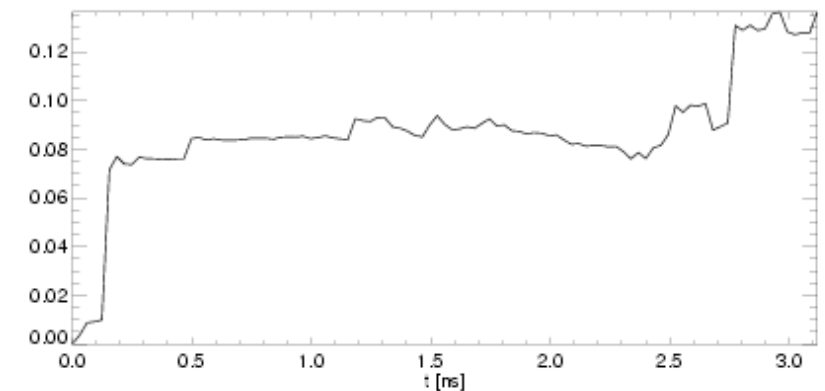
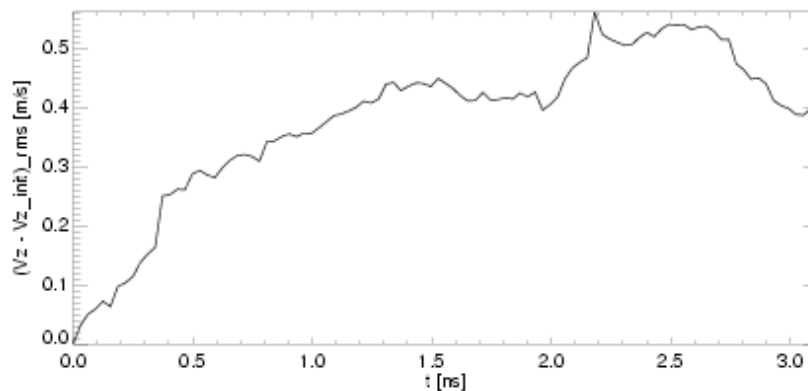
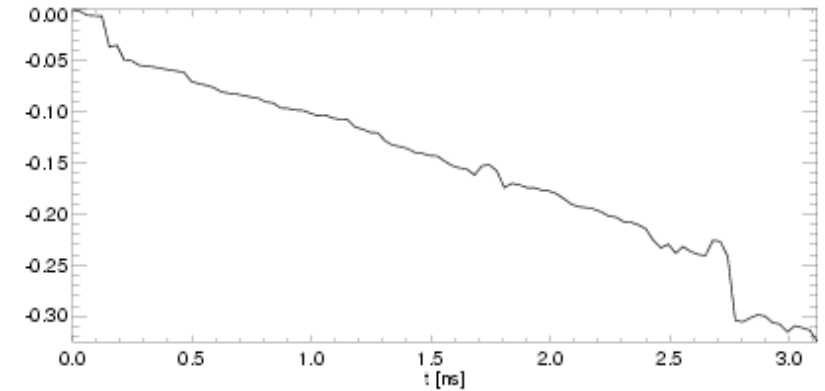
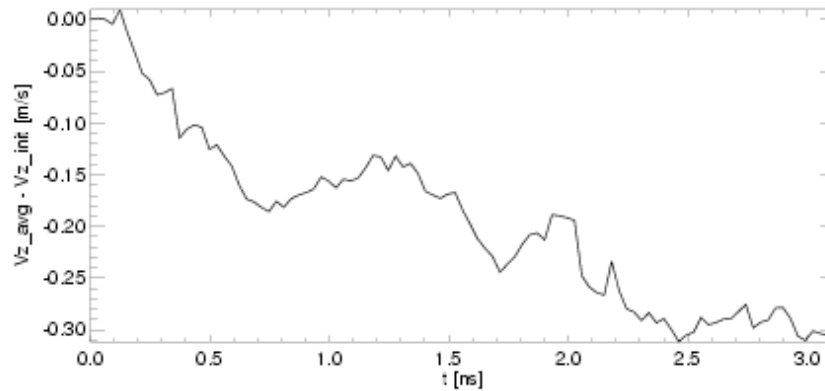
