

Towards the fabrication of phosphorus qubits for a silicon quantum computer

J. L. O'Brien,^{1,2,*} S. R. Schofield,^{1,2} M. Y. Simmons,^{1,2} R. G. Clark,^{1,2} A. S. Dzurak,^{1,3} N. J. Curson,^{1,2} B. E. Kane,⁴
N. S. McAlpine,^{2,†} M. E. Hawley,⁵ and G. W. Brown⁵

¹Centre for Quantum Computer Technology, University of New South Wales, Sydney 2052, Australia

²School of Physics, University of New South Wales, Sydney 2052, Australia

³School of Electrical Engineering & Telecommunications, University of New South Wales, Sydney 2052, Australia

⁴Laboratory for Physical Sciences, University of Maryland, College Park, Maryland 20740

⁵Los Alamos National Laboratory, Los Alamos, New Mexico 87545

(Received 21 March 2001; revised manuscript received 30 April 2001; published 21 September 2001)

The quest to build a quantum computer has been inspired by the recognition of the formidable computational power such a device could offer. In particular silicon-based proposals, using the nuclear or electron spin of dopants as qubits, are attractive due to the long spin relaxation times involved, their scalability, and the ease of integration with existing silicon technology. Fabrication of such devices, however, requires atomic scale manipulation — an immense technological challenge. We demonstrate that it is possible to fabricate an atomically precise linear array of single phosphorus bearing molecules on a silicon surface with the required dimensions for the fabrication of a silicon-based quantum computer. We also discuss strategies for the encapsulation of these phosphorus atoms by subsequent silicon crystal growth.

DOI: 10.1103/PhysRevB.64.161401

PACS number(s): 03.67.Lx, 85.35.-p, 68.37.Ef, 68.43.-h

A quantum bit (or qubit) is a two-level quantum system that is the building block of a quantum computer. To date, the most advanced realizations of a quantum computer are qubit ion trap¹ and nuclear magnetic resonance^{2–4} systems. However, scaling these systems to large numbers of qubits will be difficult,⁵ making solid-state architectures,⁶ with their promise of scalability, important. In 1998 Kane proposed a solid-state quantum computer design⁷ using phosphorus ³¹P nuclei (nuclear spin $I=1/2$) as the qubits in isotopically pure silicon ²⁸Si ($I=0$). The device architecture is shown in Fig. 1(a), with phosphorus qubits embedded in silicon approximately 20 nm apart. This separation allows the donor electron wave functions to overlap, while an insulating barrier isolates them from the surface control A and J gates. These A and J gates control the hyperfine interaction between the nuclear and electron spins and the coupling between adjacent donor electrons, respectively. For a detailed description of the computer operation refer to Kane.⁷ An alternative strategy using the electron spins of the phosphorus donors as qubits has also been proposed.⁸

One of the major challenges of this design is to reliably fabricate an atomically precise array of phosphorus nuclei in silicon — a feat that has yet to be achieved in a semiconductor system. While a scanning tunneling microscope (STM) tip has been used for atomic-scale arrangement of metal atoms on metal surfaces,⁹ rearrangement of individual atoms in a semiconductor system is not straightforward due to the strong covalent bonds involved. As a result, we have employed a hydrogen resist strategy outlined in Fig. 1(b). Here the array is fabricated using a resist technology, much like in conventional lithography, where the resist is a layer of hydrogen atoms that terminate the silicon surface. A STM tip is used to selectively desorb individual hydrogen atoms, exposing the underlying silicon surface in the required array. STM induced hydrogen desorption has been developed and refined over the past ten years¹⁰ and has been proposed¹¹ for the assembly of atomically ordered device structures. We dem-

onstrate a process to adsorb single phosphine (PH₃) molecules in a predefined array, with atomic resolution, which we have developed specifically for the fabrication of a silicon quantum computer. The incorporation of these arrays in silicon is then discussed.

