Preparing for Discovery A Bard's-Eye View

#### Stephen Mrenna

Computing Division Fermilab

Wine & Cheese, March 2006





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### **SA Perspectives**

### The Collider Calamity

For docades, the big game of American science have been the U.S. Department of Energy's metricle colliders, which investigate the nature of matter by accelerating submonit, purifies and smoothing that cognier. Colliders as the Ferrari Satiurof Academics Labornatory (Fermidal), Stenford Lanen Accelerator Conte-(SLAC) and Bookhawen Phanimal Laboratory bave discovered exons, particles such as the uppartic and



FERMILAR, home of the Towatron.

rescaled pleasurement lan him an new laws of physics, but this years Ar refraint emerging so, ble so many urbers, is now mexit g wersaus. While the furoparaus of physics build new partice stockerators, the U.S. s poised to shurt down is prentice colliders at Fermilab and SLAC over the new few years. And funding for Breeckhaver's Relative stic Heavy Ion Collider (RHC) is so right at the lab cold not have run

its full slate of experiments this year without \$15 million raised by a New York billionaire.

The sad story began in 1983, when Congress can-

provides gated a deviate called PTCV that we understandthe decay of B mesons emanating from arbitrations at the Tevanco. FTeV employed such such as demonstratic Tevanco. FTeV employed such such as the endowing duration of the second performance of Emergy semeched FTeV. Whiteout that experiments of Emergy semeched FTeV. Whiteout that experiments the Tevancer mutuing after the FTeV employer in the key the Tevancer mutuing after the FTeV employer in the key the Tevancer mutuing after the FTeV employer in the key the Tevancer mutuing after the FTeV employer in the key the National Sciences Frontaletion reasoning. While are used performent or lead RSVF that would have used Broudd averts accelerator to investigate tare particle decays that could not be observed at the LHC.

Bies des depriving researchers of putertial closcoveries, tasse curs timeaten to make the U.S. less econamically compactive. The development of lagt energy actelerators has led to advances in medicine and electronics, and American expertise in this field will wither if the U.S. ceases to build and operate colliders. Moreover, although American scientists will participate in the research at the LTIC, the Europeans will get most of the educational benefits of the textility, which will inspire and trait the rear generation of



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#### Preparing for Discovery

The New Hork Eimes 1315 Physicists Report Failure In Search for Supersymmetry The negative result illustrates

#### 'God particle' may have been seen

By Paul Rincon BBC News Online science staff

A scientist says one of the most sought after particles in physics - the Higgs boson - may have been found, but the evidence is still relatively weak.

Peter Renton, of the University of Oxford, says the particle may have been detected by



... finally

Once produced, the Higgs boson would decay very quickly

researchers at an atom-smashing facility in Switzerland.

The Higgs boson explains why all other particles have mass and is fundamental to a complete understanding of matter.



### EXPERIMENTAL EVIDENCE FOR MORE DIMENSIONS REPORTED

#### Gordon L. Kane May 2011

The worldview of physicists working on unification theories has been changing rapidly recently. That change culminated in March, at the 46th annual Recontres de Moriond conference in Les Arcs, France, with the announcement of some startling data from CERN's Large Hadron Collider (LHC).

More than two hundred years ago, Charles Augustin Coulomb showed that the electrical force had the same form as the gravitational ory. Because the work was well ahead of its time, and because of World War II, Klein's insight went largely unnoticed. See L. O'Raifeartaigh, *The Dawning of Gauge Theory*, Princeton University Press, 1977.)

The fields of the higher-dimensional theory were the gravitational tensor field, the electromagnetic vector potential field and a scalar field. Of course, the theories of electricity and magnetism were unified without extra dimensions by Maxwell, and the

#### Shakespeare's Writing Method

- Develop a large vocabulary
- Play with words
- Invent new words and phrases
- Develop the common touch
- Read great literature
- Study the great orators, actors and the popular
- Live with passion
- Write, write, write!!!

# How much does the $t\bar{t}$ cross section change from TeV to LHC?



■ 500×

[Kidonakis]



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# How much does the $t\bar{t}$ cross section change from TeV to LHC?



[Kidonakis]



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- How much does the  $\tilde{\chi}^+ \tilde{\chi}^- (m_{\chi} = 200 \text{ GeV})$  cross section change from TeV to LHC?
- 10×
- 100×
- 500×

[Pythia]



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# How much does the $\tilde{\chi}^+ \tilde{\chi}^- (m_{\chi} = 200 \text{ GeV})$ cross section change from TeV to LHC?

