# A 3D Triple-Resonance NMR Technique for Qualitative Measurement of Carbonyl- $\mathrm{H} \beta \boldsymbol{J}$ Couplings in Isotopically Enriched Proteins 

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It is well known that the magnitude of heteronuclear long-range couplings contains useful information regarding molecular conformation (1). For measurement of such couplings between protons and protonated heteronuclei, E.COSY-based techniques (2) in principle can yield quantitative results in isotopically enriched proteins (3-6). For measurement of $J$ couplings to nonprotonated heteronuclei, deconvolution of the cross-peak shape in a ${ }^{1} \mathrm{H}$-detected heteronuclear multiple-bond correlation (HMBC) spectrum (7) can be used (8-10), although for larger proteins the ${ }^{1} \mathrm{H}$ linewidth and extensive spectral overlap render this approach difficult, if not impossible.

Alternatively, qualitative information about the size of a long-range coupling can be obtained from the intensity of correlations in a 3D spectrum in which magnetization is transferred via the long-range coupling of interest. This has recently been demonstrated for couplings between ${ }^{15} \mathrm{~N}$ and intraresidue $\mathrm{H} \beta$ protons (11, 12). Here we present an analogous experiment that uses magnetization transfer via the $J$ coupling between the carbonyl $(\mathrm{CO})$ and the $\mathrm{H} \beta$ protons. The resulting information about the $\mathrm{CO}-\mathrm{H} \beta$ coupling, in combination with knowledge of the $\mathrm{H} \beta-{ }^{15} \mathrm{~N} J$ coupling and qualitative (13) or quantitative (2) data on the $\mathrm{H} \alpha-\mathrm{H} \beta$ couplings, allows unambiguous stereospecific assignment of the $\mathrm{H} \beta$ methylene protons and determines the rotomeric state of the $\mathrm{C} \alpha-\mathrm{C} \beta$ bond ( $\chi_{1}$ angle).

The pulse sequence of our new experiment, which relies on magnetization transfer via the multiple-bond ${ }^{1}{ }^{1}{ }^{13}{ }^{13} \mathrm{CO} J$ coupling, is sketched in Fig. 1. Because its main function is to correlate amide protons and ${ }^{15} \mathrm{~N}$ nuclei to $\mathrm{H} \beta$ protons, using relay via the $\mathrm{CO}^{13} \mathrm{C}$ nucleus, we name the experiment $\mathrm{HN}(\mathrm{CO}) \mathrm{HB}$. In this experiment, amide proton magnetization is transferred via its directly bonded ${ }^{15} \mathrm{~N}$ nucleus to the adjacent carbonyl spin. A subsequent $\mathrm{HMQC}(14-16)$ segment correlates ${ }^{13} \mathrm{CO}$ with the $\mathrm{H} \beta$ protons and finally ${ }^{13} \mathrm{CO}$ magnetization is transferred back to the amide proton for detection. As recently shown for the triple-resonance HNCO experiment (17, 18), magnetization transfer from the amide proton to the carbonyl is quite efficient and uniform because this transfer takes place via large and well-resolved one-bond cou-

