EXAMINATION OF MULTICOMPONENT DIFFUSION BETWEEN TWO NI-BASE SUPERALLOYS

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ABSTRACT

The interdiffusion at 1293 °C between two multicomponent Ni-base superalloys, René-N4 and René-N5, was assessed by measuring the composition vs. distance curves and by comparing the measured curves to predictions obtained using a diffusion mobility database recently published by Campbell et al. (Acta Mat.50 (2002) 775-792). Although the diffusion database was constructed primarily from binary diffusion data, the extrapolation to the multicomponent systems gave good results in the prediction of the measured composition vs. distance curves. In addition, the location of the Kirkendall porosity on the René-N4 side of the diffusion couple was successfully predicted. This initial success points to the suitability of the general approach to the development of diffusion databases and to the desirability for additional database refinements including possible efforts from the first principles community.

INTRODUCTION

Diffusion data is needed for quantitative predictions of many materials processing and phase transformation models. In the area of superalloys alone, examples are abundant. The calculation of solid diffusion during solidification is required to predict microsegregation and the amount and type of second phase particles. Diffusion data is also required to predict incipient melting temperatures during reheating for solution treatment and to determine γ' size distributions during aging or cooling. The prediction of the phase sequence during transient liquid phase bonding and during the processing and exposure of thermal barrier coatings also requires multicomponent diffusion calculations.

Recently, there has been some effort to develop multicomponent diffusion mobility databases for the FCC phase of Ni based alloys¹⁻³ using the formalism put forth by Ågren^{4,5} and Andersson and Ågren⁶. These databases are developed to reproduce

measured tracer, intrinsic and chemical diffusion data in binary and ternary systems and to permit extrapolation to higher order systems. The development and utilization of a mobility database relies on the pre-existence of a separate thermodynamic database for the phases of interest. The thermodynamic database provides the necessary thermodynamic factors to convert from chemical potential gradients to concentration gradients.

A brief summary of the diffusion formalism of $Ågren^{4,5}$ and Andersson and $Ågren^{6}$ as applied to a disordered substitutional solid solution is useful. Diffusion is assumed to occur by a vacancy exchange mechanism, in which the equilibrium vacancy concentration is maintained. The partial molar volumes of the substitutional species are assumed to be equal. For a given phase, the flux of species *i* in the *z*-direction in volume-fixed frame of reference is given by

$$J_{i} = \sum_{j=1}^{n-1} \widetilde{D}_{ij}^{n} \frac{\partial c_{j}}{\partial z}$$
(1)

where c_j is the concentration of species *j*, species *n* is the dependent (solvent) species, and \widetilde{D}_{ij}^n is the interdiffusion coefficient. The interdiffusion coefficient, also known as the chemical diffusivity, can be expressed as the difference between the intrinsic diffusivities.

$$D_{ij}^n = D_{ij} - D_{in} \,. \tag{2}$$

The intrinsic diffusivities, D_{ij} , are defined in terms of the atomic mobility, M_p , as

$$D_{ij} = \sum_{p=1}^{n} \left(\delta_{pi} - x_i \right) x_p M_p \frac{\partial \mu_p}{\partial x_j}$$
(3)

where the partial derivative of the chemical potential, μ_p , with respect to the mole fraction, x_i , defines the thermodynamic factors, which can be calculated using an appropriate multicomponent thermodynamic database, such as those by Saunders⁷ and Kattner⁸. It should be noted that the thermodynamic factors must be evaluated in the form $\mu_k(x_1, x_2, ..., x_n)$ where *n* is the dependent species.

The atomic mobility, M_p , of species p in a given phase is both composition and temperature dependent.

$$M_{p} = \Theta_{p} \frac{1}{RT} \exp\left(\frac{\Delta Q_{p}^{*}}{RT}\right)$$
(4)

where Θ_p represents the effects of the atomic jump distance (squared) and the jump frequency and has the units of m²/s. The variable ΔQ_p^* is the diffusion activation energy of specie *p* in a given phase with units of (J/mol). The variable *R* is the gas constant is and the temperature, *T*, is in Kelvin. As ΔQ_p^* and Θ_p can be combined into one parameter, it is customary¹ to let Θ_p equal 1 and only treat the temperature and composition dependence of ΔQ_p^* . Ågren and co-workers^{3,6,9-12} expressed the composition and temperature dependence of each ΔQ_p^* in terms of a Redlich-Kister¹³ polynomial, as seen in eqn (5),

$$\Delta Q_{p}^{*} = \sum_{j} x_{j} Q_{i}^{j} + \sum_{q} \sum_{j > q} x_{q} x_{j} \sum_{k} {}^{k} A_{i}^{qj} (x_{q} - x_{j})^{k}, \qquad (5)$$

where the Q_i^j and the ${}^kA_i^{qj}$ are linear functions of temperature. The expansion of the composition dependence in terms of a Redlich-Kister¹³ polynomial is similar to the CALPHAD approach^{14,15} used in the development of the thermodynamic databases. Note for a given diffusing species, *i*, that if all Q_i^j s are equal and A_i^{qj} equals zero, then ΔQ_p^* and the corresponding M_p are not concentration dependent.