#### ■ 10×





[Pythia]



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# How much does the W4j cross section change from TeV to LHC?



[MadEvent,  $k_T > 20 \text{ GeV}$ ]



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# How much does the W4j cross section change from TeV to LHC?



[MadEvent,  $k_T > 20 \text{ GeV}$ ]



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#### Heavy Colored Objects Large Kinematic Reach





- LHC phenomenology begins with rediscovering the Standard Model
- The path starts at the Tevatron

#### Understanding W+Jets is Important

- Signature Wbb + X is common to unconfirmed Standard Model processes and many new physics processes
- we "know" that Standard Model top is there

 $\mathsf{Top} \equiv \mathsf{Data} - \mathsf{Not-Top}$ 

- As JES uncertainty is reduced (CDF m<sub>t</sub>), understanding of Not-Top sets/limits understanding of Top
- Advanced (i.e. NN, DT) search techniques exploit differences in many kinematic variables
- Not-Top challenges our tools

Better tools  $\Rightarrow$  more challenging questions



#### New Physics Warm-Up

- current state of single-Top is where we will be at the LHC with a few quality fb<sup>-1</sup>
- the size of other NP signals
- it is a playground for new analysis techniques
- it challenges our tools

Not specific to NN analyses: may be more sensitive Many (11) Kinematic Variables

	Signal-Background Pairs			
	tb		tqb	
	Wbb	tŦ	Wbb	tĒ
Individual object kinematics				
$p_T(jetl_{tagged})$	$\checkmark$	$\checkmark$	$\checkmark$	_
$p_T(jetl_{untagged})$	_	_	$\sim$	$\checkmark$
$p_T(jet2_{untagged})$	_	_	_	$\sim$
$p_T(jet1_{nonbest})$	$\sim$	$\sim$	_	_
$p_T(jet2_{nonbest})$	$\checkmark$	$\sim$	_	-
Global event kinematics				
$M_T$ (jet1, jet2)		_	_	_
$p_T$ (jet1, jet2)	√.		√.	_
M(alljets)	$\checkmark$	$\sim$	√.	$\checkmark$
$H_T(\text{alljets})$	_	_	$\sim$	_
$M(\text{alljets} - \text{jet1}_{tagged})$	_	_	_	√
$H(alljets - jet1_{tagged})$	_	$\sim$	_	√
$H_T(alljets - jet1_{tagged})$	_	_	_	$\checkmark$
$p_T(alljets - jet1_{tagged})$	_	$\checkmark$	_	$\checkmark$
$M(\text{alljets} - \text{jet}_{best})$	_	$\checkmark$	_	_
$H(alljets - jet_{best})$	_	$\checkmark$	_	_
$H_T(alljets - jet_{best})$	-	$\checkmark$	-	_
$M(top_{tagged}) = M(W, jetl_{tagged})$	$\checkmark$	$\sim$	$\checkmark$	$\checkmark$
$M(top_{best}) = M(W, jet_{best})$	$\checkmark$	_	-	_
$\sqrt{\hat{s}}$	$\checkmark$	_	$\checkmark$	$\checkmark$
Angular variables				
$\Delta R(\text{jet1}, \text{jet2})$	$\checkmark$	—	$\checkmark$	—
$Q(\text{lepton}) \times \eta(\text{jet1}_{\text{untagged}})$	_	_	$\checkmark$	$\checkmark$
$\cos(\text{lepton}, Q(\text{lepton}) \times z)_{topbest}$	$\checkmark$	_	_	_
cos(lepton, jet1 <sub>untagged</sub> ) <sub>toptagged</sub>	_	—	$\checkmark$	—
cos(alljets, jet1 <sub>tagged</sub> ) <sub>alljets</sub>	_	—	$\checkmark$	$\checkmark$
cos(alljets, jet <sub>nonbest</sub> ) <sub>all jets</sub>	_	$\checkmark$	_	-



#### **Top Background Summary**





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#### Method 2

Monte Carlo ratio R = (W + b - jets)/(W + jets)• Common factors cancel Measure W + jets (no b-tag) data(W + b - jets) =  $R \times data(W + jets)$  Wcj/Wbb from Monte Carlo • Several R's



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High Multiplicity Tree Graph