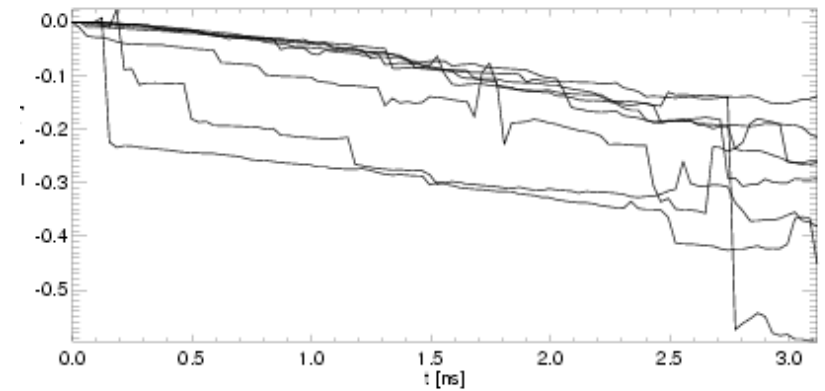
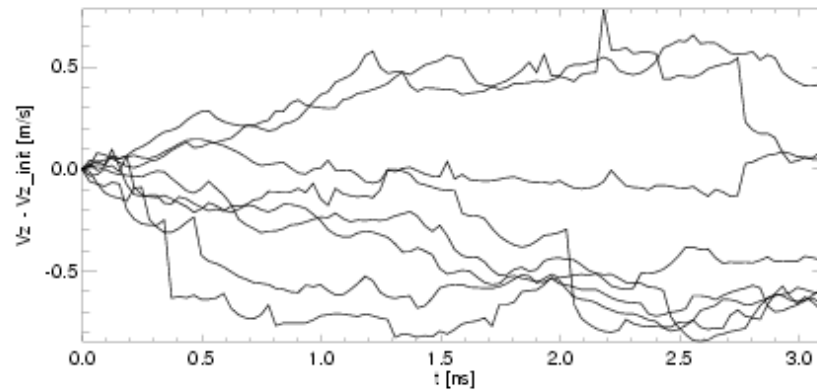
# Wiggler approach to RHIC cooler – Motivation