While hydrogen lithography at the tens of Angstrom line-width scale has been used to selectively expose the silicon surface to oxygen,¹² ammonia,¹² iron,¹³ aluminum,¹⁴ gallium,¹⁵ and cobalt,¹⁶ there has been only one recent report of adsorption at the atomic scale, in this case individual (and clusters of) silver atoms.¹⁷ Here we demonstrate the controlled adsorption of a linear array of single phosphine molecules in the extreme case of a single hydrogen atom desorption sites for direct application to the fabrication of a scalable silicon quantum computer. This technical achievement has answered the critical questions of whether a hydrogen resist is effective during exposure to phosphine and whether or not phosphine will adsorb to an STM-exposed site sufficiently small to achieve one and only one PH₃ molecule at that site.

The requirements for this quantum computer design are very stringent. In order to undertake high-resolution lithography the silicon surface must be atomically flat with a low defect density to allow the formation of a near perfect resist layer, where one hydrogen atom bonds to each surface silicon atom. The ability to then desorb individual hydrogen atoms requires a sharp, large cone angle tungsten tip in order to form < 1 nm desorption sites. These sites are subsequently exposed to high purity phosphine gas for the required phosphorus atom placement [Fig. 1(b)]. We demonstrate each of these steps below. Throughout this process particular attention must be made to avoid the introduction of any spin or charge impurities that would be fatal to the operation of the quantum computer.⁷

The final fabrication step outlined in Fig. 1(b) which is not the subject of this paper, but presents a significant challenge, is to encapsulate the phosphorus qubits in a crystalline lattice of isotopically pure silicon-28.¹⁸ The main difficulty

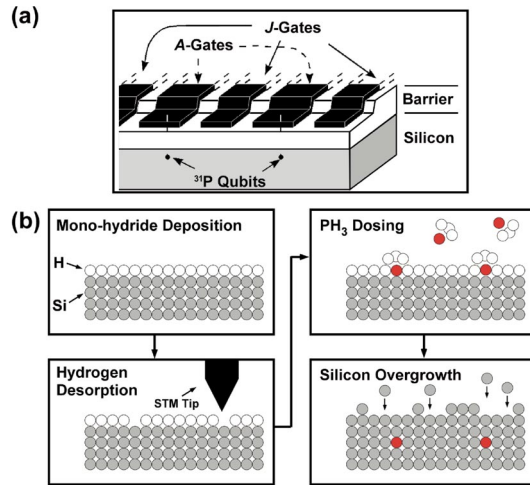


FIG. 1. (Color) A schematic of the process to fabricate the Kane architecture. (a) Detail of the Kane quantum computer architecture taken from Ref. 7, showing two phosphorus qubits in a linear array, incorporated into isotopically pure ^{28}Si and isolated from surface metal *A* and *J* gates by an insulating barrier. (b) A process to fabricate an array of phosphorus qubits in silicon. A low defect density $\text{Si}(001)2\times 1$ surface is passivated with a monolayer of hydrogen. A STM tip is used to selectively desorb hydrogen, exposing silicon on an atomic scale permitting only one phosphine molecule to adsorb at each of the required sites. Low-temperature silicon overgrowth encapsulates the phosphorus.

in this step is to ensure that the phosphorus atoms incorporate into the silicon crystal and remain in their ordered atomic array. The most direct route to achieving this is to desorb the hydrogen resist from the surface by heating to ~ 700 K followed by epitaxial silicon growth over the phosphorus array. A possible concern with heating the surface is the potential to induce lateral surface diffusion or surface segregation of the phosphorus atoms in the array. However, it is known that the phosphorus atom in the PH_2 molecule, with a single bond to the silicon, incorporates into the silicon surface with a more stable threefold coordination geometry over the lower temperature range of 550–650 K.¹⁹ We can avoid heating the surface at all during hydrogen desorption by either direct optical excitation at $\lambda = 157$ nm (Ref. 20) or electron bombardment.²¹ An alternative room-temperature approach to incorporate the phosphorus into the silicon lattice without significant diffusion involves photoinduced excitation of the attached PH_2 molecule to increase its chemical activity with the surface.²¹ Finally it is also possible to leave the hydrogen resist and grow epitaxial silicon directly on the monohydride surface.²² In all above cases 10–30 Å of epitaxial silicon can be grown at low temperatures, down to room temperature²³ to encapsulate the phosphorus, followed by elevated temperature growth at ~ 500 K to maintain crystallinity.²³ We will return to the issue of phosphorus incorporation in our final discussion.