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Tree Graph + Parton Shower



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Lower Multiplicity Tree Graph



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#### Lower Multiplicity NLO Graph



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#### Tree Graphs + Parton Showers





#### Tree Graphs + Parton Showers









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 $k_T^{1} > k_T^{2} > k_T^{3} > k_T^{4} > k_T^{cut}$ 



 $k_{\tau}^{1} > k_{\tau}^{2} > k_{\tau}^{3} > k_{\tau}^{4} > k_{\tau}^{cut}$ 





 $k_T^{1} > k_T^{2} > k_T^{3} > k_T^{4} > k_T^{cut}$ 



 $k_T^{1} > k_T^{2} > k_T^{3} > k_T^{4} > k_T^{cut}$ 











### Clever Matching of Tree Graphs and Parton Showers

#### Make Better Predictions



Stephen Mrenna Preparing for Discovery

### Clever Matching of Tree Graphs and Parton Showers

#### Address Uncertainty


## Cross check on Run2 data



## Kinematic comparisons with Run2 data





### ме-ре decomposition



# To understand the data, look at the Vista of final states

Final State	Chi2	data	bkg		
1b3j1pmiss_sumPt400+ [73]	9.0	451	374.5 +	- 18	( pyth_jj_200 = :
2b1e+2j [-]	8.0	15	6.5 +	- 1.9	( ttop0z = 2.3 ,
2j_sumPt0-400 [161]	6.0	69704	67013.6 +	- 1171.2	( pyth_jj_018 = 3
2j2mu+1pmiss [-]	-5.0	2	12.2 +	- 3	( mad_mu+mu-jj =
1b2e+2j [-]	5.0	9	3.9 +	- 1.5	( mrenna_e+e-jjj
1j1ph1pmiss [5]	4.0	2591	2470.1 +	- 37.7	( pyth_pj_045 = 7
2j1mu+1ph [-]	4.0	11	11.2 +	- 2.2	( mrenna_mu+mu-j
1e+1j1mu+ [-]	4.0	13	6.6 +	- 2.1	( ztop5i = 3.4 ,
1e+2j1ph [-]	4.0	31	20.9 +	- 2.7	( mad_aajj = 6.3
3j2mu+ [-]	4.0	34	23.2 +	- 2.7	( mrenna_mu+mu-jj
2b2j1pmiss_sumPt400+ [-]	-3.0	17	30.4 +	- 4.2	( pyth_jj_200 = :
1b2j_sumPt400+ [229]	3.0	4669	4518.6 +	- 72.7	( pyth_jj_200 = 2
4j_sumPt0-400 [253]	-3.0	2611	2736.9 +	- 42.3	( pyth_jj_040 = :
2b1j1ph1pmiss [-]	3.0	6	2.7 +	- 1.5	$( pyth_{jj}200 = 0$
1b1j1mu+ [-]	3.0	67	53.8 +	- 4.3	( pyth_jj_018 = :
1j1ph [277]	3.0	31738	31149.8 +	- 352.1	( pyth_pj_045 = :
1e+1mu+ [-]	3.0	66	53.5 +	- 3.2	( ztop5i = 38.8
4j1mu+ [-]	3.0	73	61.3 +	- 2.6	( pyth_jj_040 = :
5j [269]	3.0	448	406 +	- 14.5	( pyth_jj_040 = :
1b5j [-]	3.0	8	8.9 +	- 1.7	( pyth_jj_060 = :
1b1j1pmiss_sumPt0-400 [-]	2.0	120	104 +	- 7.2	( pyth_jj_040 = 3
2j1pmiss_sumPt0-400 [37]	2.0	2381	2281.2 +	- 73.9	( pyth_jj_018 = :

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Final State	Chi2	data	bkg	
1b2e+2j [-]	5.0	9	3.9 +-	1.5
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# Give a complete description of the Standard Model with the best tools

## Patriot



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#### Model-Independent and Quasi-Model-Independent Search for New Physics at CDF

Georgios Choudalakis,<br/>\* Khaldoun Makhoul, † Markus Klute, ‡ Conor Henderson,<br/>§ and Bruce Knuteson ¶ $_{MIT}$ 

> Ray Culbertson\*\* FNAL

#### CDF Collaboration<sup>††</sup> (Dated: February 1, 2006)

Data collected in Run II of the Farmiab Towarton are searched for indications of new electroweak scale physics. Rather than focusing on particular more physics scansarios, CDF data are analyzed for discrepancies with the Standard Model prediction. A model-independent approach considers the gress features of the data, and is sensitive to new large cross section physics. A quest-modelindependent approach emphasizes the high-pr table, and is particularly sensitive to new electroweak scale physics. This global search for new physics in  $\approx 600$  pb<sup>-1</sup> of *pp* collisions at  $\sqrt{s} = 1.96$  TeV reveals

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## Checking Assumptions Is description of Underlying Event universal?



