# Elliptic Flow of Strange-hadron $K_{s}{ }^{0}$ and $\Lambda$ in $A u+A u$ Collisions at 200 GeV — high statistics Run IV <br> Y. Lu ${ }^{1,2}$, P. Sorensen ${ }^{1}$ <br> ${ }^{1}$ Nuclear Science Division, Lawrence Berkeley National Laboratory, Berkeley, California 94720 <br> ${ }^{2}$ Institute of Particle Physics, Central China Normal University, Wuhan, China 430079 



Figure 1: v2 as a function of transverse momentum $\mathrm{p}_{\mathrm{T}}$ for $\mathrm{K}_{\mathrm{S}}{ }^{0}$ (filled circles) and $\Lambda$ (filled triangles) from $200 \mathrm{GeV} \mathrm{Au}+\mathrm{Au}$ minimum bias collisions. The charged hadron $v_{2}$ is from [4]. Only statistical uncertainty is included in the errors.

Elliptic flow, due to its self quenching nature, is sensitive to the early stage of the dense system created in ultrarelativistic heavy ion collisions [1]. Because (multi-) strange hadrons have smaller hadronic cross-sections than nonstrange hadrons [2], the flavor dependence of elliptic flow can provide insight into the partonic stage of heavy ion collisions. Recent work established that at low $\mathrm{p}_{\mathrm{T}} \mathrm{v}_{2}$ depends on particle mass while at intermediate $\mathrm{p}_{\mathrm{T}}$ it saturates and appears to depend on the number of constituent quarks [3, 5, 6]. A scaling law is naturally expected from coalescence or recombination models. Above $\mathrm{p}_{\mathrm{T}} \sim 5 \mathrm{GeV} / c$, with large error, $\mathrm{V}_{2}$ of $\mathrm{K}_{\mathrm{S}}{ }^{0}$ and $\Lambda$ seems to take on a value supported by $\mathrm{R}_{\mathrm{CP}}$ of $\mathrm{K}_{\mathrm{S}}{ }^{0}$ and $\Lambda$ close to that of charged hadrons [3].

Using the STAR TPC at RHIC, $\mathrm{K}_{\mathrm{S}}{ }^{0}$ and $\Lambda$ are reconstructed via their decay topology $\mathrm{K}_{\mathrm{S}}{ }^{0} \rightarrow \pi^{+}+\pi^{-}$and $\Lambda \rightarrow p+$ $\pi^{-}$. Cuts on geometry, kinematics and particle identification via specific ionization are applied to reduce combinatorial background. The remaining background is counted for in our analysis procedure. We calculate the elliptic flow parameter $\mathrm{v}_{2}$ from the particle momentum-space azimuthal angle with respect to the reaction plane. The basic idea of this method is that $1 . v_{2}$ of real and background particles is averaged by real particle $\mathrm{v}_{2}$ and background $\mathrm{v}_{2}$ according to following equation:

$$
v_{2}^{R E A L+B G}(m)=v_{2}^{R E A L} \times \frac{R E A L}{R E A L+B G}(m)+v_{2}^{B G}(m) \times \frac{B G}{R E A L+B G}(m)
$$

2. $v_{2}$ of background is smoothly changing with invariant mass. By measuring the ratio of real and background counts
vs. invariant mass and $\mathrm{v}_{2}$ vs. invariant mass, we can extract $\mathrm{v}_{2}$ for the real counts. For this analysis, about 11 million minimum bias events are used from $\mathrm{Au}+\mathrm{Au}$ collisions at 200 GeV .
Figure 1 shows elliptic flow $\mathrm{v}_{2}$ as a function of $\mathrm{p}_{\mathrm{T}}$ for charged hadrons (open diamonds), $\mathrm{K}_{\mathrm{S}}{ }^{0}$ (filled circles) and $\Lambda$ (filled triangles) for the minimum bias data set. $v_{2}$ of charged hadrons, $\mathrm{K}_{\mathrm{S}}{ }^{0}$ and $\Lambda$ first increases with transverse momentum and then saturates or decreases. Up to $6 \mathrm{GeV} / c$, there is a clear particle type dependence. At low $\mathrm{p}_{\mathrm{T}}$, smaller $\mathrm{v}_{2}$ of $\Lambda$ than that of $\mathrm{K}_{\mathrm{S}}{ }^{0}$ is consistent with mass ordering as predicted by the hydrodynamic model [7], which assume local thermal equilibrium. At $\mathrm{p}_{\mathrm{T}} \sim 2 \mathrm{GeV} / c \mathrm{v}_{2}$ of $\mathrm{K}_{\mathrm{S}}{ }^{0}$ and $\Lambda$ cross over and continue to increase. $\mathrm{v}_{2}$ of $\mathrm{K}_{\mathrm{S}}{ }^{0}$ starts to decrease at $\mathrm{p}_{\mathrm{T}} \sim 3.6 \mathrm{GeV} / c$ while $\mathrm{v}_{2}$ of $\Lambda$ stays large up to 6 $\mathrm{GeV} / c$. At intermediate $\mathrm{p}_{\mathrm{T}}$ the $\mathrm{K}_{\mathrm{S}}{ }^{0}$ and $\Lambda$ measurements seem to indicate that $\mathrm{v}_{2}$ doesn't depend on particle mass at this region. These more precise measurements supports the previously suggested baryon to meson dependence which is qualitatively consistent with quark coalescence or recombination models of hadronization. Above $\mathrm{p}_{\mathrm{T}} \sim 3.6 \mathrm{GeV} / c, \mathrm{~K}_{\mathrm{S}}{ }^{0}$ $\mathrm{v}_{2}$ decreases with increasing $\mathrm{p}_{\mathrm{T}} . \Lambda \mathrm{v}_{2}$ increases with $\mathrm{p}_{\mathrm{T}}$ until it reaches its maximum measured value at $\mathrm{p}_{\mathrm{T}} \sim 4 \mathrm{GeV} / c$. The $\mathrm{K}_{\mathrm{S}}{ }^{0}$ and $\Lambda \mathrm{v}_{2}$ may merge for $\mathrm{p}_{\mathrm{T}}>6 \mathrm{GeV} / c$ but with the current statistical errors the measurements remain inconclusive.

## REFERENCES

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