The first step in fabricating a phosphorus qubit array is to prepare a clean, defect free silicon surface. Figure 2(a) demonstrates an optimally prepared $\text{Si}(001)2\times 1$ surface consisting of rows of σ -bonded silicon-silicon dimers. By heating the surface to 1200 °C in a variable-temperature ultrahigh vacuum STM system and performing a controlled cool-down²⁴ we have achieved large defect-free regions. The

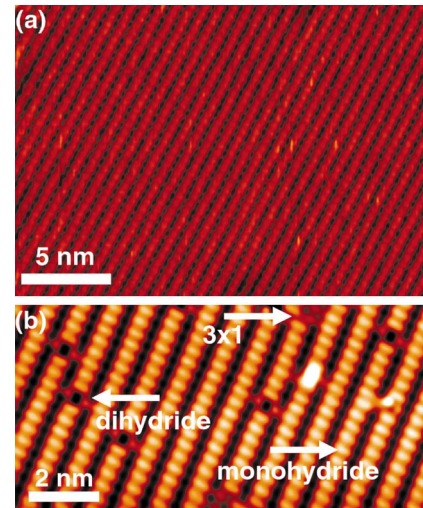


FIG. 2. (Color) Preparation of clean and hydrogen terminated surfaces. (a) Filled state STM image of a clean, very low defect density $\text{Si}(001)2\times 1$ surface. This image was acquired at a sample bias of -1.0 V and a tunneling current of 0.4 nA. (b) Fully hydrogen terminated $\text{Si}(001)2\times 1$ surface. Image acquired at -2.5 V and 0.4 nA. The surface is almost entirely 2×1 monohydride, with sites of 1×1 dihydride and 3×1 trihydride also identified.

“bean” shaped protrusions in this filled state image correspond to the charge overlap of the electrons in the dangling bonds on each silicon surface atom giving rise to a weak π -bond. These dangling bonds make the surface reactive, allowing the subsequent adsorption of species such as hydrogen and phosphine.

The adsorption of atomic hydrogen on the $\text{Si}(001)2\times 1$ surface can occur in three ways depending on the surface temperature, forming either a mono-, di-, or 3×1 mono- or dihydride.²⁵ An atomic hydrogen source consisting of a heated tungsten filament and water-cooled heat shroud has been used for this work. Figure 2(b) shows an image taken at room temperature of a low defect density $\text{Si}(001)2\times 1$ surface after exposure to high purity atomic hydrogen at 600 K. It can clearly be seen that a near uniform coverage of the

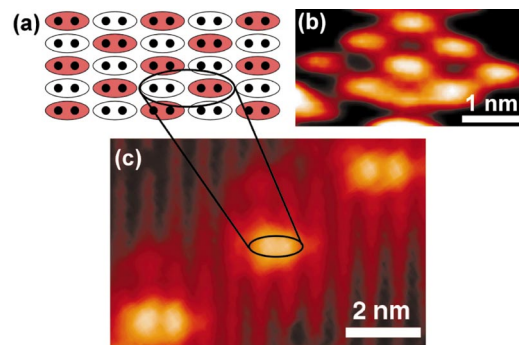


FIG. 3. (Color) Bonding structure of phosphine on $\text{Si}(001)2\times 1$. (a) Schematic of the $c(4\times 2)$ structure of PH_3 bonded to the $\text{Si}(001)2\times 1$ surface, where red dimers indicate PH_3 bonding sites. (b) A STM image of this structure, acquired at a sample bias of -3.0 V and tunneling current 0.2 nA. (c) STM image of three hydrogen desorption sites on a mono-hydride $\text{Si}(001)2\times 1:\text{H}$ surface. The highlighted regions in (a) and (c) indicate that only one phosphine molecule can adsorb at each desorption site.