- The first New Physics to find is the Standard Model
- Need complete description of most important processes
- Understanding comes from looking at consistency of full dataset
- Then, how do we find New Physics?



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e<sup>+</sup> e<sup>-</sup> bb Final State



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## Generate Process Code On-Line

Quarks: d u s c b t d~ u~ s~ c~ b~ t~

Leptons: e- mu- ta- ve vm vt e+ mu+ ta+ ve~ vm~ vt~

Bosons: A Z W+W-hg

 Special: P j (sums over d u s c d ~ u ~ s ~ c ~ g)

 Process:
 PP > W+ > e+ ve jijj

 Submit
 EXAMPLES

 Max QCD Order:
 4

 Max QED Order:
 2

To improve our web services we now request that you register. Registration is quick and free. You may register for a password by clicking <u>here</u>



Generic Particles and Vertices

$$\mathcal{L}_{\rm FFV} = \bar{f}' \gamma^{\mu} \left( \mathsf{G}(1) \frac{1 - \gamma_5}{2} + \mathsf{G}(2) \frac{1 + \gamma_5}{2} \right) f V_{\mu}^*$$

$$\mathcal{L}_{\rm FFS} = \bar{f}' \left( {\rm GC}(1) \frac{1-\gamma_5}{2} + {\rm GC}(2) \frac{1+\gamma_5}{2} \right) f S^*$$

$$\begin{split} \mathcal{L}_{\text{VVV}} &= -i \mathbf{G} \quad \left\{ \begin{array}{l} (\partial_{\mu} V_{1\nu}^{*}) (V_{2}^{\mu*} V_{3}^{\nu*} - V_{2}^{\nu*} V_{3}^{\mu*}) \\ &+ (\partial_{\mu} V_{2\nu}^{*}) (V_{3}^{\mu*} V_{1}^{\nu*} - V_{3}^{\nu*} V_{1}^{\mu*}) \\ &+ (\partial_{\mu} V_{3\nu}^{*}) (V_{1}^{\mu*} V_{2}^{\nu*} - V_{1}^{\nu*} V_{2}^{\mu*}) \right\} \end{split}$$

$$\mathcal{L}_{\rm VVS} = {\rm G} V_1^{\mu*} V_{2\mu}^* S^*$$

$$\mathcal{L}_{\text{SSS}} = \text{G} \, S_1^* S_2^* S_3^* \qquad \qquad \mathcal{L}_{\text{VSS}} = i \text{G} V_{\mu}^* S_2^* \stackrel{\leftrightarrow \mu}{\partial} S_1^*$$

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#### A general search for new phenomena in *ep* scattering at HERA

H1 Collaboration

#### Abstract

A model-independent search for deviations from the Standard Model prediction is performed in  $\epsilon^{-1}p$  collisions at HERA using HI data corresponding to an integrated luminosity of 117 pb<sup>-1</sup>. For the first time all event topologies involving isolated electrons, photom, muons, neutrinos and jets with high transverse momenta are investigated in a single analysis. Events are assigned to exclusive classes according to their final state. A statistical algorithm is developed to search for deviations from the Standard Model in the distributions of the scalar sum of transverse momenta or invariant mass of final state particles and to quantify their significance. A good agreement with the Standard Model in prediction is observed in most of the event classes. The most significant deviation is found for a topology containing an isolated muon, missing transverse momentum and a jet, consistent with a previously reported observation.

# arXiv:hep-ex/0408044 v1 12 Aug 2004



Figure 1: The data and the SM expectation for all event classes with a SM expectation greater than 0.01 events. The analysed data sample corresponds to an integrated luminosity of 117 pb<sup>-1</sup>. The error bands on the predictions include model uncertainties and experimental systematic errors added in quadrature.



