

Lattice2004 @ Fermilab, Jun. 21-26, 2004

Pentaquarks

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Outline of this talks

- Brief review of the pentaquarks
 - Theoretical interests
 - Experimental facts
- Survey of available lattice results
 - Strategy of studying pentaquarks
 - Five studies, all quenched
- Summary and Outlook

Multi-quark hadrons

- 📌 **QCD** (color confinement) may not exclude the presence of the multi-quark hadrons such as

$$Q\bar{Q}Q\bar{Q}, \quad QQQQ\bar{Q}, \quad (QQQ)^2, \quad Q\bar{Q}Q\bar{Q}Q\bar{Q}, \quad \dots$$

- 📌 **“Exotic”** multi-quark hadrons should have quantum numbers which can not be accommodated by $Q\bar{Q}$ or QQQ

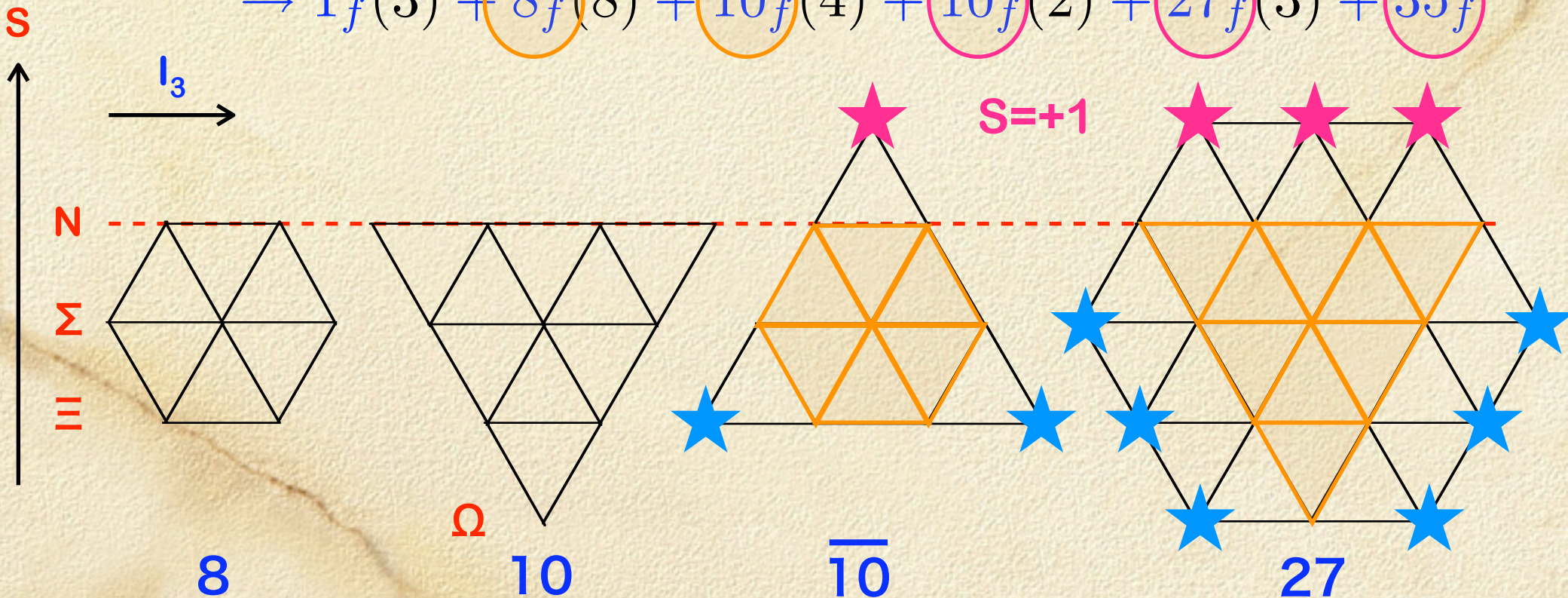
Pentaquark baryons



Flavor structure (SU(3) case)

$$3_f \otimes 3_f \otimes 3_f \otimes 3_f \otimes \bar{3}_f$$

$$\rightarrow 1_f(3) + 8_f(8) + 10_f(4) + \bar{10}_f(2) + 27_f(3) + 35_f$$



Exotic anti-decuplet baryons

📌 **SU(3) Skyrmion** in the rigid rotator approach

☑ The **third** baryons in the $(10^*, 1/2^+)$ as “Rotational Band”

cf. $(8, 1/2^+)$, $(10, 3/2^+)$

Manohar(84), Chemtob(85)

📌 A narrow exotic **S=+1 baryon $\Theta^+(Z^+)$** predicted by the chiral quark-soliton model

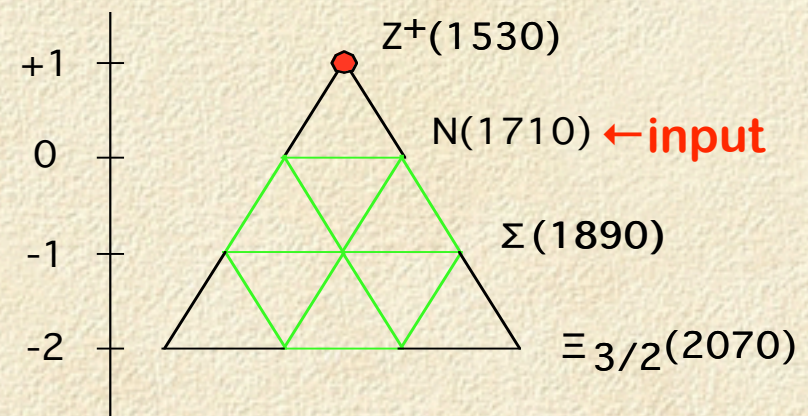
Diakonov et al. Z. Phys. A359 (97) 305

☑ Exotic: **S=+1** in the $10^*(I=0)$

☑ Low mass: **1530 MeV**

☑ Narrow width: ~~<15~~ → **30 MeV**
(Jaffe)

☑ **$J^P=1/2^+$**

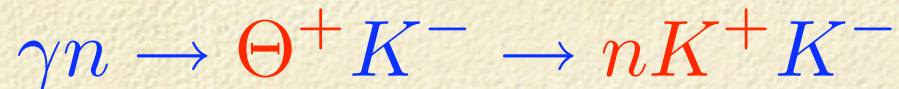


Discovery of Exotic $S=+1$ Baryon

T. Nakano et al.

Phys.Rev.Lett.91 (2003) 012002

Laser-Electron Photon facility (LEPS)@Spring-8



$Mass = 1540 \pm 10 \text{ MeV}$

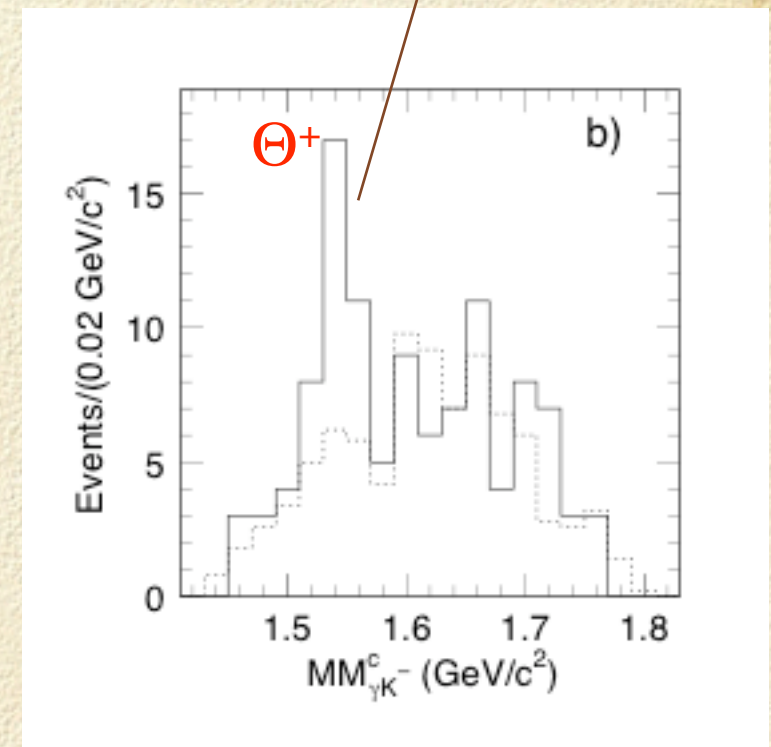
$Width \leq 25 \text{ MeV}$

☑ Positive Strangness ($uudd\bar{s}$)

☑ Very narrow width

☑ $I=0$ (no pK^+ partner)

☑ Spin and Parity are undetermined.