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Fig. 1. Pulse scheme of the $\mathrm{HN}(\mathrm{CO}) \mathrm{HB}$ experiment. Narrow pulses correspond to a $90^{\circ}$ flip angle and wide pulses correspond to $180^{\circ}$. Pulses for which the phase is not indicated are applied along the $x$ axis. Phases and phase cycling of the other pulses are as follows: $\phi_{1}=x,-x ; \phi_{2}=y,-y ; \phi_{3}=x ; \phi_{4}=4(x)$, $4(y), 4(-x), 4(-y) ; \phi_{5}=16(x), 16(-x) ; \phi_{6}=8(x), 8(-x) ; \phi_{7}=x, x,-x,-x ; \phi_{8}=16(x), 16(-x) ;$ $\phi_{9}=4(x), 8(-x), 4(x)$; Acq. $=4(x,-x,-x, x), 4(-x, x, x,-x)$. Quadrature in the $t_{1}$ dimension is obtained by altering the phase $\phi_{3}$ in a States-TPPI manner (27), and similarly, $\phi_{8}$ is altered in a StatesTPPI manner for quadrature detection in the $t_{2}$ dimension. Delay durations used are as follows: $\tau=2.25$ $\mathrm{ms} ; T=13.5 \mathrm{~ms} ; T^{\prime}=T+\tau_{180^{\circ}}(\mathrm{C} \alpha)$, where $\tau_{180^{\circ}}(\mathrm{C} \alpha)$ is the duration of the $180^{\circ} \mathrm{C} \alpha$ pulse $(162 \mu \mathrm{~s}) ; \epsilon=$ $25 \mathrm{~ms} ; \delta_{0}=10 \mathrm{~ms}+\tau_{180^{\circ}}\left({ }^{\prime} \mathrm{H}\right)$, with $\tau_{180^{\circ}}\left({ }^{\prime} \mathrm{H}\right)=21 \mu \mathrm{~S} ; \delta_{1}=7.25 \mathrm{~ms} ; \delta_{2}=2.75 \mathrm{~ms}$. C $\alpha$ pulses were applied at a frequency corresponding to $56 \mathrm{ppm} ; \mathrm{C} \alpha / \beta$ composite ( $90_{x}^{\circ} 180^{\circ} 90_{x}^{\circ}$ ) $180^{\circ}$ pulses were applied at 44.5 ppm , with the RF power adjusted to give a $48.5 \mu \mathrm{~S} 90^{\circ}$ pulse (yielding a null at the CO frequency for a ${ }^{13} \mathrm{C}$ spectrometer frequency of 151 MHz ) and the carbonyl carrier was set at 177 ppm . The ${ }^{1} \mathrm{H}$ carrier is positioned on the $\mathrm{H}_{2} \mathrm{O}$ resonance and $\mathrm{H}_{2} \mathrm{O}$ presaturation is used.
plings. As will be discussed below, correlating the carbonyl with the $\mathrm{H} \beta$ resonances results in an approximate $\left({ }^{3} J_{\mathrm{CO}-\mathrm{H} \beta}\right)^{2}$ dependence for the cross-peak intensity, thereby providing the information of interest.

First, we briefly outline the relevant magnetization-transfer steps, using the productoperator formalism to describe the spin system. For clarity, relaxation terms are not included and constant multiplicative factors are omitted. Only terms that result in observable magnetization during the detection period, $t_{3}$, are retained. The spin operators used are I for the amide proton, N for the amide ${ }^{15} \mathrm{~N}, \mathrm{~S}$ for the carbonyl ${ }^{13} \mathrm{C}$, and B for the $\mathrm{H} \beta$ proton. $J_{\mathrm{XY}}$ denotes the $J$ coupling between spins X and Y , which is always a one-bond coupling except for $J_{\mathrm{SB}}$ and $J_{\mathrm{S} k}$, which denote a three-bond $\mathrm{H} \beta-$ CO $J$ coupling and a multiple-bond $J$ coupling between CO and a proton $k$, respectively. RF phases correspond to the first step of the phase cycle given in the legend to Fig. 1.

Longitudinal amide-proton magnetization $\left(I_{z}\right)$, present at the start of the experiment, is transferred by means of an INEPT sequence into antiphase ${ }^{15} \mathrm{~N}$ magnetization, yielding

$$
\begin{equation*}
\sigma_{a}=N_{y} I_{z} \cos \left(2 \pi J_{\mathrm{NI}} \tau\right) . \tag{1}
\end{equation*}
$$

The following ${ }^{15} \mathrm{~N}$ evolution period, $t_{1}$, is of the constant-time variety (19-21). Because the ${ }^{1} \mathrm{H}$ and ${ }^{15} \mathrm{~N} 180^{\circ}$ pulses are applied simultaneously during this "constant-time" interval ( $T+T^{\prime}$ ), the evolution due to $J_{\mathrm{NH}}$ coupling is not affected. Its total duration ( $T+T^{\prime}$ ) is chosen to be an odd multiple of $1 /\left(2 J_{\mathrm{NH}}\right)$ ( 27 ms in our experiment), such that the ${ }^{15} \mathrm{~N}$ magnetization at the end of the $t_{1}$ evolution period (time $b$ ) is in phase with respect to spin I. The $180^{\circ}{ }_{\phi 5}(\mathrm{CO})$ pulse is applied simultaneously with the $180_{\phi 4}^{\circ}\left({ }^{15} \mathrm{~N}\right)$ pulse to ensure that the ${ }^{15} \mathrm{~N}$ magnetization becomes antiphase with respect to the carbonyl spin, S . The power and duration of the carbonyl $180^{\circ}$ pulse are carefully adjusted in such a manner that the $\mathrm{C} \alpha$ resonances do not experience this pulse; i.e., for $t_{1}=0$, the ${ }^{15} \mathrm{~N}$ spin remains in phase with respect to $J$-coupled C $\alpha$ nuclei. A $180^{\circ}$ $\mathrm{C} \alpha$ pulse applied during the constant-time evolution period ensures that the ${ }^{15} \mathrm{~N}$ magnetization remains in phase with respect to $\mathrm{C} \alpha$, independent of the $t_{1}$ value. At the end of the constant-time evolution period (time $b$ ) the relevant magnetization is described by