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## ■ ALEPH and L3 analyses underway



## Quaero: DØ, hep-ex/0106039

Process	$\epsilon_{\rm sig}$	ĥ	$N_{ m data}$	$\sigma^{95\%}  imes \mathcal{B}$
$WW  ightarrow e\mu  ot\!\!\!/ _T$	0.14	$19.0\pm4.0$	23	1.1 pb
ZZ  ightarrow ee 2j	0.12	$19.7 \pm  4.1$	19	0.8 pb
$tar{t}  ightarrow e  ot\!$	0.13	$3.1\pm~0.9$	8	0.8 pb
$tar{t}  o e \mu  ot\!$	0.14	$0.6\pm\ 0.2$	2	0.4 pb
$h_{175} \rightarrow WW \rightarrow e \not\!\!\! E_T 2j$	0.02	$29.6\pm6.5$	32	11.0 pb
$h_{200}  ightarrow WW  ightarrow e  ot\!$	0.07	$66.0 \pm 13.8$	69	4.4 pb
$h_{225}  ightarrow WW  ightarrow e  ot\!$	0.06	$43.1\pm9.2$	44	3.6 pb
$h_{200}  ightarrow ZZ  ightarrow ee 2j$	0.15	$17.9\pm3.7$	15	0.6 pb
$h_{225}  ightarrow ZZ  ightarrow ee 2j$	0.15	$18.8\pm3.8$	12	0.4 pb
$h_{250}  ightarrow ZZ  ightarrow ee 2j$	0.17	$18.1\pm~3.7$	18	0.6 pb
$W'_{200}  ightarrow WZ  ightarrow e  ot\!$	0.05	$27.7\pm6.3$	29	3.4 pb
$W'_{350}  ightarrow WZ  ightarrow e  ot\!$	0.23	$22.7\pm5.2$	27	0.7 pb
$W'_{500}  ightarrow WZ  ightarrow e  ot\!$	0.26	$2.1\pm~0.8$	2	0.2 pb
$Z'_{350}  ightarrow t \overline{t}  ightarrow e  ot\!$	0.11	$18.7\pm4.0$	20	1.1 pb
$Z'_{450}  ightarrow t \overline{t}  ightarrow e  ot\!$	0.14	$18.7\pm4.0$	20	0.9 pb
$Z_{550}^\prime  ightarrow t \overline{t}  ightarrow e  ot\!$	0.14	$3.8\pm1.0$	2	0.3 pb
$Wh_{115}  ightarrow e  ot\!$	0.08	$37.3\pm8.2$	32	2.0 pb
$Zh_{115}  ightarrow ee 2j$	0.20	$19.5\pm4.1$	25	0.8 pb
$LQ_{225}\overline{LQ}_{225} ightarrow$ ee 2 $j$	0.33	$0.3\pm0.1$	0	0.07 pb
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#### **BARD: Interpreting New Frontier Energy Collider Physics**

Bruce Knuteson<sup>\*</sup> MIT

Stephen Mrenna<sup>†</sup> FNAL

No systematic procedure currently exists for inferring the underlying physics from discrepancies observed in high energy collider data. We present BARD, an algorithm designed to facilitate the process of model construction at the energy frontier. Top-down scans of model parameter space are discarded in favor of bottom-up diagrammatic explanations of particular discrepancies, an explanation space that can be exhaustively searched and conveniently tested with existing analysis tools.



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## New Particles

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SSS	SSS	S	d	npm(1)	npW	S	Xsss	31
ssf	ssf~	f	S	npm(2)	npW	s	Xssf	32
SZS	SZS	s	d	npm(3)	npW	s	Xszs	33
szf	szf~	f	S	npm(4)	npW	s	Xszf	34
sas	sas~	S	d	npm(5)	npW	S	Xsas	35
saf	saf~	f	S	npm(6)	npW	S	Xsaf	36
sbs	sbs~	S	d	npm(7)	npW	S	Xsbs	37
sbf	sbf~	f	S	npm(8)	npW	S	Xsbf	38
SCS	scs~	S	d	npm(9)	npW	S	Xscs	39
scf	scf~	f	S	npm(10)	npW	S	Xscf	40
• • •								
OSS	OSS	S	d	npm(51)	npW	0	Xoss	81
osf	osf~	f	S	npm(52)	npW	0	Xosf	82
ozs	ozs	S	d	npm(53)	npW	0	Xozs	83
ozf	ozf~	f	S	npm(54)	npW	0	Xozf	84
SSV	SSV	v	W	npm(55)	npW	S	Xssv	85
osv	osv	v	w	npm(56)	npW	S	Xosv	86
scv	scv~	v	W	npm(57)	npW	S	Xscv	87