Confirmation from other experiments



DIANA/ITEP (hep-ex/0304040)
Mass = 1539 ± 2 MeV,
Width < 9 MeV



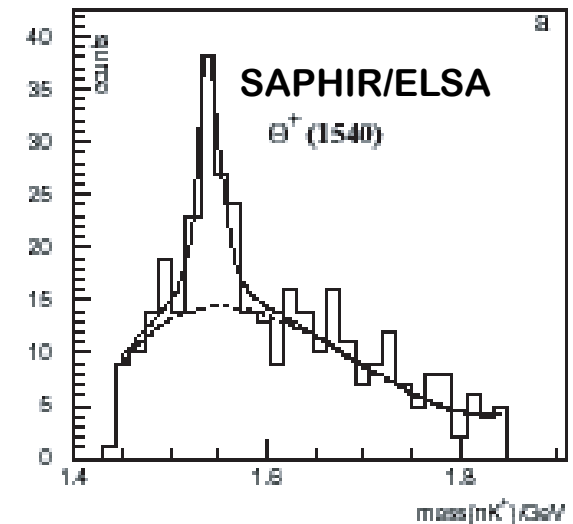
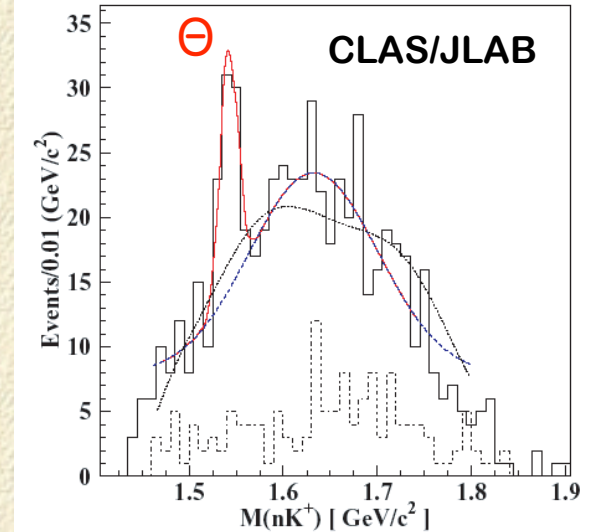
CLAS/JLAB (hep-ex/0307018)
Mass = 1542 ± 5 MeV,
Width < 21 MeV



SAPHIR/ELSA (hep-ex/0307083)
Mass = 1540 ± 4 MeV,
Width < 25 MeV



HERMES/DESY (hep-ex/0312044)
Mass = 1528 ± 2.6 MeV,
Width < 19 ± 5 MeV



But, spin and parity are still undetermined.

Summary of Experiments

Where	Reaction	Mass	Width	σ 's*
LEPS	$\gamma C \rightarrow K^+ K^- X$	1540 +- 10	< 25	4.6
DIANA	$K^+ X e \rightarrow K^0 p X$	1539 +- 2	< 9	4.4
CLAS	$\gamma d \rightarrow K^+ K^- p(n)$	1542 +- 5	< 21	5.2
SAPHIR	$\gamma p \rightarrow K^+ K^0(n)$	1540 +- 6	< 25	4.8
ITEP	$\nu A \rightarrow K^0 p X$	1533 +- 5	< 20	6.7
CLAS	$\gamma p \rightarrow \pi^+ K^- K^+(n)$	1555 +- 10	< 26	7.8
HERMES	$e^+ d \rightarrow K^0 p X$	1528 +- 3	13 +- 9	~5
ZEUS	$e^+ p \rightarrow e' K^0 p X$	1522 +- 3	8 +- 4	~5
COSY	$pp \rightarrow K^0 p \Sigma^+$	1530 +- 5	< 18	4-6

*Gaussian statistical significance: estimated background fluctuation

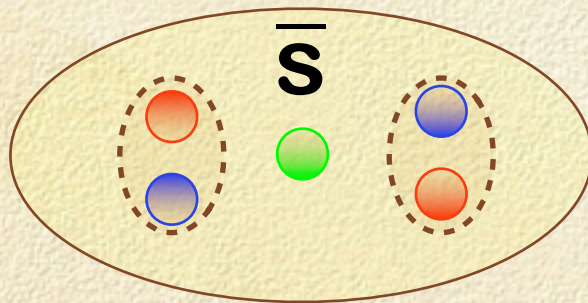
The existence of the Θ has been established.

Correlated quark (diquark) model

- ☑ Quark models for Θ^+ : natural J^P -assignment $\rightarrow 1/2^-$
 but, “fall-apart” into KN in a **S-wave**

- ☑ **Diquark** correlation : $l=J=0$, diquark $(3_c^*, 3_f^*)$

Jaffe-Wilczek, Phys.Rev.Lett. 91 (03) 232003



$$[ud]_0 \leftrightarrow [ud]_0$$

$L=1$

Relative P-wave is necessary
between the pairs of diquarks

$$\Rightarrow J^P = 1/2^+$$

$$(3_f \otimes 3_f) \otimes (3_f \otimes 3_f) \otimes \bar{3}_f \rightarrow \bar{3}_f \otimes \bar{3}_f \otimes \bar{3}_f = 1_f + 8_f + 8_f + \bar{10}_f$$

Diquark Model vs. Chiral Soliton

$\Theta(1540) \Leftrightarrow J^\pi=1/2^+, I=0, S=+1$ pentaquark

Diquarks: $8_f \oplus 10_f^*$

Chiral Soliton: **ONLY** 10_f^*

Ideally Mixed

$qqss\bar{s}$ — Σ_s

$qqqs\bar{s}$ — $N_s, \Xi_{3/2} \Rightarrow 1.75\text{GeV}$

$qqqq\bar{s}$ — Θ, Σ, Λ

$qqqq\bar{q}$ — $N \Rightarrow \text{Roper: } N(1440)$

— $\Xi_{10} \Rightarrow 2.07\text{GeV}$

— Σ_{10}

— $N_{10} \Leftarrow N(1710)$

— Θ

The linear-in- m_s treatment

Ξ^{--} ($dds\bar{s}^{\text{bar}}u$) baryon candidate

NA49/CERN SPS

Phys.Rev.Lett.92:042003,2004

Exotic $S=-2$ and $Q=-2$

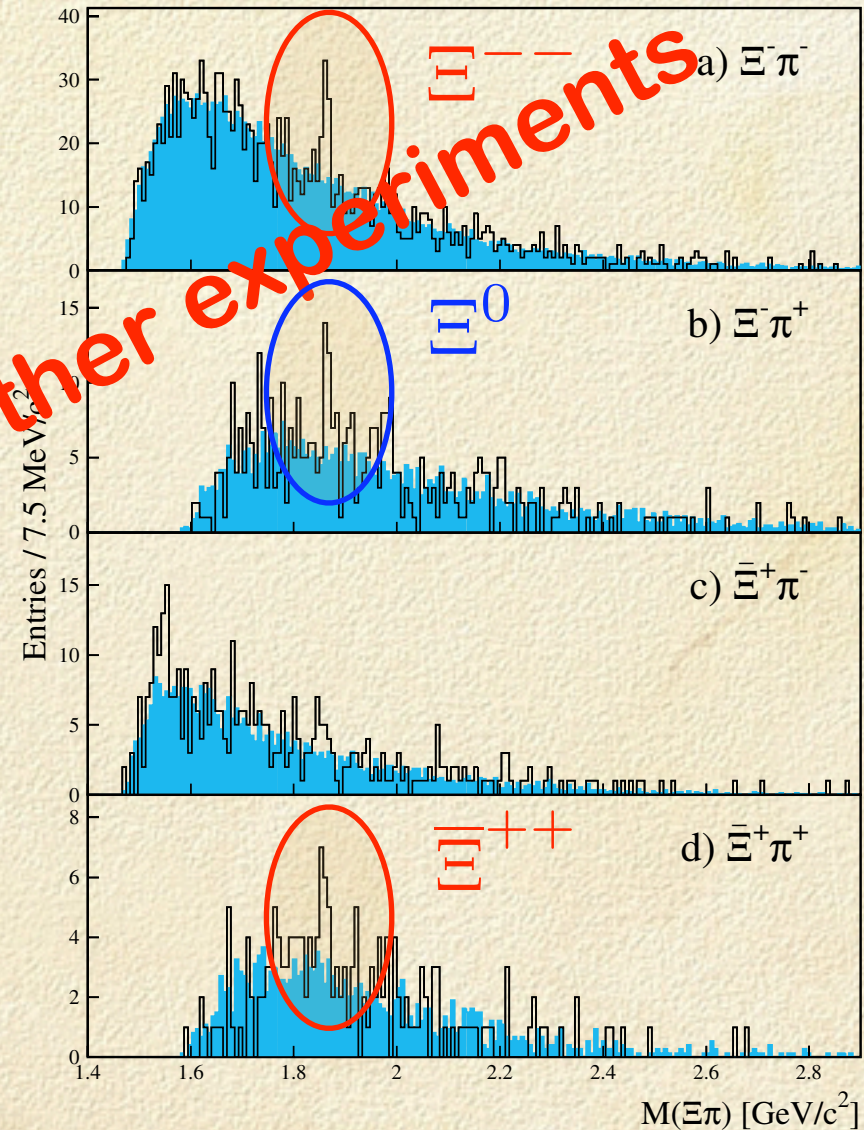
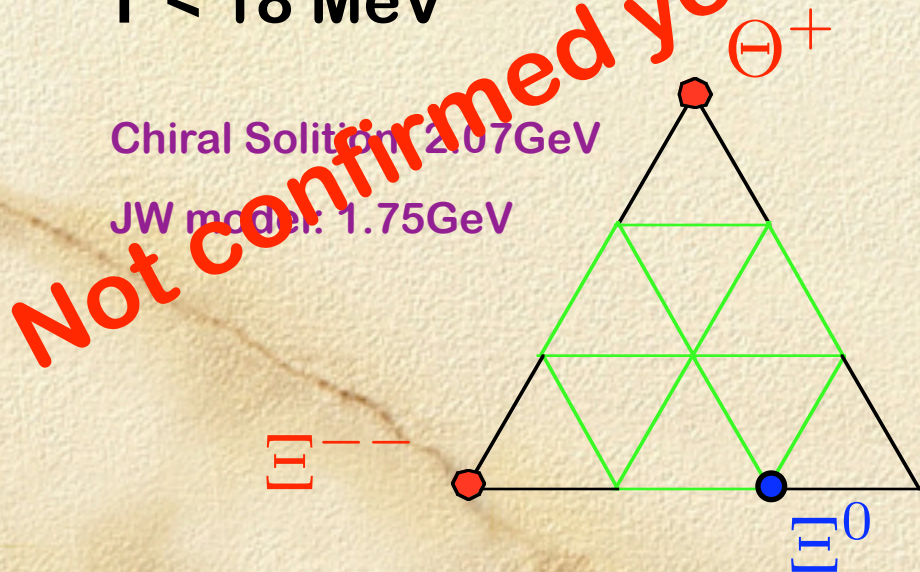
baryon resonance in $\Xi^- \pi^-$