Optimized mobility functions² were obtained using the PARROT¹⁶ optimization code to evaluate the composition and temperature dependence of ΔQ_p^* using the available experimental diffusion tracer, intrinsic, and chemical diffusivity data. The experimental data were weighted giving preference to the tracer diffusivity, D_p^* , data, which are independent of concentration, as seen by equation (6)

$$D_p^* = RTM_p \quad . \tag{6}$$

The current approach requires defining metastable end-member diffusion mobilities, such as the self-diffusion in fcc-W. Determination of these end-member quantities follows approaches similar to those used to determine the lattice stabilities of the metastable thermodynamic quantities of the elements^{14,15}. This determination of diffusion activation energies for metastable end-member phases enables the extrapolation to higher order systems where diffusion data may be limited. Presently, the only available check on these assessed values is the application of a diffusion correlation, which states the following for pure FCC metals¹⁷

$$\frac{-Q}{RT_M} \approx 17\tag{7}$$

where Q is the diffusion activation energy, T_M is the melting temperature. For Ni-W the activation energy, Q, for fcc-W was optimized to be -311420 J/mol and the metastable fcc melting temperature is 2229 K. This gives an acceptable ratio equal to 16.8. First principles calculations of these metastable quantities would be useful.

To examine the validity of this diffusion mobility database, diffusion simulations are compared to results from an experimental multicomponent diffusion couple profile.

EXPERIMENTAL AND COMPUTATIONAL PROCEDURES

René-N4 and René-N5, first-generation and second-generation superalloys respectively, were chosen as the two sides of the diffusion couple. The sample geometry of each half of the diffusion couple consisted of a 2.25 cm square with a thickness of 6.35 mm. The diffusion couples were pre-bonded at 1277 °C under a load of 180 N for 2 h. After bonding, the diffusion couples were diffusion heat-treated at 1293 °C for 10 h and 100 h in a vacuum furnace and gas-cooled. The 1293 °C temperature was chosen to ensure that both superalloys would be in the single phase γ region. The diffusion couples were characterized using scanning electron microscopy and microprobe analysis using standard ZAF correction and elemental standards. Gas cooling of the diffusion couples from 1293 °C produced coarse γ' precipitates (see Figure 1.) which resulted in significant scatter in the experimental composition profiles. This scatter is due to the chance of measuring the composition of either the γ or γ' phase, while performing the line scan. Therefore, the diffusion couples were re-heated for 1 hr at 1293 °C and water quenched to reduce the experimental scatter. The compositions of initial René-N4 and René-N5 were

determined by an average of ten microprobe measurements obtained near the ends of the diffusion couples. The values are given in Table 1.

Alloy	Al	Со	Cr	Hf	Мо	Nb	Та	Ti	Re	W
René-N4	4.23	7.79	10.29		1.48	0.47	4.64	3.46		6.38
René-N5	6.18	7.72	7.48	0.15	1.4		7.13		2.85	6.38

Table 1. Alloy Percent Mass Fraction Compositions. Balance Ni.



Figure 1. High and low magnification back scatter images of 10 h René-N5/René-N4 diffusion couple gas cooled from 1293 °C showing the precipitation of γ' . This coarse structure was eliminated with a post diffusion heat-treatment that included a rapid quench.

The experimental diffusion couples were simulated using the 1-D finite-difference diffusion code, DICTRA^{*18}, in conjunction with the thermodynamic database, Ni-Data⁷, and the Ni diffusion mobility database of Campbell et al². The simulation considered only single phase γ with a planar interface between René-N4 and René-N5. A 200-point geometric grid, which consisted of a higher density of grid points at the center, was used to describe the 6.35 mm couple. The calculations were done with concentration dependent diffusion coefficient matrices as described in eqns. (1-5). The matrices \widetilde{D}_{ij}^{Ni} for the initial compositions are shown in Tables 2 and 3.

^{*} The use of any commercial product does not constitute an endorsement by the National Institute Standards and Technology.

	Al	Со	Cr	Мо	Nb	Та	Ti	W
Al	+119.5	+13.93	+34.83	+34.34	+42.43	+51.50	+49.51	+53.22
Со	-11.37	+17.00	-8.25	-5.67	-5.55	-1.83	-7.10	-9.69
Cr	-4.26	-5.37	+13.67	-3.21	+8.93	+9.91	+8.25	+2.49
Mo	-8.33	-0.280	-0.426	+7.57	-0.55	-0.36	-0.17	-0.45
Nb	+0.31	+0.25	+0.66	+0.27	+24.05	+0.74	+0.85	+0.31
Та	-0.68	+0.33	+0.53	0.24	+0.26	+0.76	+0.50	+0.23
Ti	+1.63	+1.35	+4.94	+4.94	+6.25	+6.57	+23.62	+5.41
W	-1.81	-0.62	-0.55	-0.60	-1.22	-0.83	-0.70	+3.40