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## New Interactions

a	sas	sas~	np_coupl_c(453)	QNP
a	sbs	sbs~	np_coupl_c(455)	QNP
a	SCS	scs~	np_coupl_c(457)	QNP
• • •				
b	u	scs~	np_coupl_cLR(261)	QNP
b	u	scv~	np_coupl_rLR(41)	QNP
d	b	OSS	np_coupl_cLR(408)	QNP
d	b	osv	np_coupl_rLR(27)	QNP
d	Ъ	ozs	np_coupl_cLR(418)	QNP
d	Ъ	SSS	np_coupl_cLR(183)	QNP
• • •				
Z	tss	tzs~	np_coupl_c(466)	QNP
Z	tzs	tzs~	np_coupl_c(474)	QNP
Z	w+	scs~	np_coupl_c(438)	QNP

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- **1** NP must couple to the initial state or an annihilation particle
- **2** SM particles must couple to either the initial or final state
- 3 No more than n NP particles can appear in a given diagram
- 4 NP particles can appear twice only in separate chains
- 5 • •

## **Pmiss Final States**







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Stephen Mrenna Preparing for Discovery

Compose:	New Physic	s <2>						
<u>Eile</u>	t <u>V</u> iew	Options T	ools <u>W</u> in	dow <u>H</u> elp				
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-	To:	🔝 witten@	Plas.princet	on.edu				
Subject:	New Physi	cs						

Dear Prof Witten,

```
I have analyzed the excesses observed in the data,
and have determined the following stories, ranked in descending
log likelihood:
```

```
Story 1
Particles (SU(3),Q,type)
                                      t4/3f
          555
                        OSV
Mass (GeV)
          251+/-12
                        1043+/-102
                                      341+/-73
Interactions
          sss b b
                                      sss t4/3f t4/3f~ ....
                        SSS WH W-
Coupling
          .1+/-.03
                        .3+/-.1
                                      1.0+/-0.3 ....
Story 2
Could you please tell us the correct string vacuum?
Sincerely,
```

the Bard

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Table 1. Quantum numbers of scalar and vector leptoquarks with SU(3) x SU(2) x U(1) invariant couplings to quark-lepton pairs (Y =  $Q_{em} - T_{\chi}$ ).

		Spin	F = 3B + L	SU(3) <sub>C</sub>	SU(2)W	U(1)Y
	s,	0	-2	3*	1	$\frac{1}{3}$
	ŝ,	0	-2	3*	1	4 3
	s,	0	-2	3*	3	$\frac{1}{3}$
	· v.	1	-2	3*	2	5
	ĩ,	1	-2	3*	2	- 1/6
	Rs	0	0	3	2	7
	$\widetilde{R}_{1}$	0	0	3	2	1 6
	υ,	1	0	3	1	23
	ũ,	1	0	3	1	53
	Ũ,	1	0	3	3	23
		1				

Büchmuller et al.

Table 2.	Couplings	of scalar	and	vector	leptoqua	rks to	quark	-lepton
	pairs. The	subscripts	L,R	of the	couplings	refer	to the	lepton
	chirality.							

	F = -2	, scal	ars	F ≖ -2,	vectors	
channe1	s,	ĩ,	<u>s</u> ;	V.	ĩ,	
eī, u	911.R	22	-93L	91R	Ĩ	
vid	-916	-	-936	9 <sub>31</sub>	-	
e LR d	100	91R	- 12 gsL	92L,R		
νĽu	12.0	-	12 g.	-	Ĩ.	
	F = 0.	vector	s	F = 0,	scalars	
channe 1	U,	ũ,	Ū,	Rz	Ĩ,	
e <sub>LR</sub> d <sub>b</sub>	his,e	-	-hau	-h <sub>28</sub>	ñar	
YL UT	h.L	-	hat	hau	-	
e, R ut	-	ñ	12 has	haL,R	-	
Y, dp		-	V2 has	-	ñ	