$M = 1.862 \text{ GeV}$

$\Gamma < 18 \text{ MeV}$

Chiral Soliton: 2.07 GeV

JW model: 1.75 GeV



Charm (bottom) pentaquarks

$(uudd\bar{s}) \rightarrow (uudd\bar{c})$ or $(uudd\bar{b})$

Very simple mass estimations:

$$M(\Theta_c^0) = M(\Theta^+) + M(\Lambda_c) - M(\Lambda)$$

$$\approx 2710 \text{ MeV} < M(D) + M(N) \approx 2810 \text{ MeV}$$

$$M(\Theta_b^+) = M(\Theta^+) + M(\Lambda_b) - M(\Lambda)$$

$$\approx 6050 \text{ MeV} < M(B) + M(N) \approx 6200 \text{ MeV}$$

Jaffe-Wilczek

Suggest: charm (bottom) pentaquarks may be **bound states** ???

Θ_c ($uudd^{\text{bar}}c$) baryon candidate

H1/DESY (hep-ex/0403017)

Exotic $C=-1$ and $Q=0$
baryon resonance in D^*p

$M = 3.099 \text{ GeV}$

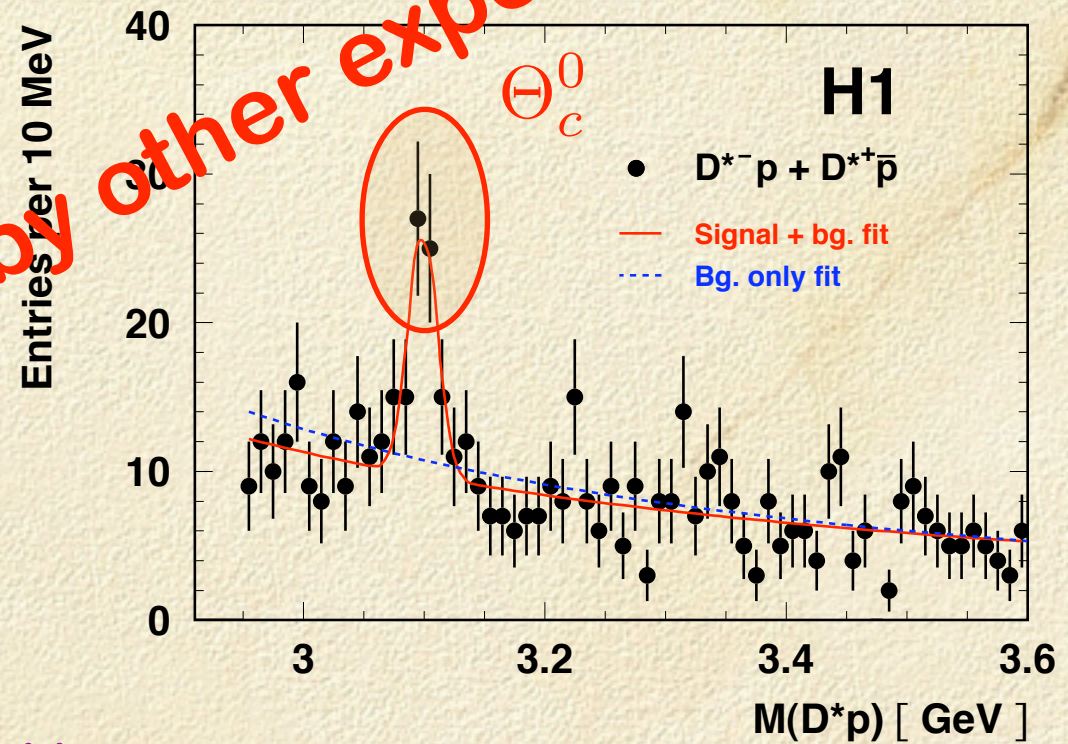
$\Gamma < 12 \text{ MeV}$

Negative charm $\rightarrow (uudd\bar{c})$

Note:

DN threshold = 2.808 GeV

D^*N threshold = 2.948 GeV



Not confirmed yet by other experiments

What can lattice QCD do?



Does the spectrum of QCD possess the $\Theta^+(1540)$?



What is **spin** and **parity** of the $\Theta^+(1540)$?



Are there other pentaquark baryons?

Maximal knowledge about those matters is essential to understanding the structure of the pentaquark state.



Lattice QCD has a chance to answer the last two questions before experimental efforts.

Lattice studies of pentaquarks



Four studies posted to HEP-archive

- Csikor-Fodor-Katz-Kovacs, [hep-lat/0309090](https://arxiv.org/abs/hep-lat/0309090), [JHEP 0311 \(03\) 070](https://arxiv.org/abs/hep-lat/0309090).
- Sasaki, [hep-lat/0310014](https://arxiv.org/abs/hep-lat/0310014).
- Chiu-Hsieh, [hep-lat/0403020](https://arxiv.org/abs/hep-lat/0403020), [0404007](https://arxiv.org/abs/hep-lat/0403020).
- Kentucky Collab., [hep-ph/0406196](https://arxiv.org/abs/hep-ph/0406196).



One preliminary result reported at some conferences

- Sigaev-Jahn-Negele (MIT Collab.), [Quark-Nuclear-Physics04 etc.](https://arxiv.org/abs/hep-ph/0406196)



New results will be presented at this conference

- Ishii et al. (TIT Collab.), [June 24 \(Thu\) 9:20](#)
- Chiu-Hsieh, [June 24 \(Thu\) 9:40](#)
- Mathur et al. (Kentucky Collab.), [June 24 \(Thu\) 10:00](#)
- Koutsou, [Poster](#)

Main difficulty in lattice study

A simple minded study of pentaquark state with

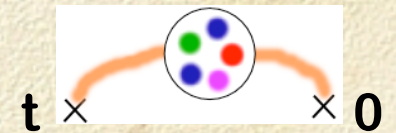
$$\Theta^+ \sim \underbrace{\varepsilon_{abc} d_a d_b u_c}_N \times \underbrace{\bar{s}_e u_e}_K$$

How can we distinguish between

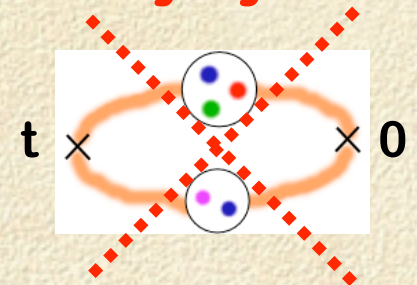
the mass of the pentaquark state

and

the total energy of the interacting KN two-body system



Choose a specific operator with as little overlap
 The 2-pt function $\langle \Theta(t)\Theta(0) \rangle$ should be
 with the KN scattering state as possible
 dominated by the latter if $M_\Theta > M_N + M_K$
 $\langle \Theta^+ | 0 | 0 \rangle \gg \langle K + N | 0 | 0 \rangle$



KN threshold (1)

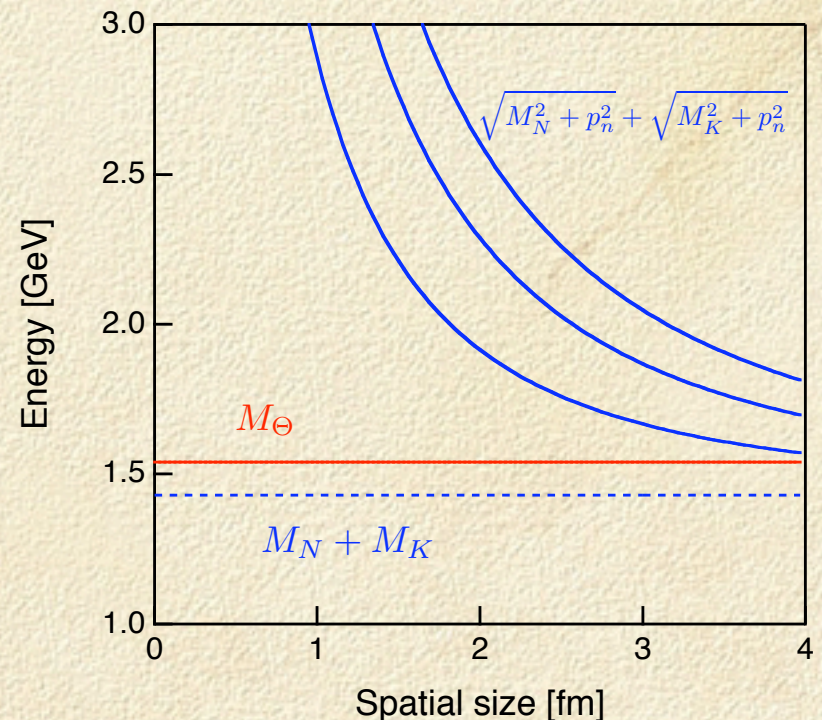
$\Theta(1/2^+) \rightarrow (\text{KN})_{\text{P-wave}} : \sqrt{M_N^2 + p_{\text{min}}^2} + \sqrt{M_K^2 + p_{\text{min}}^2}$

$\Theta(1/2^-) \rightarrow (\text{KN})_{\text{S-wave}} : M_N + M_K$ **non-interacting**

All momentum are quantized

$$|\vec{p}| = \sqrt{n} |\vec{p}_{\text{min}}|, \quad |\vec{p}_{\text{min}}| = 2\pi/L$$

⇒ the P-wave KN threshold can be lifted by changing spatial size L



KN threshold (2)

How valid is the previous estimation ?