$$
\begin{equation*}
\sigma_{b}=N_{y} S_{z} \cos \left(2 \pi J_{\mathrm{NI}} \tau\right) \sin \left(2 \pi J_{\mathrm{NS}} T\right) \cos \left(\Omega_{\mathrm{N}} t_{1}\right) \tag{2}
\end{equation*}
$$

where $\Omega_{\mathrm{N}}$ denotes the ${ }^{15} \mathrm{~N}$ angular offset frequency. Omitting the trigonometric terms, at time $c$ one obtains

$$
\begin{equation*}
\sigma_{c}=N_{z} S_{y} \tag{3}
\end{equation*}
$$

During the following intervals, between times $c$ and $d$, the carbonyl magnetization remains in the transverse plane, but because of the $180^{\circ}$ pulse applied at the middle of $t_{2}$, its chemical shift does not need to be taken into account. During the time $\epsilon$, between time points $c$ and $d$, S-spin magnetization becomes antiphase with respect to its long-range coupled protons. During the $\epsilon$ delay, the S -spin magnetization also dephases because of $J$ coupling to the adjacent $\mathrm{C} \alpha$ and ${ }^{15} \mathrm{~N}$ spins. However, the ${ }^{13} \mathrm{CO}$ dephasing caused by these interactions is also refocused at the end of the second $\epsilon$ interval and therefore may be ignored. Thus, at time $d$ one obtains

$$
\begin{equation*}
\sigma_{d}=N_{z} S_{x} B_{z} \sin \left(\pi J_{\mathrm{SB}} \epsilon\right) \prod_{k} \cos \left(\pi J_{\mathrm{S} k} \epsilon\right) \tag{4}
\end{equation*}
$$

which is converted into $S-B$ multiple-quantum coherence by the subsequent ${ }^{1} \mathrm{H}$ $90_{\phi 8}^{\circ}$ pulse,

$$
\begin{equation*}
\sigma_{e}=N_{z} S_{x} B_{y} \tag{5}
\end{equation*}
$$

where the trigonometric terms again have been omitted temporarily.
During both the first half and the second half of the $t_{2}$ evolution period, dephasing of the $B_{y}$ term in expression [5] by coupling to its directly attached ${ }^{13} \mathrm{C}$ nucleus is prevented by the (composite) $180^{\circ}$ pulses applied to the $\mathrm{C} \alpha$ and $\mathrm{C} \beta$ resonances. To avoid the necessity for large first-order phase corrections in the $F_{2}$ dimension of the spectrum, these pulses are applied only for durations of $t_{2} / 2$ that are larger than the duration of the composite $180^{\circ} \mathrm{C} \alpha / \beta$ pulse ( $194 \mu \mathrm{~s}$ ). The effect of S-spin resonance offset is removed by the $180^{\circ}(\mathrm{CO})$ pulse applied at the midpoint of $t_{2}$. Note that one cannot use a single $180^{\circ} \mathrm{C} \alpha / \beta$ pulse to remove $J$ coupling to proton B, because this would reintroduce the effect of the $\mathrm{CO}-\mathrm{C} \alpha J$ coupling. At the end of the $t_{2}$ evolution
period the $S-B$ multiple-quantum coherence is converted back into antiphase S-spin magnetization,

$$
\begin{equation*}
\sigma_{f}=N_{z} S_{x} B_{z} \cos \left(\Omega_{\mathrm{B}} t_{2}\right) \tag{6}
\end{equation*}
$$

where $\Omega_{\mathrm{B}}$ denotes the angular offset frequency of proton B . The transfer of $\sigma_{f}$ back into observable amide-proton magnetization follows the reverse pathway from that described above.