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	X	r,s,h)				$X_{(r,s,h)}$			
	l	q		ν	e	и	d		
ē	(1, 1, 0) (1, 3, 0)	$(3, 1, \frac{2}{3})$	v	(1, 1, 0)	(1, 1, -1)	$(3, 1, \frac{2}{3})$	$(3, 1, -\frac{1}{3})$		
$\bar{q}$	$(\bar{3}, 1, -\frac{2}{3})$	(1,1,0)(1,3,0)(8,1,0)	ē	(1, 1, 1)	(1, 1, 0)	$(3, 1, \frac{5}{3})$	$(3, 1, \frac{2}{3})$		
$\vec{v}^c$	$(1, 2, -\frac{1}{2})$	$(3, 2, \frac{1}{6})^{B}$	ū	$(\bar{3}, 1, -\frac{2}{3})$	$(\bar{3}, 1, -\frac{5}{3})$	(1, 1, 0)(8, 1, 0)	(1, 1, -1)		
$\overline{e}^{c}$	$(1, 2, -\frac{3}{2})$	$(3, 2, -\frac{5}{6})^{B}$	ā	$(\bar{3}, 1, \frac{1}{3})$	$(\bar{3}, 1, -\frac{2}{3})$	(1, 1, 1)	(1, 1, 0)  (8, 1, 0)		
$\bar{u}^c$	$(3, 2, \frac{1}{6})^{B}$	$(\bar{3}, 2, \frac{5}{6})^{B}$							
$\bar{d}^{c}$	$(3, 2, -\frac{5}{6})^B$	$\left(\bar{3},2,-\frac{1}{6}\right)^{\mathrm{B}}$							
$\Phi_{(r,s,h)}$				$\Phi_{(r,s,h)}$					
	Ł	q		ν	е	u	d		
<i>ē</i> c	$(1, 1, 1)^* (1, 3, 1)$	$(3, 1, -\frac{1}{3})^{B}(3, 3, -\frac{1}{3})$	ē	$(1, 2, \frac{1}{2})$	$(1, 2, -\frac{1}{2})$	$(3, 2, \frac{7}{6})$	$(3, 2, \frac{1}{6})$		
$\bar{q}^{c}$	$\left(\bar{3},1,\tfrac{1}{3}\right)^{\mathrm{B}}\left(\bar{3},3,\tfrac{1}{3}\right)$	$(\overline{3}, 1, \frac{1}{3})^{B} (\overline{3}, 3, \frac{1}{3})^{*}$ $(6, 1, \frac{1}{2})^{*} (6, 3, \frac{1}{2})$	$\bar{q}$	$(\bar{3}, 2, -\frac{1}{6})$	$(\bar{3}, 2, -\frac{7}{6})$	$(1, 2, \frac{1}{2})(8, 2, \frac{1}{2})$	$(1, 2, -\frac{1}{2})(8, 2, -\frac{1}{2})$		
		<i>,</i> , , , , , , , , , , , , , , , , , , ,	$\bar{v}^c$	(1, 1, 0)	(1, 1, -1)	$(3, 1, \frac{2}{3})^{B}$	$(3, 1, -\frac{1}{3})^{B}$		
			$\bar{e}^{c}$	(1, 1, 1)	(1, 1, -2)	$(3, 1, -\frac{1}{3})^{B}$	$(3, 1, -\frac{4}{3})^{B}$		
			ūc	$(3, 1, \frac{2}{3})^{B}$	$(3, 1, -\frac{1}{3})^{B}$	$(\bar{3}, 1, \frac{4}{3})^* (6, 1, \frac{4}{3})$	$(\bar{3}, 1, \frac{1}{3})^{B}(6, 1, \frac{1}{3})$		
			$\bar{d}^{c}$	$(3, 1, -\frac{1}{3})^{B}$	$(3, 1, -\frac{4}{3})^{B}$	$(\bar{3}, 1, \frac{1}{3})^{B}(6, 1, \frac{1}{3})$	$(\bar{3}, 1, -\frac{2}{3})^{*B}(6, 1, -\frac{2}{3})$		

**Table 2.** Standard model representations carried by the scalars and vector bosons when coupling to two light fermions; (r, s, h) denotes a boson of hypercharge h carrying the r and s representations of SU(3) and SU(2), respectively. Entries with a \* vanish when both light fermions belong to the same family; those marked by a  $\frac{6}{5}$  generate baryon-number violating operators.