- Why look for alternatives to solenoid design?
  - solenoid design & beam requirements are challenging
  - accelerator physics group of the RHIC electron cooling project is now considering a wiggler-based approach
- Advantages of a wiggler
  - provides focussing & suppresses recombination
    - Modest fields (~10 Gauss) effectively reduce recombination via ‘wiggle’ motion of electrons:
    - $\rho_w = \frac{\Omega_{gyro}}{k_w^2 \gamma} \sim 1.4 \times 10^{-3} \lambda_w^2 [m] B_w [G] / \gamma$
  - e- bunch is easier: less charge and un-magnetized
- What’s the effect of ‘wiggle’ motion on cooling?
  - increase minimum impact parameter of Coulomb logarithm: ...needs to be simulated

$$\rho_{min} \textcircled{R} \rho_w$$



# Unmagnetized simulations for “wiggler” param.’s





## Wiggler fields in the beam frame

- Tests in absence of wiggler field look promising
  - unmagnetized dynamical friction agrees with theory
  - numerical e-/e+ trick suppresses diffusion by 4x
    - D. Bruhwiler et al., Proc. 33rd ICFA Beam Dynamics Workshop (2004)

- Wiggler field has following form (lab frame)

$$B_x = B_w \cosh(k_w x) \cos(k_w z)$$

$$B_y = B_w \cosh(k_w y) \sin(k_w z)$$

$$B_z = B_w \sinh(k_w y) \cos(k_w z) - B_w \sinh(k_w x) \sin(k_w z)$$

$$k_w = 2\pi / z_w$$

- Lorentz transformation to beam frame
  - yields circularly-polarized EM wave
  - relatively strong, rapidly-oscillating external fields
    - not well-suited for Hermite algorithm; need operator splitting





## Operator Splitting Approach

- Numerical technique used for ODE's & PDE's
- Consider Lorentz force equations

$$\dot{\mathbf{x}} = \mathbf{v} \quad \dot{\mathbf{v}} = \frac{q}{m} \mathbf{E}_{Coulomb} + \frac{q}{m} (\mathbf{E}_{ext} + \mathbf{v} \times \mathbf{B}_{ext})$$

- Robust 2<sup>nd</sup>-order 'Boris' uses operator splitting
  - J. Boris, Proc. Conf. Num. Sim. Plasmas, (1970), p. 3.

- Add external  $\mathbf{E}$ ,  $\mathbf{B}$  fields via operator splitting
  - Hermite algorithm: drift + coulomb fields
  - Boris 'kick': all external  $\mathbf{E}$ ,  $\mathbf{B}$  fields
- Benchmark w/ pure Hermite alg. for constant  $B_{\parallel}$



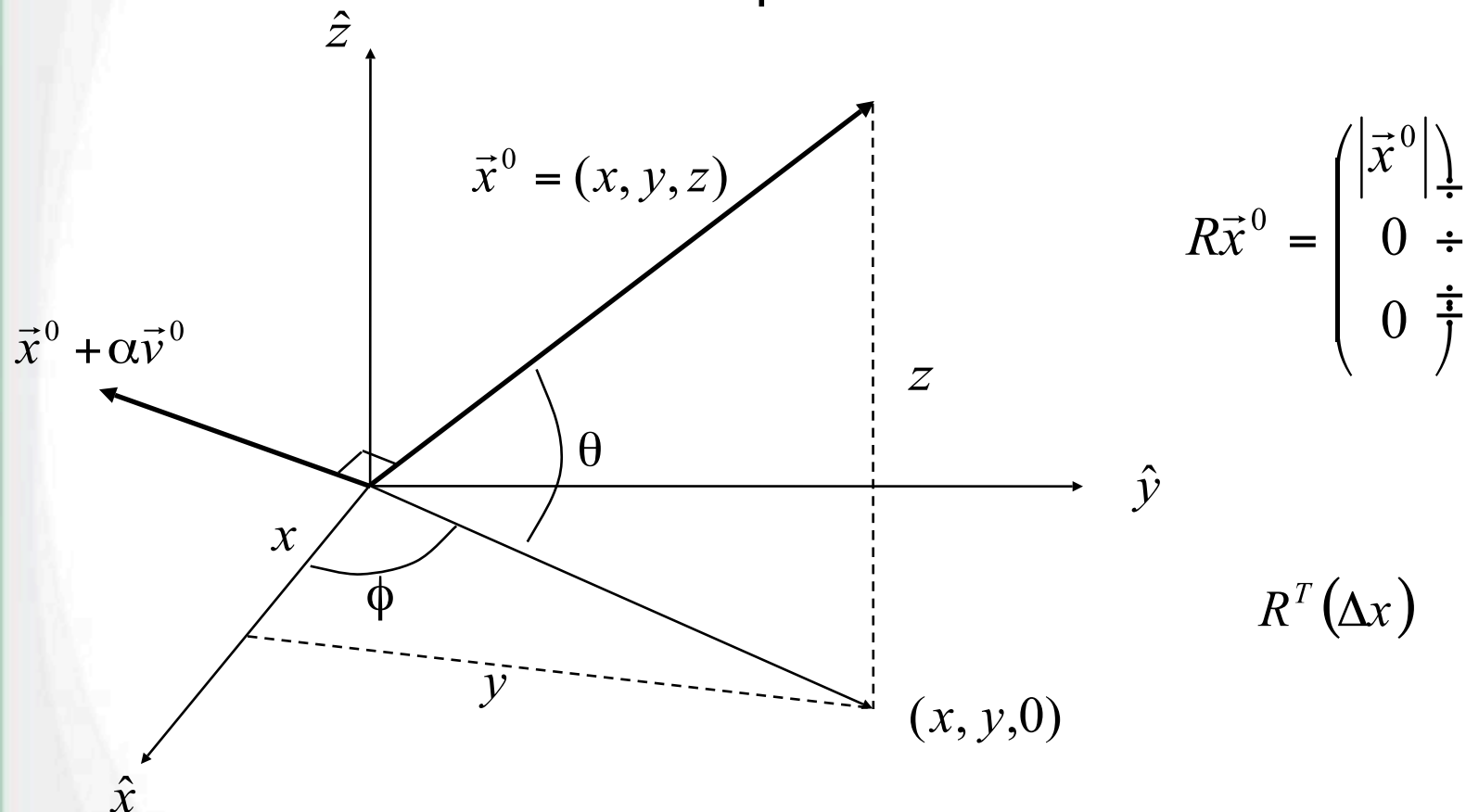
## Operator splitting is implemented for 'reduced' model