After the relevant trigonometric terms are reintroduced, the magnetization at the beginning of the detection period is described by

$$
\begin{align*}
& \sigma_{g}=I_{x} \cos ^{2}\left(2 \pi J_{\mathrm{N} \mathrm{I}} \tau\right) \sin \left(2 \pi J_{\mathrm{NS}} T\right) \cos \left(\Omega_{\mathrm{N}} t_{1}\right) \cos \left(\Omega_{\mathrm{B}} t_{2}\right) \\
& \times \sin \left[\pi\left(\delta+\delta_{1}+\delta_{2}\right) J_{\mathrm{NS}}\right] \sin ^{2}\left(\pi J_{\mathrm{SB}} \epsilon\right) \prod_{k} \cos ^{2}\left(\pi J_{\mathrm{S} k} \epsilon\right) \tag{7}
\end{align*}
$$

${ }^{1} \mathrm{H}-{ }^{15} \mathrm{~N} J$ couplings in proteins are all close to $\sim 92 \mathrm{~Hz}$ and $J_{\mathrm{NS}}$ values are also quite uniform, varying from 13 to 16.5 Hz (22). Therefore, the sine and cosine terms containing $J_{\mathrm{NS}}$ and $J_{\mathrm{NI}}$ in expression [7] may be considered constants. To evaluate the attenuation caused by the $\cos ^{2}\left(\pi J_{\mathrm{S} k} \epsilon\right)$ terms in Eq. [7], some knowledge of the magnitude of these couplings is required. Hansen et al. (23) report ${ }^{2} J_{\mathrm{CO}-\mathrm{H} \alpha}$ couplings in the range of $4.2-7.1 \mathrm{~Hz}$. Interresidue ${ }^{3} J_{\mathrm{CO}-\mathrm{H} \alpha}$ couplings are usually smaller than 4 Hz , except for the rare case where the intervening $\mathrm{N}-\mathrm{C} \alpha$ dihedral angle is close to $+60^{\circ}$ [ corresponding to a trans orientation of $\mathrm{CO}(i)$ relative to $\mathrm{H} \alpha(i+1)$ ]. For the intraresidue ${ }^{3} J_{\mathrm{CO}-\mathrm{H} \beta}$ coupling, Hansen et al. suggest a trans ${ }^{3} J_{\mathrm{CO}-\mathrm{H} \beta}$ value of 11.9 Hz and a gauche value of 0.4 Hz . Bermel et al. (8) utilize less extreme values of 8.5 and 1.4 Hz , respectively. Considering that, to the best of our knowledge, no ${ }^{3} J_{\mathrm{CO}-\mathrm{H} \beta}$ coupling larger than 8 Hz has ever been measured in a polypeptide, an upper limit of $\sim 8 \mathrm{~Hz}$ for this coupling appears reasonable. Using an $\epsilon$ delay of 25 ms , passive couplings of 4 and 8 Hz cause an attenuation of 9 and $35 \%$ in cross-peak intensity, respectively. Hence, significant attenuation can be caused by these passive couplings. When the $\mathrm{CO}-\mathrm{H} \beta$ connectivity for the case of nonequivalent $\beta$-methylene protons is studied, the passive coupling effect accentuates the difference in $J$ coupling constant because a large active ${ }^{3} J_{\mathrm{CO}-\mathrm{H} \beta 1}$ coupling invariably combines with a small passive ${ }^{3} J_{\mathrm{CO}-\mathrm{H} \beta 2}$ coupling, and vice versa.

The absolute intensity of an ${ }^{n} J_{\mathrm{OO}-\mathrm{H}}$ correlation also depends on an additional number of factors. First, presaturation of the $\mathrm{H}_{2} \mathrm{O}$ resonance may decrease the amide-proton magnetization via cross relaxation or via exchange of the amide proton with water, thus attenuating the observed correlations correspondingly. Second, ${ }^{15} \mathrm{~N}$ and ${ }^{13} \mathrm{CO}$ transverse relaxation times may vary substantially for different amino acids in a protein, significantly affecting the absolute intensity of an observed ${ }^{n} J_{\mathrm{CO}-\mathrm{H}}$ correlation. Unless both these relaxation rates and the effect of $\mathrm{H}_{2} \mathrm{O}$ presaturation are known, it is most prudent to use the relative intensities of the correlations observed for a given amide proton to $\mathrm{H} \alpha$ and $\mathrm{H} \beta$ resonances to make a qualitative estimate of the size of the ${ }^{3} J_{\mathrm{CO}-\mathrm{H} \beta}$ coupling. If, for example, a strong correlation to $\mathrm{H} \alpha$ but no correlations to $\mathrm{H} \beta$ are observed, this indicates that the CO-H $\beta$ coupling is small. However, if no correlations to either $\mathrm{H} \alpha$ or $\mathrm{H} \beta$ are observed, this may be caused by other factors such as broad ${ }^{13} \mathrm{CO}$ or ${ }^{15} \mathrm{~N}$ linewidths or $\mathrm{H}_{2} \mathrm{O}$ presaturation.