Table 2. Interdiffusion Coefficients for René-N4 $(x10^{-14} \text{ m}^2/\text{s})$ at 1293 °C

Table 3. Interdiffusion Coefficients for René-N5 $(x10^{-14} \text{ m}^2/\text{s})$ at 1293 °C

	Al	Со	Cr	Hf	Мо	Re	Та	W
Al	+93.16	+13.92	+33.46	-6.51	+33.42	+25.44	+48.63	+50.87
Co	-6.51	+27.22	-8.56	-27.64	+4.95	-5.11	+3.87	-9.21
Cr	+4.51	-4.23	+21.02	-6.25	-0.22	-0.78	+13.81	+6.89
Hf	+0.86	+0.07	+1.70	+262.1	+1.52	+0.87	+2.37	+1.84
Mo	-0.35	-0.30	-0.30	-1.91	+7.71	-0.25	-0.13	-0.19
Re	-0.75	-0.32	-0.36	-2.59	-0.25	+0.08	-0.51	-0.32
Та	-0.03	+0.33	+0.98	-4.17	+0.64	+0.86	+7.75	+0.87
W	-1.18	-0.57	-0.54	-4.51	-0.39	-0.11	-0.76	+0.59

RESULTS

The experimental diffusion couples were compared to the diffusion simulations using both the composition profiles and the location of the porosity.

To compare the experimental composition profiles with the calculated profiles, the experimental Matano interface was determined by an averaging process. For three of the elements with large concentration differences (Cr, Re, and Ti), the Matano interface was independently located. With respect to the average position, the coordinates for Cr, Re, and Ti profiles for the 100 h treatment were $-15 \mu m$, $-1 \mu m$ and $+17 \mu m$, respectively. These variations are small compared to the large diffusion distances.

Figure 2 shows the agreement between the experimental and calculated diffusion profiles after 10 h and 100 h at 1293 °C. In general, the calculations are in good agreement with the experiments; however, there are some discrepancies especially in the Cr and Re profiles. In the René-N4, the Cr diffuses more slowly than predicted and the Re diffuses more rapidly than predicted. To further evaluate the results, individual profiles are plotted as a function of distance over the square root of time (z/\sqrt{t}), as seen in Figure 3. (Note the composition differences in the Co, Mo, and Hf profiles are within the experimental error of the measurements and thus, the profiles are not evaluated.) For the given length scale and times of 10 h and 100 h for the Al, Cr, Ta, Ti, Re and W profiles the simulated profiles are independent of time.

The overall agreement is confirmed by a simple error analysis for the six profiles in Figure 3 at 100 h. This simple error analysis, shown in Table 4, consists of calculating the difference between the simulated and experimental value at each grid point, summing the differences, and then averaging over the sum of the differences by the total number of grid points. The W profile for 100 h had the smallest average error (1.8 %) and the Re

profile for 100 h had the largest average (5.8 %). These errors are within the experimental scatter. Both 10 h and 100 h experimental Ta profiles show non-monotonic behavior, which is also observed in the calculated profiles but not to the same magnitude.

Al	Cr	Та	Ti	Re	W
3.1 %	2.0 %	2.7 %	2.4 %	5.8 %	1.8 %

Table 4. Error Analysis of 100 h composition profiles



Figure 2. Calculated and experimental composition profiles for René-N4/René-N5 diffusion couples after 10 h and 100 h at 1293 °C.



Figure 3. Composition profiles plotted as function of distance (μm) divided by the square root of time (s).

The location of porosity formed in the experimental couples can also be compared to diffusion simulations. In the lattice-fixed frame of reference, the sum of the net atom fluxes equals the vacancy flux. The negative gradient of the vacancy flux give the number of sinks (sources) necessary to maintain local equilibrium. The expected location of Kirkendal porosity is given by the position of the largest negative value of the gradient, as demonstrated by Höglund and Ågren¹⁹. Figure 4(a) shows this calculation for the René-N4/René-N5 couple. The location of maximum porosity is predicted to be at 65 μ m to the left of the Matano interface on the René-N4 side of the couple. Porosity is observed on the René-N4 side of the interface of the experimental couple as shown in Figure 4(b). Thus, the diffusion simulation and experiment are in qualitative agreement.



Figure 4. (a) Predicted location of where maximum pore density is expected for René-N4/René-N5 at 1293 °C. (b) Back scatter image of René-N4/René-N5 diffusion couple after 100 h at 1293 °C. Thin white line indicates position of microprobe scan. The dashed white line corresponds to the Matano interface. The dashed black line is location of the predicted maximum porosity.

CONCLUSIONS

Experimental diffusion data obtained by isothermal anneals of diffusion couples made of René-N4 and René-N5 in the single phase γ region were compared to calculations using a multicomponent diffusion mobility database. The results compare favorably. The calculations also correctly predict the location of Kirkendall porosity. This initial success points to the suitability of the general approach and the desirability of further database refinements.

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