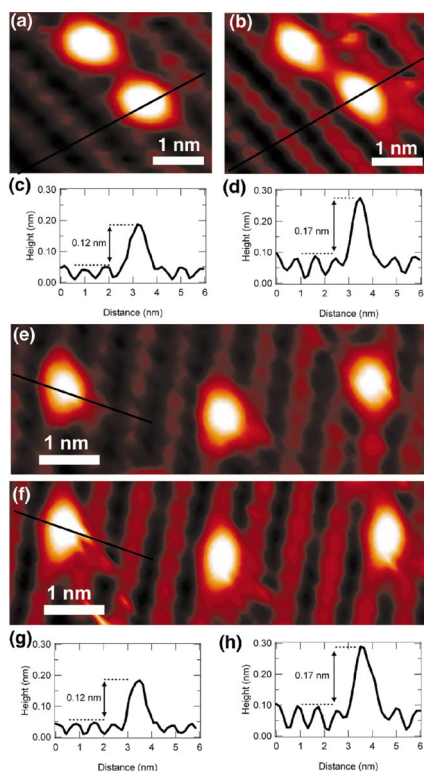


FIG. 4. (Color) Demonstration of single PH_3 adsorption through a STM patterned hydrogen resist. STM images of two desorption sites before (a) and after (b) phosphine dosing. The corresponding line profiles (c and d) show a characteristic height increase of ~ 0.05 nm. Three desorption sites before (e) and after (f) phosphine dosing and corresponding line profiles (g and h). All images were acquired at a sample bias of -1.8 V and tunneling current of 0.4 nA.

monohydride phase has occurred, where one hydrogen atom bonds to each silicon atom. Sites of di-hydride, where two hydrogen atoms bond to each silicon, and 3×1 mono- and di-hydride, a mixture of the other two phases, are also indicated. We have found that all three phases passivate the surface and act as an effective resist during subsequent exposure to phosphine. Comparison of current-voltage spectroscopy before and after hydrogen dosing (not shown) confirms the existence of a hydrogen passivated surface, with the silicon π^* -antibonding peak evolving into the silicon-hydrogen antibonding peak after hydrogenation.²⁶

A STM tip is then used to desorb hydrogen from the surface by application of a controlled voltage pulse between the tip and sample.¹⁰ Optimization of the geometry of an oxide-free tungsten tip and controlled voltage pulses (sample bias ~ 6 V and tunneling current ~ 1 nA for ~ 1 ms) makes atomic resolution desorption possible. In order to allow the adsorption of one phosphine molecule, and therefore only one phosphorus atom, it is necessary to desorb an area that exposes less than or equal to two silicon dimers as shown schematically in Fig. 3(a). This is because phosphine bonds to the $\text{Si}(001)2\times 1$ surface with a $c(4\times 2)$ surface periodicity as demonstrated in Fig. 3(b) where we have dosed a clean $\text{Si}(001)2\times 1$ surface with phosphine.¹⁹ The STM image in Fig. 3(c) shows three < 1 nm diameter hydrogen desorption sites in a row with a pitch of ~ 4 nm on a hydrogen terminated $\text{Si}(001)2\times 1$ surface. This image, with such a close

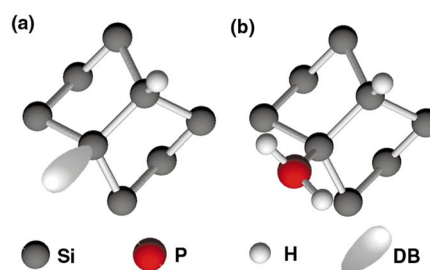


FIG. 5. (Color) Phosphine adsorption on the patterned $\text{Si}(001)2\times 1$:H surface. Schematic of a single silicon dangling bond (DB) (a) and the Si-PH_2 structure expected after dissociation of the adsorbed PH_3 (b). These models are shown looking down onto the surface, where the two silicon atoms in the center are surface atoms and the others are subsurface atoms.

spacing between sites, highlights the atomic resolution desorption achieved. The distance between sites can easily be increased to the required qubit spacing of 20 nm, and we have performed controlled lithography of single desorption sites in a line > 100 nm in length. The desorption sites in Fig. 3(c) appear as bright protrusions as a result of the extension of electron density out of the surface due to the silicon dimer surface states of the exposed silicon dangling bonds.¹⁰ Figure 3 demonstrates that these desorption sites are sufficiently small to allow only one phosphine molecule to bond to the surface at each site.