The $\mathrm{HN}(\mathrm{CO}) \mathrm{HB}$ experiment is demonstrated for the protein calmodulin ( $M_{\mathrm{r}} \sim 16.7$ kDa ) complexed with $\mathrm{Ca}^{2+}$. The experiment was carried out at a $600 \mathrm{MHz}{ }^{1} \mathrm{H}$ frequency on an unmodified Bruker AMX-600 spectrometer, using a $1.5 \mathrm{~m} M$ sample concentration in $95 \% / 5 \% \mathrm{H}_{2} \mathrm{O} / \mathrm{D}_{2} \mathrm{O}, \mathrm{pH} 6.3,35^{\circ} \mathrm{C}$. On average, ${ }^{15} \mathrm{~N}$ transverse relaxation times in the ${ }^{1} \mathrm{H}$-coupled mode are $\sim 70 \mathrm{~ms}$, as judged by the ${ }^{15} \mathrm{~N}$ linewidth observed in an Overbodenhausen (HSQC) spectrum, and ${ }^{13} \mathrm{CO}$ transverse relaxation times (in the ${ }^{1} \mathrm{H} /{ }^{15} \mathrm{~N} /{ }^{13} \mathrm{C}$-coupled mode) were measured to be $\sim 50 \mathrm{~ms}$. Significant variations of both ${ }^{15} \mathrm{~N}$ and ${ }^{13} \mathrm{CO}$ transverse relaxation, however, are present within the protein.

Figure 2 shows sections of two $F_{2}\left({ }^{1} \mathrm{H}\right)-F_{3}(\mathrm{HN})$ slices taken from the 3D $\mathrm{HN}(\mathrm{CO}) \mathrm{HB}$ spectrum, displaying connectivities between amide and aliphatic protons. As discussed before, the carbonyl nucleus is used for magnetization relay purposes only, and its frequency is not measured in this particular experiment. However, the intensity of a correlation between an amide proton $\left(F_{3}\right)$ and $\mathrm{H} \alpha$ and $\mathrm{H} \beta$ protons ( $F_{2}$ ) depends primarily on the magnitude of the heteronuclear $J$ coupling between the carbonyl that precedes the observed amide proton and $\mathrm{H} \alpha / \beta$ (cf. Eq. [7]). For example, on the left side of Fig. 2B, a correlation is observed between the amide proton of Phe-65 and the upfield $\mathrm{H} \beta^{\prime}$ proton of Asp-64; the position of the second Asp-64 proton, $\mathrm{H} \beta$, is marked by the empty box. Clearly, the $\mathrm{H} \beta^{\prime}$ proton is trans relative to the carbonyl, and $\mathrm{H} \beta$ is in a gauche position. For Asp-64, the correlation to its $\mathrm{H} \alpha$ proton (at 5.40 ppm ) falls outside the window shown but also has very low intensity, indicative of a small ${ }^{2} J_{\mathrm{CO}-\mathrm{H} \alpha}$ coupling. For the majority of residues, however, relatively intense correlations via the intraresidue ${ }^{2} J_{\mathrm{CO}-\mathrm{H} \alpha}$ are present. For example, all other residues for which correlations are present in Fig. 2 exhibit relatively intense correlations to $\mathrm{H} \alpha$ protons. These resonances all fall in the $3.7-4.4 \mathrm{ppm}$ range and correspond to transfer via ${ }^{2} J_{\mathrm{CO}-\mathrm{H} \alpha}$; i.e., they represent correlations between the amide of residue $i$ and $\mathrm{H} \alpha$ of residue $i-1$.