- Energy shift in the finite volume:

$$\Delta M = E_{KN} - (M_N + M_K)$$

$$= -\frac{2\pi(M_N + M_K)a_0}{M_N M_K L^3} \left[1 - 2.834 * \frac{a_0}{L} + 6.375 * \left(\frac{a_0}{L}\right)^2 \right]$$

$L \gg a_0$ Lüscher formula

- KN scattering length(volume) is quite small in $l=0$ channel

	$l=0$	$l=1$
S-wave (fm)	0.0 ± 0.03	-0.32 ± 0.02
P-wave (fm ³)	0.08 ± 0.01	-0.16 ± 0.1

Choice of operator (1)

Nucleon spectroscopy

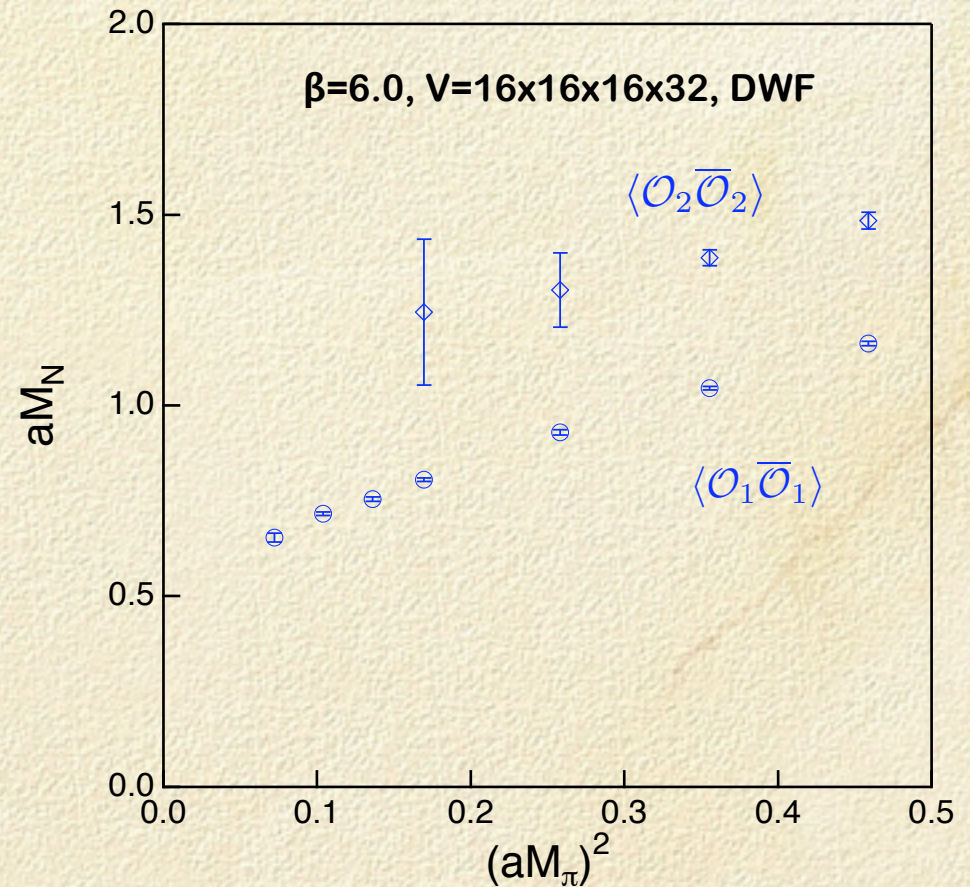
$$\mathcal{O}_1 = \varepsilon_{abc} [u_a^T C \gamma_5 d_b] u_c$$

$$\mathcal{O}_2 = \varepsilon_{abc} [u_a^T C d_b] \gamma_5 u_c$$

The second operator vanishes in the non-relativistic limit.

⇒ small overlap with the nucleon

$$|\langle N | \mathcal{O}_2 | 0 \rangle| \approx 0$$



Bowler et al., Nucl. Phys. B240 (1984) 213.

Leinweber, Phys. Rev. D51 (1995) 6383.

Sasaki-Blum-Ohta, Phys. Rev. D65 (2002) 074503.

Choice of operator (1)

Nucleon spectroscopy

$$\mathcal{O}_1 = \varepsilon_{abc} [u_a^T C \gamma_5 d_b] u_c$$

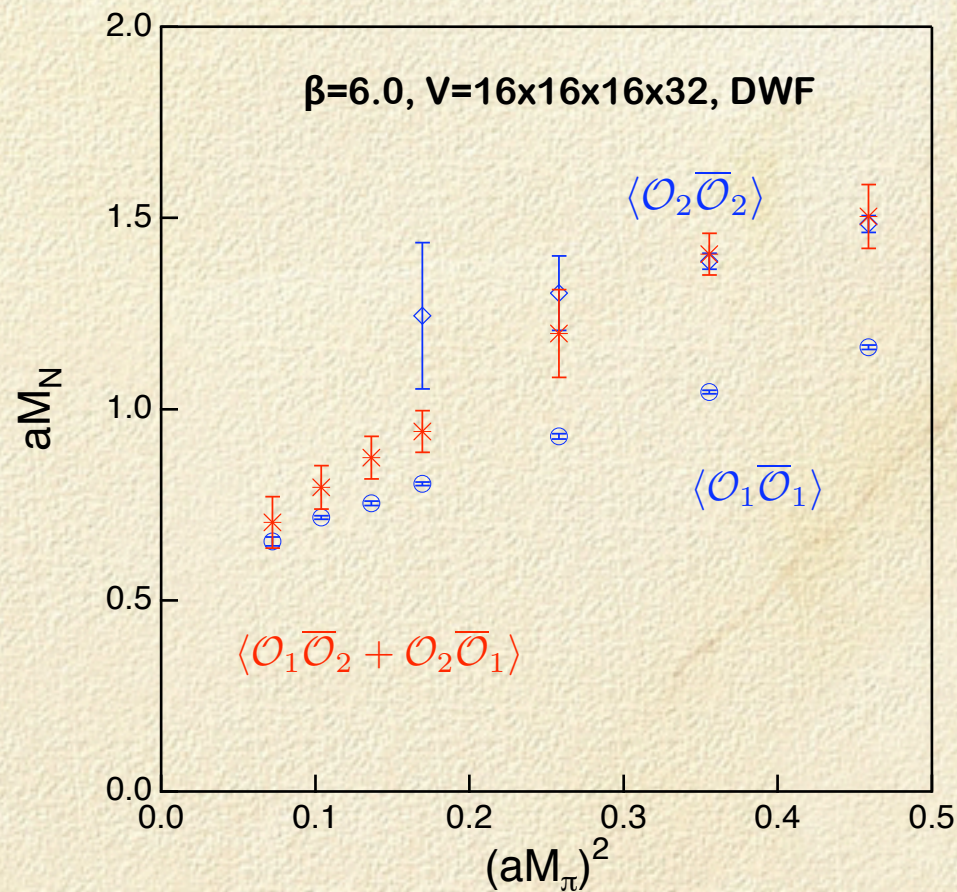
$$\mathcal{O}_2 = \varepsilon_{abc} [u_a^T C d_b] \gamma_5 u_c$$

The second operator vanishes in the non-relativistic limit.

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The cross correlation suggests that it is evident in the heavy quark regime, but it might be no longer robust in light quark regime.



Choice of operator (2)

- Simple minded operator as **Nucleon** ⊗ **Kaon**

$$\mathcal{O}_{I=0/1} = \varepsilon_{abc} [u_a^T C \gamma_5 d_b] \{u_c (\bar{s}_e \gamma_5 d_e) \mp (u \leftrightarrow d)\} \quad \text{Kentucky}$$

- **Color variant** of Nucleon ⊗ Kaon

$$\mathcal{O}_{I=0/1} = \varepsilon_{abc} [u_a^T C \gamma_5 d_b] \{u_e (\bar{s}_e \gamma_5 d_c) \mp (u \leftrightarrow d)\} \quad \text{Csikor et al.}$$

- Exotic description as **diquark-diquark-antiquark**

$$\mathcal{O}_{I=0} = \varepsilon_{abc} \varepsilon_{aef} \varepsilon_{bgh} [u_e^T C \Gamma_1 d_f] [u_g^T C \Gamma_2 d_h] C \bar{s}_c^T \quad \text{Sasaki}$$

$$\Gamma_1 \neq \Gamma_2, \quad \Gamma_{1,2} = 1, \gamma_5, \gamma_\mu \gamma_5$$

Choice of operator (2)

- Simple minded operator as **Nucleon** \otimes **Kaon**

$$\mathcal{O}_{I=0/1} = \varepsilon_{abc} [u_a^T C \gamma_5 d_b] \{u_c (\bar{s}_e \gamma_5 d_e) \mp (u \leftrightarrow d)\} \quad \text{Kentucky}$$

- Color variant** of **Nucleon** \otimes **Kaon**

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In the non-relativistic limit (**heavy** quark regime),

→→ Small overlap with KN two-hadron state is expected.