Following STM lithography to expose small regions of the $\text{Si}(001)2\times 1$ surface it is then necessary to adsorb the individual phosphorus atoms or qubits using phosphine gas. Figures 4(a) and 4(b) show atomic resolution images of two hydrogen desorption sites ~ 1.5 nm apart on a dimer row both before and after exposure to phosphine gas. The effectiveness of the hydrogen resist as a barrier to phosphine adsorption is demonstrated by the uniform hydrogen coverage after phosphine dosing except at the previously desorbed hydrogen sites. In order to observe any changes after phosphine exposure we have specifically created single hydrogen desorption sites, rather than larger desorption sites (as in Fig. 3) and present high-resolution images where the spacing between sites is very small. Here the bright protrusion at each of the desorption sites in Fig. 4(a) is the signature of the single silicon dangling bond, after desorption of just one hydrogen atom, in this case from the left side of the dimer. The remaining hydrogen on the silicon dimer is known to be transient²⁷ and we have observed it diffusing from one side of the dimer to the other.

In order to obtain high purity phosphine gas delivery, the PH_3 microdosing system and its connection to the UHV STM employed internally electropolished gas lines assembled in a clean-room environment. Mass spectra taken in the chamber during the exposure at a pressure of 10^{-8} mbar reveal no significant increase in the partial pressure of any other species. The sticking coefficient of phosphine on the clean silicon surface is 1 .¹⁹ Figure 4(b) shows the same area as Fig. 4(a) after exposure to phosphine gas at room temperature. Analysis of the line profiles in Figs. 4(c) and 4(d) shows a characteristic increase of ~ 0.05 nm in the height of the protrusion after phosphine dosing.²⁸ Figures 4(d) and 4(e) show three desorption sites, in a line perpendicular to the dimer rows, before and after phosphine dosing. The associ-

ated line profiles [Figs. 4(f) and 4(g)] again show a ~ 0.05 nm height increase. This increase, observed at all adsorption sites over several images, was calibrated against an atomic step edge on the same surface (not shown) both before and after phosphine dosing. This reproducible increase confirms the adsorption of a PH_3 molecule and corresponds to the difference between the exposed silicon dangling bond and the adsorbed phosphine. The transient nature of the hydrogen atom on the silicon dimer can account for the asymmetry of the image in Fig. 4(b) where one phosphine molecule has bonded to the left silicon in the dimer (upper) and another phosphine molecule has bonded to the right silicon in the dimer (lower).

Studies of the interaction of phosphine with the clean $\text{Si}(001)2\times 1$ surface¹⁹ suggest that PH_3 molecularly adsorbs to one end of a silicon dimer and can then dissociate to PH_2 provided enough silicon dangling bonds are available nearby for the re-adsorption of the dissociated hydrogen. Figure 5(a) shows a schematic of a single dangling bond site on a hydrogen terminated surface before phosphine dosing and Fig. 5(b) shows the proposed structure of Si-PH_2 after dissociation of the adsorbed PH_3 . In this configuration the adsorbed phosphorus atom maintains the tetrahedral geometry and threefold coordination, with the PH_2 bonded to one end of the dimer.¹⁹ However, it should be noted that the absence of available dangling bond sites on the hydrogen-terminated surface could inhibit this dissociation step. A similar dissociative process and Si-XH_2 configuration is understood to occur in the adsorption of NH_3 (Ref. 29) and AsH_3 (Ref. 30) on the $\text{Si}(001)2\times 1$ surface, where both nitrogen and arsenic are isoelectronic with phosphorus.

In the above discussion we have demonstrated the process of controlled single phosphine molecule adsorption for the extreme case of single hydrogen atom desorption sites. In future work it will be important to maintain the ordered

atomic array during the final silicon growth step [final panel of Fig. 1(b)]. Here the optimization of the size and crystallographic orientation of the desorption sites will be critical. Dissociation of singly bonded PH_2 to threefold coordinate $\text{P} + 2\text{H}$ can proceed above 500 K,²³ provided there are suitable dangling-bond sites that overcome competition with the recombinative desorption of PH_3 . An alternative option is to dose the hydrogen-patterned surface with phosphine gas at elevated temperatures (~ 550 K) to directly induce phosphorus incorporation into the first atomic layer. We have independently performed elevated temperature studies which indicate that at these temperatures, the silicon surface is stable.