In the entire 3D spectrum, only 11 intraresidue $\mathrm{HN}-\mathrm{H} \alpha$ correlations (via ${ }^{3} J_{\mathrm{CO}-\mathrm{H} \alpha}$ ) were observed and none are present in Fig. 2. An intense intraresidue $\mathrm{NH}-\mathrm{H} \alpha$ correlation in the $\mathrm{HN}(\mathrm{CO}) \mathrm{HB}$ spectrum is indicative of a positive $\phi$ angle near $60^{\circ}$, except for glycines, where $\phi$ angles of either $+60^{\circ}$ or $-60^{\circ}$ are expected to give rise to intense correlations. Four of the eleven observed correlations correspond to glycine residues, G23, G96, G132, and G59; they all have positive $\phi$ angles in the X-ray crystal structure (24) and all four are located in homologous positions in the four calciumbinding loops of calmodulin. For these residues the correlation is observed to the proR H $\alpha$ proton, not to the pro-S proton which replaces $\mathrm{C} \beta$. The remaining intraresidue correlations correspond to residues for which the crystal structure indicates that the $\mathrm{H} \alpha$ is close to a cis arrangement with respect to the carbonyl $\left(\phi \approx-120^{\circ}\right)$, or for which ${ }^{15} \mathrm{~N}$ relaxation studies indicate a high degree of mobility (25).

The $\mathrm{HN}(\mathrm{CO}) \mathrm{HB}$ experiment is considerably more complex than its counterpart, HNHB. In part this is due to the extra magnetization relay step, transferring magnetization via the carbonyl carbon. However, an added degree of difficulty is caused by the fact that the present experiment requires a protein that is highly enriched in both ${ }^{13} \mathrm{C}$ and ${ }^{15} \mathrm{~N}$. Hence, it becomes necessary to remove the effect of both ${ }^{1} \mathrm{H} \alpha / \beta-{ }^{13} \mathrm{C} \alpha /$ $\beta,{ }^{15} \mathrm{~N}-{ }^{13} \mathrm{C} \alpha$, and ${ }^{13} \mathrm{CO}-{ }^{13} \mathrm{C} \alpha J$ couplings. Despite the large number of RF pulses used in the $\mathrm{HN}(\mathrm{CO}) \mathrm{HB}$ experiment, a high-quality spectrum can be recorded in a


FIG. 2. Two sections of ( $F_{2}, F_{3}$ ) slices taken from the $\mathrm{HN}(\mathrm{CO}) \mathrm{HB}$ spectrum of calmodulin, showing correlations between the amide of one residue (marked below each dotted line) and $\mathrm{H} \alpha$ (unlabeled) and $\mathrm{H} \beta$ resonances of the preceding amino acid. Square boxes mark the positions of $\mathrm{H} \beta$ protons with vanishing intensity, i.e., with a small ${ }^{3} J_{\mathrm{CO}-\mathrm{H} \beta}$. The "invisible" $\mathrm{H} \beta$ peak positions have been taken from Ikura et al. (28). The spectrum results from a ( 36 complex $) \times(64$ complex $) \times(512$ complex) data matrix, recorded with 128 scans per hypercomplex ( $t_{1}, t_{2}$ ) pair. Acquisition times were $27 \mathrm{~ms}\left(t_{1},{ }^{15} \mathrm{~N}\right), 12.8 \mathrm{~ms}\left(t_{2},{ }^{1} \mathrm{H}\right)$, and $55.3 \mathrm{~ms}\left(t_{3}, \mathrm{HN}\right)$. Zero filling and $60^{\circ}$-shifted sine-bell filtering were used in the $t_{2}$ dimension and zero filling combined with a squared $60^{\circ}$-shifted sine bell was used for $t_{3}$. After $t_{2}$ and $t_{3}$ Fourier transformation, the $t_{1}$ time domain was extended to 72 complex data points using linear prediction with a mirror image constraint (29), followed by a squared $90^{\circ}$-shifted sine-bell filter and zero filling to 128 points, prior to the final $t_{1}$ Fourier transform.
reasonable amount of measuring time ( $\sim 3.5$ days) using a relatively low sample concentration ( $1.5 \mathrm{~m} M$ ) of a protein which is large by NMR standards. The information that it provides is important for determining stereospecific assignments of $\mathrm{H} \beta$ methylene protons and $\chi_{1}$ torsion angles, where it complements other techniques that rely on measurement of $\mathrm{H} \alpha-\mathrm{H} \beta$ and ${ }^{15} \mathrm{~N}-\mathrm{H} \beta$ couplings. In addition, the experiment
provides important sequential $\mathrm{NH}(i)-\mathrm{H} \alpha(i-1)$ connectivity information, complementing results obtainable with the $\mathrm{H}(\mathrm{CA}) \mathrm{NNH}(26)$ experiment.

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