Csikor-Fodor-Katz-Kovacs (1)

hep-lat/0309090, JHEP 0311 (03) 070

Wilson fermions (Quench) + Wilson gauge action

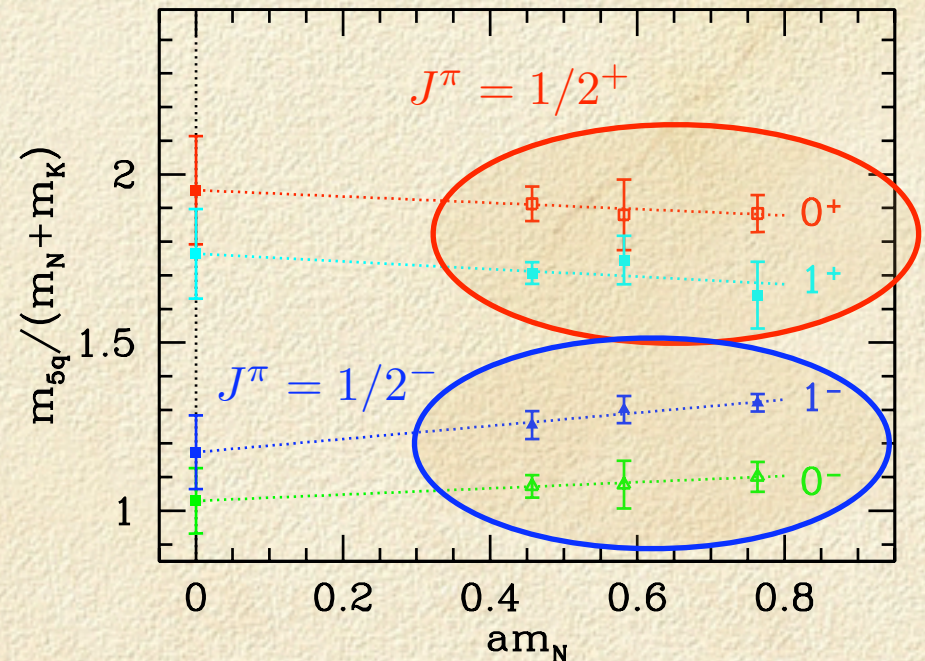
■ $L \sim 2.0$ fm, $a = 0.17, 0.12, 0.09$ fm ($\beta = 5.7, 5.85, 6.0$)

○ linear continuum extrapolation

■ m_π from 0.4 - 0.6 GeV

■ $J = 1/2, l = 0, 1$, parity projection

■ Color variant of N times K



Csikor-Fodor-Katz-Kovacs (2)

[hep-lat/0309090](#), [JHEP 0311 \(03\) 070](#)



All analysis is done by single-exp. fits and single operator

$$\frac{M_{5Q}}{M_K + M_N} = 1.073(34) \quad \text{at } \beta = 6.0$$

negative parity

Csikor-Fodor-Katz-Kovacs (2)

hep-lat/0309090, JHEP 0311 (03) 070

- ☑ All analysis is done by single-exp. fits and single operator

$$\frac{M_{5Q}}{M_K + M_N} = 1.073(34) \quad \text{at } \beta = 6.0$$

negative parity

- ☑ Variational approach (at the **heaviest** quark mass)

- Mix two operators $\mathcal{O}_1 + \alpha\mathcal{O}_2$ $\begin{cases} \mathcal{O}_1: \text{color variant of } \mathcal{O}_K \times \mathcal{O}_N \\ \mathcal{O}_2: \mathcal{O}_K \times \mathcal{O}_N \end{cases}$
- Choose α to isolate **the first excited state**
- At the optimized α

$$\frac{E_0}{M_K + M_N} = 0.994(18)$$

KN scattering state

$$\frac{E_1}{M_K + M_N} = 1.074(20)$$

Sasaki (1)

hep-lat/0310014

Wilson fermions (Quench) + Wilson gauge action

■ $L \sim 2.2$ fm ($32^3 \times 48$), $a = 0.07$ fm ($\beta = 6.2$)

■ m_π from 0.6 - 1.0 GeV

■ $J = 1/2$, $l = 0$, parity projection

■ Includes charm sector

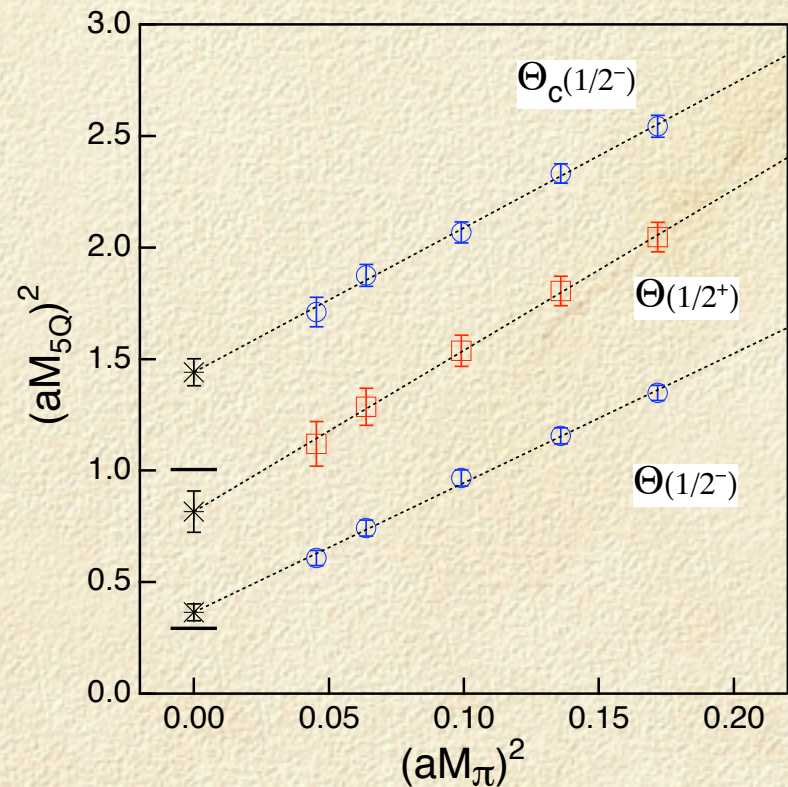
■ Diquark-diquark-antiquark

$$J^\pi = 1/2^-$$

$$M_{\Theta^+} = 1.76(9) \text{ GeV}$$

$$J^\pi = 1/2^+$$

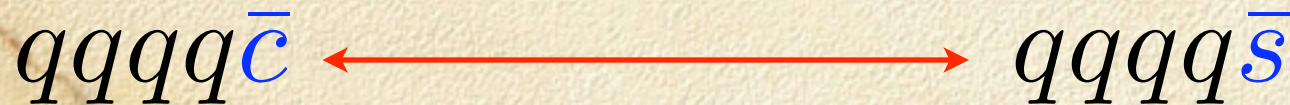
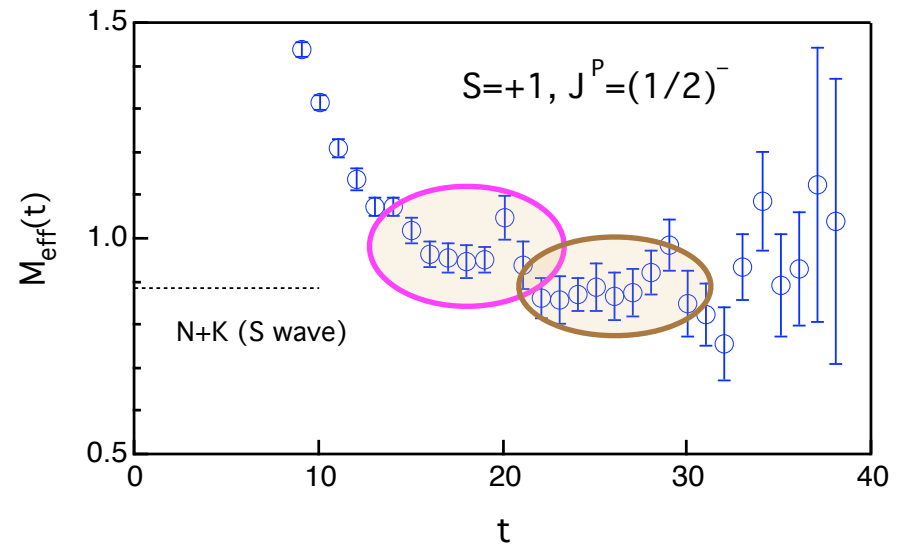
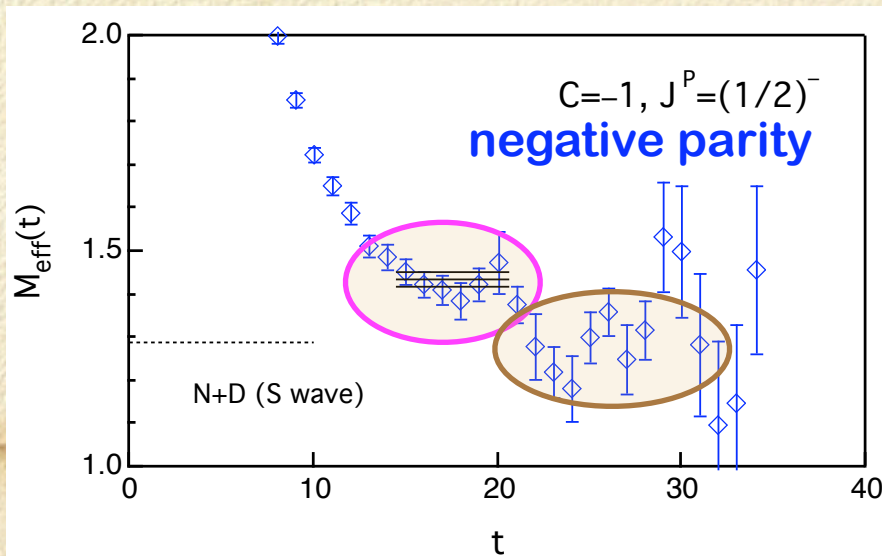
$$M_{\Theta^+} = 2.62(9) \text{ GeV}$$