We have demonstrated the effectiveness of a hydrogen resist as a barrier to phosphine adsorption and used STM lithography for the controlled placement of single phosphorus bearing molecules on a silicon surface - central to the construction of a silicon based quantum computer. This process, shown for closely spaced controlled doping, demonstrates the achievability of more widely spaced (~ 20 nm), precisely positioned phosphorus qubit arrays over large areas. While it is not yet possible to guarantee the subsequent overgrowth steps required for the Kane architecture, this paper demonstrates a significant advance in the solid-state qubit fabrication process and bodes well for the realization of a scalable silicon quantum computer. Finally, it is important to note that the fabrication strategy demonstrated here is also directly applicable to other silicon based quantum computer architectures.⁸

We would like to thank I. Andrienko, T. I. Kamins, G. J. Milburn, S. Prawer, C. Thirstrup, and R. S. Williams for useful discussions. This work was funded by the Australian Research Council Special Research Center scheme, the Australian Government, the U.S. National Security Agency, and the U.S. Advanced Research and Development Activity.

*E-mail: job@phys.unsw.edu.au

[†]Present address: School of Mathematical and Physical Sciences, University of Newcastle, Callaghan, New South Wales, 2308, Australia.

¹C. A. Sackett *et al.*, Nature (London) **404**, 256 (2000).

²I. L. Chuang *et al.*, Nature (London) **393**, 143 (1998).

³I. L. Chuang *et al.*, Phys. Rev. Lett. **80**, 3408 (1998).

⁴J. A. Jones and M. Mosca, Phys. Rev. Lett. **83**, 1050 (1999).

⁵W. S. Warren, Science **277**, 1688 (1997).

⁶D. Loss and D. P. DiVincenzo, Phys. Rev. A **57**, 120 (1998).

⁷B. E. Kane, Nature (London) **393**, 133 (1998).

⁸R. Vrijen *et al.*, Phys. Rev. A **62**, 012 306 (2000).

⁹M. F. Crommie *et al.*, Science **262**, 218 (1993).

¹⁰C. Thirstrup *et al.*, Surf. Sci. **411**, 203 (1998), and references therein.

¹¹J. R. Tucker and T.-C. Shen, Solid-State Electron. **42**, 1061 (1998).

¹²J. W. Lyding *et al.*, J. Vac. Sci. Technol. B **12**, 3735 (1994).

¹³D. P. Adams *et al.*, J. Vac. Sci. Technol. B **14**, 1642 (1996).

¹⁴T.-C. Shen *et al.*, Phys. Rev. Lett. **78**, 1271 (1997).

¹⁵T. Hashizume *et al.*, Surf. Sci. **386**, 161 (1997).

¹⁶G. Palasantzas *et al.*, J. Appl. Phys. **185**, 1907 (1999).

¹⁷M. Sakurai *et al.*, Phys. Rev. B **62**, 16 167 (2000).

¹⁸H. Bracht *et al.*, Phys. Rev. Lett. **81**, 393 (1998).

¹⁹D. S. Lin *et al.*, Phys. Rev. B **61**, 2799 (2000), and references therein.

²⁰T. Vondrak and X.-Y. Zhu, Phys. Rev. Lett. **82**, 1967 (1999).

²¹H. Guo *et al.*, Prog. Surf. Sci. **62**, 239 (1999).

²²M. Copel and R. M. Tromp, Phys. Rev. Lett. **72**, 1236 (1994).

²³D. J. Eaglesham *et al.*, Phys. Rev. Lett. **65**, 1227 (1990).

²⁴B. S. Swartzentruber *et al.*, J. Vac. Sci. Technol. A **7**, 2901 (1989).

²⁵K. Oura *et al.*, Surf. Sci. Rep. **35**, 1 (1999).

²⁶R. J. Hamers *et al.*, Phys. Rev. Lett. **59**, 2071 (1987).

²⁷E. Hill *et al.*, Phys. Rev. B **60**, 15 896 (1999).

²⁸A slightly larger height in the line profile of the dimer rows is observed after PH_3 dosing. Such a difference can frequently occur due to minor changes in imaging conditions between scans, which results in the STM tip extending further into the gap between dimer rows. However the height difference due to PH_3 adsorption is measured from the top of the dimer rows to the top of the protrusion and is not therefore affected by this.

²⁹Zhi-Heng Loh and H. C. Kang, J. Chem. Phys. **112**, 2444 (2000), and references therein.

³⁰J. E. Northrup, Phys. Rev. B **51**, 2218 (1995).