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### Useful Theorists





Stephen Mrenna Preparing for Discovery

#### Debunking Anomalies Unexpected Consequences



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# The Bard at the LHC

- LHC phenomenology begins with the Standard Model
- Need complete description of most important processes
- Understanding comes from looking at consistency of full dataset
- Discrepancies can and will arise in specific final states
- Bard can write a series of ranked stories to describe each
  - bottom-up
- Can test this on Run2 data

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- Discrepancies can and will arise in specific final states
- Bard can write a series of ranked stories to describe each
  bottom-up
- Can test this on Run2 data
- It works
- No, we haven't found anything · · · yet

# Extra Slides



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Preparing for Discovery

W+4 partons							
TEVA	TRON	LHC					
Graph	Cross Sect(fb)	Graph	Cross Sect(pb)				
Sum	Sum 1035.004		577.948				
ug_e+vedggg	<u>112.250</u>	gu_e+vedggg	<u>89.815</u>				
gux_e-vexdxggg	<u>112.040</u>	ug_e+vedggg	<u>89.603</u>				
uux_e-vexudxgg	<u>112.010</u>	gd_e-vexuggg	<u>45.522</u>				
uux_e+veuxdgg	<u>111.900</u>	dg_e-vexuggg	<u>45.342</u>				
dux_e-vexddxgg	46.423	uu_e+veudgg	<u>34.174</u>				
udx_e+veuuxgg	46.388	dxg_e+veuxggg	15.346				
dux_e-vexuuxgg	46.349	gdx_e+veuxggg	<u>15.341</u>				
udx_e+veddxgg	<u>46.330</u>	uxg_e-vexdxggg	10.868				
gdx_e+veuxggg	<u>40.234</u>	gux_e-vexdxggg	<u>10.866</u>				
dg_e-vexuggg	<u>40.122</u>	gg_e+veuxdgg	<u>9.920</u>				
udx_e+vegggg	30.906	gg_e+vescxgg	<u>9.907</u>				
dux_e-vexgggg	30.867	gg_e-vexsxcgg	<u>9.907</u>				
ddx_e-vexudxgg	<u>15.189</u>	gg_e-vexudxgg	<u>9.842</u>				
ddx_e+veuxdgg	<u>15.171</u>	du_e+veddgg	<u>8.903</u>				



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#### MCFM vs MEPS





Matched Datasets have consistently steeper slopes (note: MCFM steeper than LO)

Truncated Datasets contain only  $Wb\bar{b} + Wb\bar{b}j$ 

Slopes more consistent with MCFM

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#### Quaero Automated Searches Bruce Knuteson

#### Algorithm

Signal events

### 1) The signal Monte Carlo is processed

- (events are generated using Pythia, if requested)
- events are smeared with a fast detector simulation
- selection criteria are applied for desired final state
- particle identification efficiencies are considered

#### This gives

- total number of expected signal events in final state
- Monte Carlo signal events as they would look in the detector



- 2) An optimal region is chosen in the variables provided
  - a) Estimate signal and background densities using kernels



1) place "bumps of probability" around each Monte Carlo point 2) sum these bumps into a continuous distribution

$$p(x) = \sum_{i=1}^{N} \text{gauss}(x - x_i)$$

The multivariate generalization is immediate

b) Define a *discriminant* 

$$D(x) = \frac{p(x \mid s)}{p(x \mid s) + p(x \mid b)}$$

and choose a cut on D(x) that minimizes

the 95% CL cross section limit you would expect to set assuming the data contains no signal.

We call 1/this quantity the "sensitivity"

Note that so far we have made no use of the data



3) Comparing number of observed events in the data to expected bkg, set 95% CL cross section limit on signal

4) Result is returned by email



#### Total elapsed time $\approx$ 1 hour



#### Algorithm

Result

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From: quaero@fnal.gov Subject: Quaero Request #29

 $W_R \rightarrow t \overline{b} \rightarrow e \not\!\!\! E_T 2 j$ 

Result

Pythia cross section x branching ratio = 1.68 pb.

Upper limits on the cross section to this process at confidence levels of 50%, 90%, and 95% are found to be 0.8 pb, 1.8 pb, and 2.1 pb, respectively. Maximal sensitivity (0.73 pb<sup>-1</sup>) is achieved in a region of variable space with 17.6 signal events expected, 32.7 +- 7.1 background events expected, and 36 events observed in the data.

#### Plots

Plots of the variables that you used are available for viewing at <u>http://quaero.fnal.gov/quaero/requests/plots/29.ps</u>. The red curve is the expected background; the green curve is your signal multiplied by a factor of 10; the black dots are D0 data.





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