Sasaki (2)

hep-lat/0310014

- ✓ Test diquark-diquark-antiquark op. in **heavy** quark regime
- Anti-charmed pentaquarks $uudd\bar{c}$



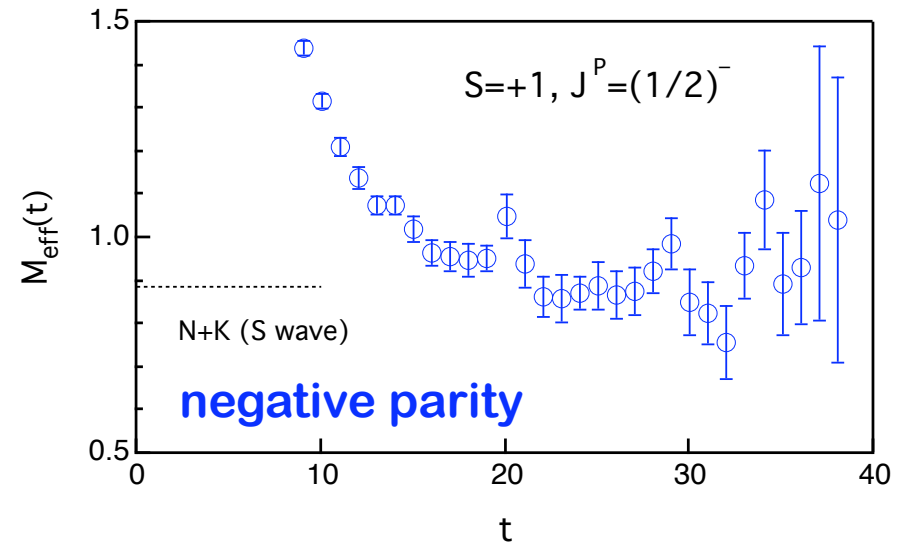
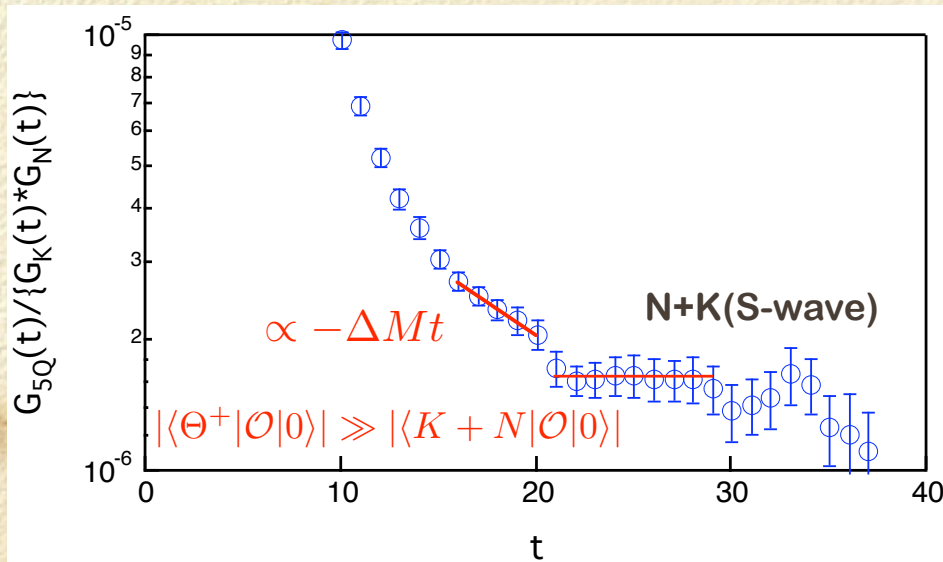
Sasaki (3)

hep-lat/0310014



Double-exp. fits accept the presense of two indep. states, which are close to each other. ($A_{5Q} \gg A_{KN}$)

- Ground state should be the S-wave KN scattering state.



- First excited state might be the pentaquark state. (need confirmation)

Summary of exploratory studies

- CFKK and Sasaki both claim :
 - Can not accommodate the pentaquark state near the KN threshold in the **positive** parity channel
 - Some indications for the presence of the pentaquark state near the KN threshold in the **negative** parity channel
 - The spin-parity of the $\Theta^+(1540)$ is most likely **$1/2^-$**

Chiu-Hsieh (1)

hep-lat/0403020, 0404007

Chiral fermions (Quench) + Wilson gauge action

■ $L \sim 1.6$ fm ($20^3 \times 40$), $a = 0.08$ fm ($\beta = 6.1$)

■ m_π from 0.4 - 1.0 GeV

■ $J = 1/2, l = 0$, parity projection

■ Includes charm sector

■ Same operator as Sasaki

	5Q baryon	J^Π	mass	Expt.
10	$\Theta([ud][ud]\bar{s})$	$1/2^+$	1539(95)	$\Theta^+(1540)$
	$N_s([ud][us]_+\bar{s})$	$1/2^+$		
	$\Sigma_s([us][us]\bar{s})$	$1/2^+$		
	$\Xi([ds][ds]\bar{u})$	$1/2^+$	1826(87)	$\Xi_{3/2}^-(1860)$
8	$N([ud][ud]\bar{d})$	$1/2^+$	1460(51)	$N(1440)P_{11}$
	$\Sigma([ud][us]_+\bar{d})$	$1/2^+$		
	Λ	$1/2^+$		
	Ξ	$1/2^+$		
1 + 8		$1/2^-$		
3	$T_c([ud][us]\bar{c})$	$1/2^-$	2785(46)	
	$T_c^*([ud][us]\bar{c})$	$1/2^+$	3243(66)	
	$T_{cs}([us][ds]\bar{c})$	$1/2^-$		
6	$\Theta_c([ud][ud]\bar{c})$	$1/2^+$	2977(109)	$\Theta_c(3099)$
	$\Theta_c^*([ud][ud]\bar{c})$	$1/2^-$	3353(89)	search ?
	$N_c([ud][us]_+\bar{c})$	$1/2^+$	3180(69)	search ?
	$\Xi_c([ds][ds]\bar{c})$	$1/2^+$	3650(95)	search ?

Chiu-Hsieh (2)

hep-lat/0403020, 0404007

 They claim:

the lower-lying pentaquark; **positive parity**

Chiu-Hsieh (2)

hep-lat/0403020, 0404007

- They claim:
 - ☑ the lower-lying pentaquark; **positive parity**
 - In contradiction to Wilson results (CFKK/Sasaki)

Chiu-Hsieh (2)

hep-lat/0403020, 0404007



They claim:



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- **The discrepancy is due to the lattice artifacts of the Wilson fermion**

Chiu-Hsieh (2)

hep-lat/0403020, 0404007



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- Any NK scattering state is not seen**

Chiu-Hsieh (2)

hep-lat/0403020, 0404007



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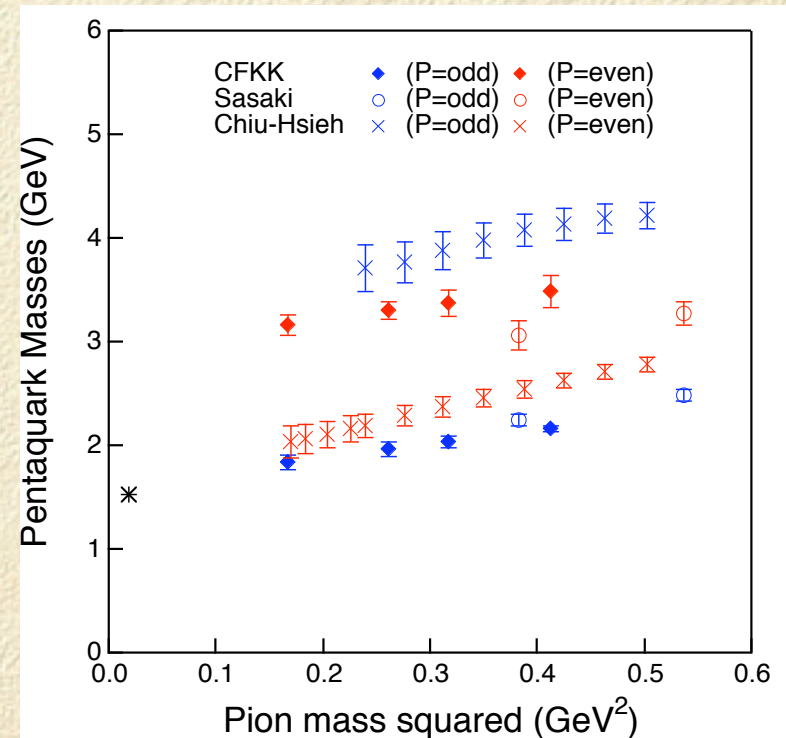
- Strongly support the Jaffe-Wilczek model**
 - **the N(1440) can be identified with a pentaquark**

Comments on Chiu-Hsieh



The lightest pion mass ~ 0.4 GeV

- One would not expect any qualitative difference between Wilson and Chiral fermions in their quark mass range.