- Use analytical two-body theory for ion/e- pairs
  - handle each pair separately in center-of-mass frame
  - calculate initial orbit parameters in relevant plane
  - advance dynamics for a fixed time step
    - electron's new position and velocity are known
    - changes to ion position/velocity are small perturbations
  - total ion shift is sum of individual changes
- Initial algorithm is in place and partially tested
  - Speed and stability need to be improved
  - Initial comparisons with Hermite algorithm look good
  - Value of operator-splitting approach is verified
    - Hermite implementation in VORPAL will be modified so it can also be used with operator splitting



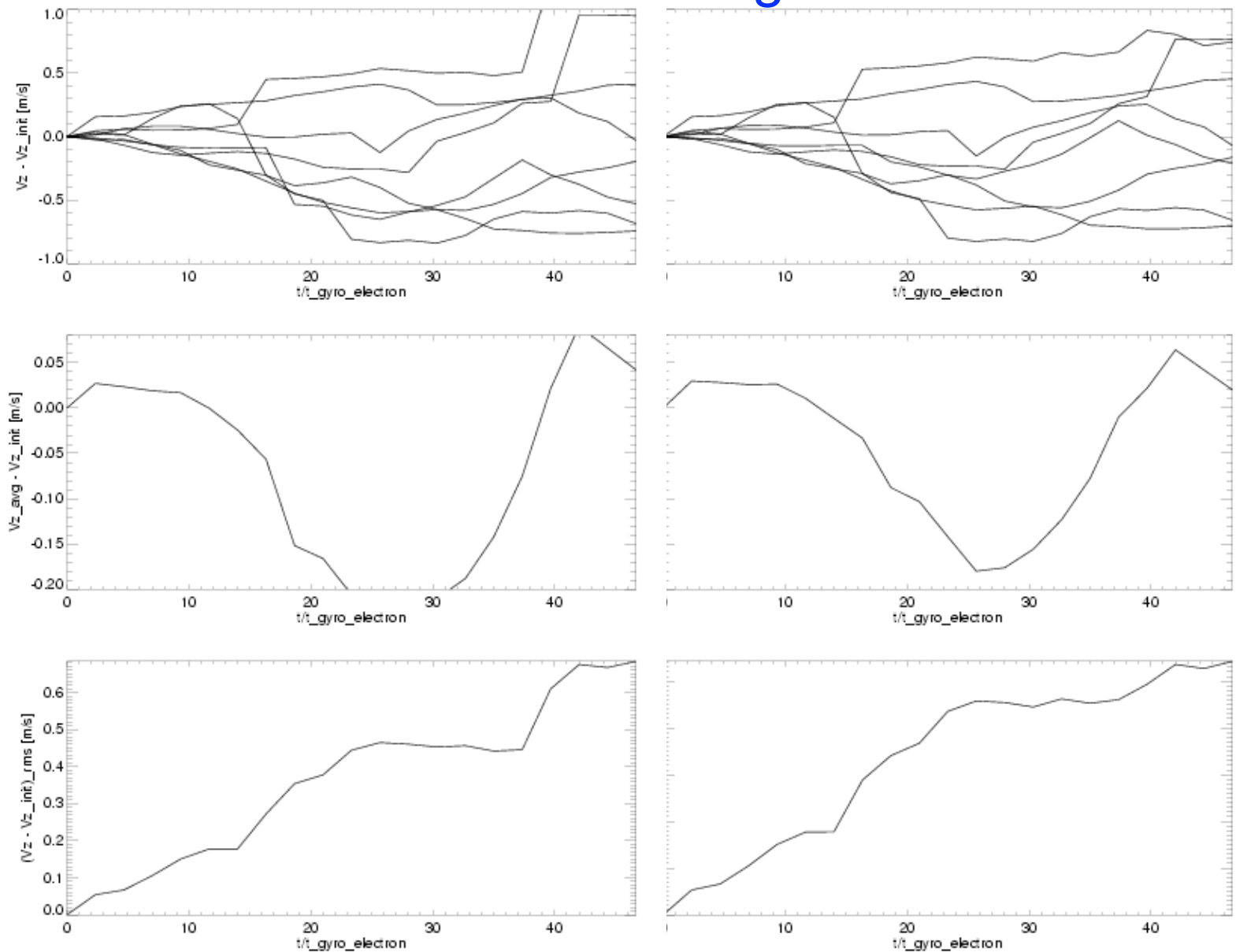
## Semi-analytic 'Reduced' Model for Binary Collisions

- Must find the plane in which partial orbit occurs
  - necessary rotations (yaw, pitch, roll) are complete
  - transformations are messy, but straightforward
  - “initial” positions & velocities obtained in this plane
- Then standard orbital parameters are calculated





# Hermite & Reduced Model agree well for B=0



Hermite Algorithm

Reduced Model



## Hermite algorithm & Reduced Model compare well for RHIC parameters w/ 5 Tesla B-field

- Agreement validates reduced model & operator splitting approach

Results for initial ion speeds:  $V_x=0.0$  m/s;  $V_z=3.0E+05$  m/s; 800 trajectories

	Hermite	Binary Collision
<delta_Vz_ion> [m/s]	-0.067	-0.066
dVz_rms/sqrt(# traj)	0.007	0.005
Time steps / TSPG*	3978 / 70	576 / 10
Processor - min/run	145	25

Results for initial ion speeds:  $V_x=2.83E+05$  m/s;  $V_z=1.0E+05$  m/s; 800 trajectories

	Hermite	Binary Collision
<delta_Vx_ion> [m/s]	-0.033	-0.043
dVx_rms/sqrt(# traj)	0.008	0.003
<delta_Vz_ion> [m/s]	-0.068	-0.062
dVz_rms/sqrt(# traj)	0.007	0.004
Time steps / TSPG*	3978 / 70	1017 / 20
Processor - min/run	144	44

\* TSPG = Time steps per gyroperiod



## Acknowledgements

We thank A. Burov, P. Schoessow, P. Stoltz, V. Yakimenko & the Accelerator Physics Group of the RHIC Electron Cooling Project for many helpful discussions. Discussions with A.K. Jain regarding solenoid field errors were especially helpful.

Work at Tech-X Corp. was supported by the U.S. Dept. of Energy under grants DE-FG03-01ER83313 and DE-FG02-04ER84094. We used resources of NERSC.