Comments on Chiu-Hsieh

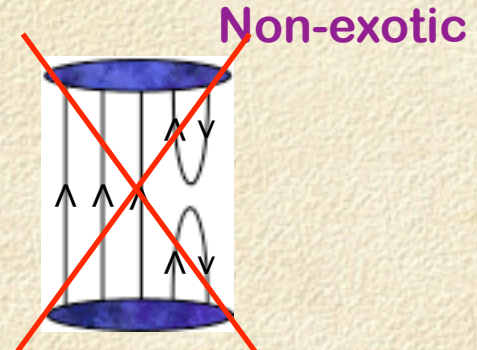
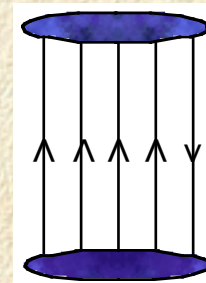
☑ The lightest pion mass ~ 0.4 GeV

- One would not expect any qualitative difference between Wilson and Chiral fermions in their quark mass range.

☑ Really observe the N-type lightest pentaquark?

$$\cancel{M_N([ud][ud]\bar{d}) = 1460(51)\text{MeV} \rightarrow \cancel{N(1440), 1/2^+}}$$

$M_\Theta([ud][ud]\bar{s}) \text{ in } SU(3)_f \text{ limit}$



Kentucky Collab. (1)

Mathur et al., hep-ph/0406196

- Overlap fermions (Quench) + Iwasaki gauge action
 - $L \sim 2.4$ and 3.2 fm ($12^3 \times 40$, $16^3 \times 40$), $a = 0.20$ fm
 - m_π down to 0.18 GeV from 1.0 GeV
 - $J = 1/2$, $l = 0, 1$, parity projection
 - Simple minded operator as N times K
 - Sequential empirical Bayes method
 - Check the volume dependence of spectral weight.

Kentucky Collab. (2)



They claim:

Mathur et al., hep-ph/0406196



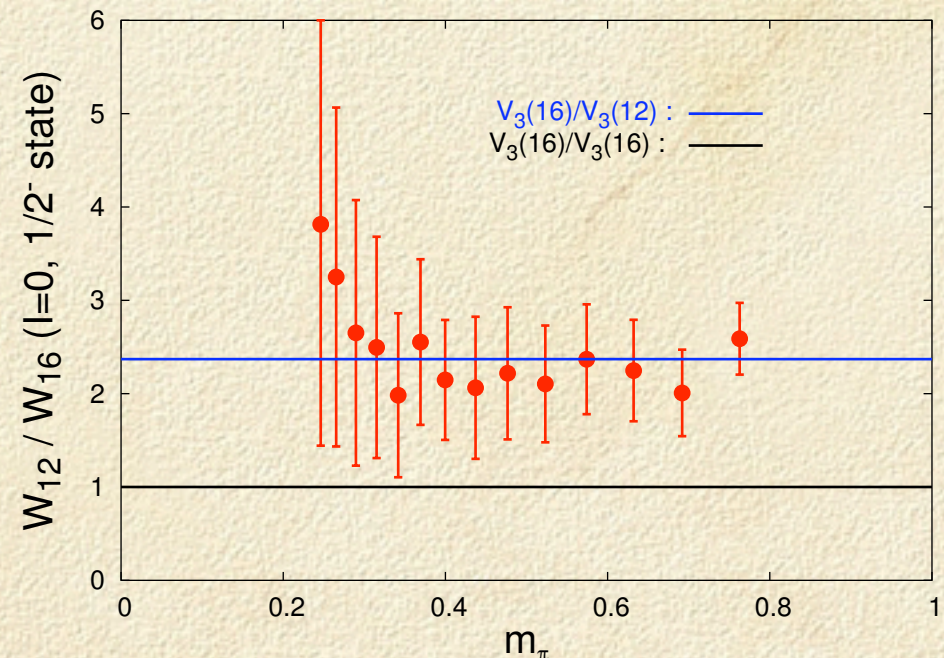
See only KN scattering state in negative parity

- Ground state in either parity channel has a characteristic volume dependence on the spectral weight.

$$|N, \vec{p}, s\rangle_{\text{latt}} \propto \sqrt{\frac{M}{VE}} |N, \vec{p}, s\rangle_{\text{cont}}$$

The spectral weight should have
 $1/V$ dependence for two particles

$$W_{12}/W_{16} = 16^3/12^3 = 2.37$$



Kentucky Collab. (2)



They claim:

Mathur et al., hep-ph/0406196

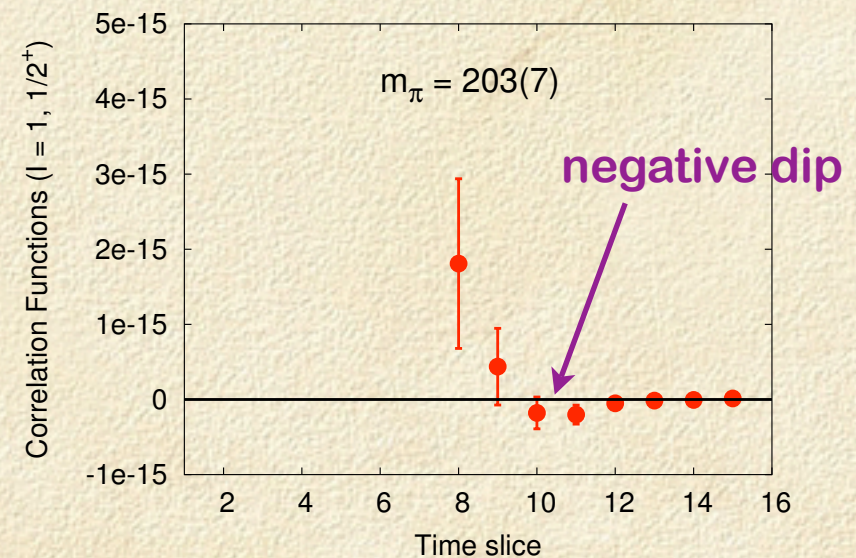


See only KN scattering state in negative parity

- Ground state in either parity channel has a characteristic volume dependence on the spectral weight.



Confirm the ghost $KN\eta'$ state in positive parity channel at $M_\pi < 266$ MeV



Kentucky Collab. (3)

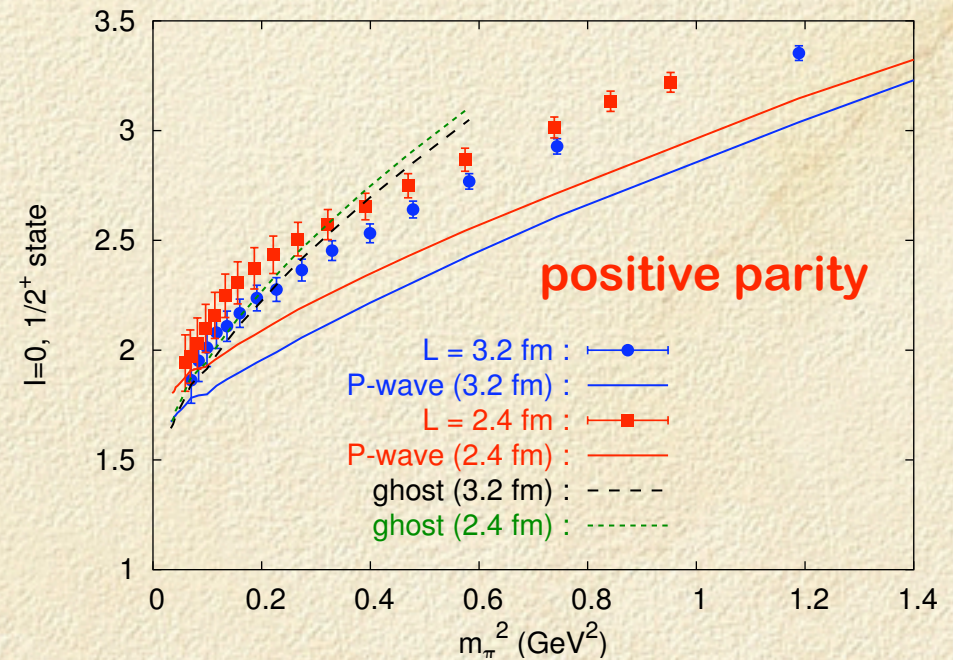
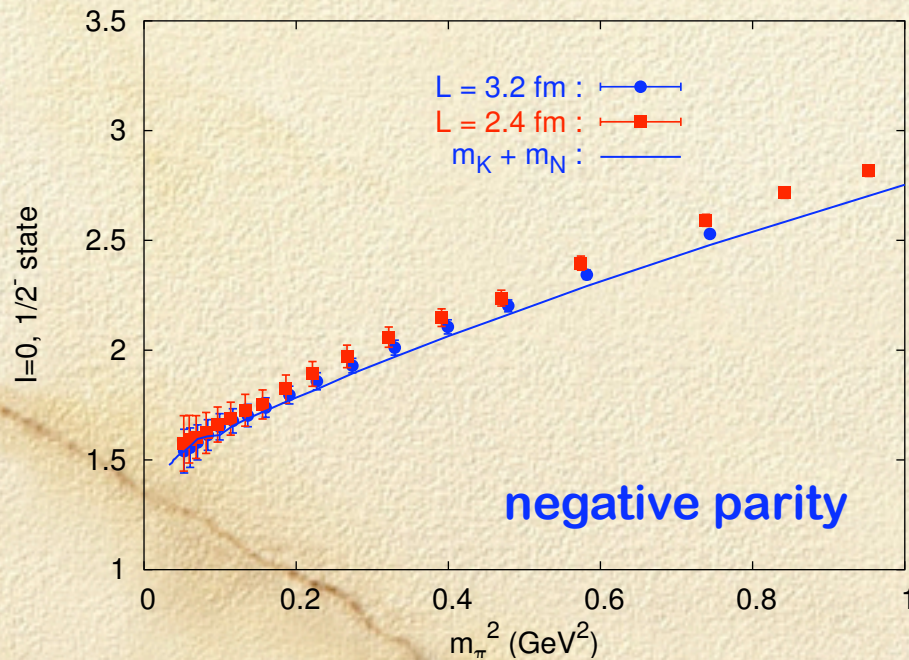


They claim:

Mathur et al., hep-ph/0406196



No sign of pentaquark signal in either parity channel at present.



Sigaev-Jahn-Negele (1)



Eight possible local sources: **extention of Sasaki's proposal**

$$\mathcal{O} = \varepsilon_{abc}\varepsilon_{aef}\varepsilon_{bgh}[u_e^T C\Gamma_1 d_f][u_g^T C\Gamma_2 d_h]\Gamma_3 C\bar{s}_c^T$$

I=0

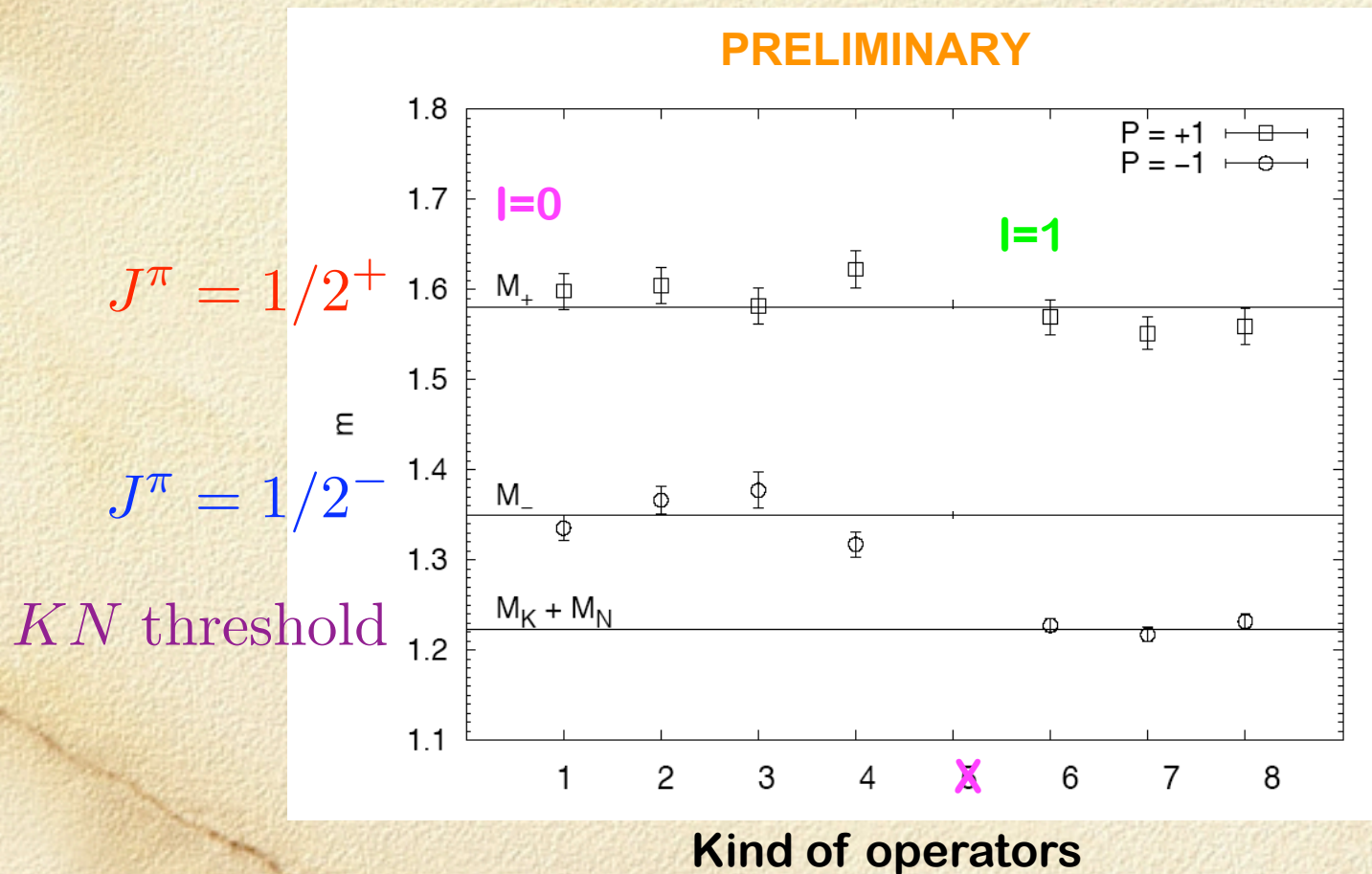
I=1

	Γ_1	Γ_2	Γ_3	Isospin	Lorentz
1	1	γ_5	1	S	P
2	γ_5	$\gamma_5\gamma_\mu$	$\gamma_5\gamma_\mu$	S	V
3	1	$\gamma_5\gamma_\mu$	γ_μ	S	A
4	$\gamma_5\gamma_\mu$	$\gamma_5\gamma_\nu$	$\varepsilon_{\mu\nu\rho\lambda}\sigma_{\rho\lambda}$	S	T
5	$\gamma_\nu\tau_n$	$\sigma_{\mu\nu}\tau_n$	$\gamma_5\gamma_\mu$	V	V
6	$\gamma_\mu\tau_n$	$\varepsilon_{\mu\nu\rho\lambda}\sigma_{\nu\lambda}\tau_n$	γ_μ	V	A
7	$\gamma_\mu\tau_n$	$\gamma_\nu\tau_n$	$\varepsilon_{\mu\nu\rho\lambda}\sigma_{\rho\lambda}$	V	T
8	$\sigma_{\mu\nu}\tau_n$	$\sigma_{\nu\lambda}\tau_n$	$\varepsilon_{\mu\nu\rho\lambda}\sigma_{\rho\lambda}$	V	T

⇒ Compute 4x4 cross correlator in each isospin channel.

Sigaev-Jahn-Negele (2)

The results are currently from double-exp. fits to each operator.



$\beta=6.0$ $16^3 \times 32$
Wilson (Quench)

Comparison of lattice results

Q1: Observe pentaquarks in lattice QCD?

Q2: Which parity is assigned to the lowest state of the pentaquark?

Q3: Find anti-charmed pentaquark as bound state or near threshold?

	A1	A2	A3	Op.
Csikor et al.	YES	negative	–	color variant of KN
Sasaki	YES	negative	NO	diquark-diquark-antiquark
Kentucky	NO	not positive	–	simple KN
Chiu-Hsieh	YES	positive	YES	diquark-diquark-antiquark
MIT	YES	negative	–	diquark-diquark-antiquark

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Summary (1)

All results should be regarded as **exploratory** at present.

- ★ To confirm the presence of the pentaquarks
 - Try some other possible operators
 - The $M \times M$ correlation matrix analysis is required to separate the KN scattering state

- ★ To disentangle the pentaquark signal from the KN scattering states
 - Identify the lowest few scattering states in both parity channels
 - Verify the volume dependence

Detail studies are in progress by CFKK, Sasaki, Chiu-Hsieh, Kentucky, MIT

Summary (2)

 The following questions still remain open:



Does the spectrum of QCD possess the $\Theta^+(1540)$?



What is **spin** and **parity** of the $\Theta^+(1540)$?



Are there other pentaquark baryons?

- other member of the anti-decuplet, especially $\Xi_{3/2}$
- the **charm (bottom)** pentaquark
- the **spin-orbit partner** of the $\Theta^+(1540)$; $J=3/2$

There are many exciting issues to be explored.

Many thanks to

- Zoltan Fodor & Sandor Katz
- Keh-Fei Liu & Nilmani Mathur
- Ting-Wai Chiu
- John Negele

for their correspondences