SUSTAINMENT OF PLASMA ROTATION BY ICRF

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BACKGROUND AND MOTIVATION

- Alcator C-Mod and JET observe development of co-current plasma rotation in ICRF-heated discharges.
- ICRF heating introduces zero (or negligible) angular momentum to the plasma.
 - Experiments have a symmetric k_{\parallel} spectrum and contribute no net angular momentum.
 - Even if the k_{\parallel} spectrum lauched is one sided, the angular momentum input is negligible ($k_{\parallel} = n/R$; $n \approx 12$ for C-Mod)

$$(\Delta T)_{RF} = RF \text{ Torque } = M \Delta v_{\parallel} R = n \Delta E / \omega$$

$$(\Delta T)_{NBI}$$
 = typical NBI torque = $\Delta E R v_{beam}^{-1}$

$$\frac{\Delta T_{\rm RF}}{\Delta T_{\rm NBI}} = \frac{n v_{\rm beam}}{\omega R} \approx \frac{1}{15} <<1$$

• What is the mechanism for developing toroidal rotation and how does it scale?

REPRESENTATIVE EXPERIMENTS

1. Paper by J. E. Rice, et al. [Nuclear Fusion 39 (1999) 1175] reports rotation observations and scaling.



MODEL OVERVIEW - 1

- 1. Even though ICRF heating introduces no net torque, there remains the possibility of creating positive and negative torque density regions.
- 2. Describe plasma response to torque density by an angular momentum diffusion equation.
 - Separated torque density regions lead to finite central rotation
- **3.** Model ICRF heating by the introduction of energetic particles on the equatorial plane and the removal of an equal number of cold particles.
 - Particles are introduced at a particular flux surface the resonance location— with equal numbers of co- and counter- velocities so there is no angular momentum input.

REPRESENTATIVE INITIAL ORBITS



• Fast-wave refraction leads to midplane heating.

MODEL OVERVIEW - 2

- 4. Follow particles by ORBIT with ion-ion pitch angle and drag collisions
 - Record particle's flux-surface when Energy \rightarrow 0.
 - Particle's displacement from originating flux-surface drives a radial neutralizing current and a $j_r B_{\theta}$ torque density in the background plasma
 - Continuous creation of energetic particles drives steady $j_{r}\,\mbox{current}$
 - ORBIT also computes the torque density imparted to the background by energetic-ion collisions
 - Total volume-integrated applied torque vanishes to 2.10-4 accuracy.
- 5. Compute non-vanishing central rotation from torque density
- 6. Investigate scaling of central rotation and sensivity to initial conditions
 - Particle energy and pitch, resonance location and q.

ANGULAR MOMENTUM DIFFUSION EQUATION-1

1. General Form of angular rotation rate Ω response to torque density τ

$$\frac{\partial}{\partial t} \left(M \ n R^2 \Omega \right) \ = \ \boldsymbol{\nabla} \cdot \left\{ n \ M \ R^2 \ \boldsymbol{\chi}_{M} \ \boldsymbol{\nabla} \Omega \right\} \ + \tau$$

2. Steady-state axisymmetric version

$$\frac{1}{\mathbf{V'}} \frac{\partial}{\partial \psi} \left\{ \mathbf{V'} \left\langle n \ M \ R^2 \chi_M \left(\nabla \ \psi \right)^2 \right\rangle \frac{\partial \Omega}{\partial \psi} \right\} = - \left\langle \tau \right\rangle$$

$$V' = \oint \frac{d\ell 2\pi R}{\nabla \psi} = \frac{\partial V}{\partial \psi}$$
 and <> denotes magnetic surface average

- 3. $\langle \tau \rangle$ is torque density on bulk plasma and has two sources:
 - $j_r B_\theta$ torque arising from radial curents which neutralize energetic particle displacements
 - Collisional Angular momentum transfer from energetic particles.

ANGULAR MOMENTUM DIFFUSION EQUATION-2

4. First integral of angular momentum equation

$$\mathbf{V'}\left\langle n \mathbf{M} \mathbf{R}^2 \chi_{\mathbf{M}} \left(\nabla \psi \right)^2 \right\rangle \frac{\partial \Omega}{\partial \psi} = -\int_0^{\psi} \left\langle \tau \right\rangle \mathbf{V'} \, d\psi = \mathbf{T}(\psi)$$

- $T(\psi)$ = torque exerted inside ψ -surface
- No net torque condition: $T(\psi_{max}) = 0$
- 5. Apply no-slip boundary condition at surface
 - Field lines outside separatrix line-tied to vessel; toroidal rotation not permitted
- 6. Torque proportional to rate of creation of energetic particles N and angular momentum transferred per particle.

 $\left< \tau \right> \propto \dot{N}$

ANGULAR MOMENTUM DIFFUSION EQUATION-3

7. Angular rotation rate (use toroidal flux Φ as independent variable)

$$\Omega(\Phi) = \int_{\Phi}^{\Phi_{max}} \frac{d\Phi}{q \, V'} \frac{T(\Phi)}{\left\langle n \, M \, R^2 \, \chi_M \, \left(\nabla \psi \right)^2 \right\rangle}$$

8. Conclude:

For regions of separated postive and negative torque density, T(Φ) is non-zero and toroidal rotation can develop, even though the total torque T(Φ_{max})=0

ION-CYCLOTRON HEATING MODEL

- 1. The ion-cyclotron heating process changes a particle's prependicular energy, while leaving v_{\parallel} and the canonical angular momentum unchanged.
 - No net angular momentum is introduced
- 2. Our ICRF model replaces cold particles by energetic particles constrained to have an equal number of co- and counter velocity particles.
 - Particles are created on the same flux surface on the midplane
 - Mimics ICRF heating for particles whose orbit is tangent to the resonant surface at the midplane where the fast wave intensity is high.
 - Pitch at creation is fixed to be low: $v_{\parallel}/v = (0.25 0.40)$
- **3.** Energetic particles will spatially diffuse via banana diffusion and will collisionally transfer angular momentum to the bulk plasma.
 - ORBIT code follows these processes via a Monte Carlo approach.

ORBIT CODE

- **1. ORBIT** code has been developed to follow energetic particle orbits in toroidal confinement geometries of arbitrary cross section.
- 2. Hamiltonian formalism developed
 - Rigorous Hamiltonian form found:
 - R. B. White and M.S. Chance, Phys. Fluids 27, 2455 (1984)
 - R. B. White, Phys. Fluids B2, 845 (1990)

 $\frac{dP_{\zeta}}{dt} = -\frac{\partial H}{\partial \zeta} \qquad \frac{d\zeta}{dt} = \frac{\partial H}{\partial P_{\zeta}}$ $\frac{dP_{\theta}}{dt} = -\frac{\partial H}{\partial \theta} \qquad \frac{d\zeta}{dt} = \frac{\partial H}{\partial P_{\theta}}$

- 3. Monte- Carlo collisions after A. Boozer et al. Phys. Fluids 24, 851 (1981)
- 4. Collision model: Energetic proton ion-ion collisions with cold deuterons.

$$\frac{d\left\langle \theta^{2}\right\rangle}{dt} = \nu_{o} \left(\frac{E_{o}}{E}\right)^{3/2} \qquad \frac{1}{E} \frac{dE}{dt} = -\nu_{o} \left(\frac{E_{o}}{E}\right)^{3/2} \frac{M_{\text{proton}}}{M_{\text{deuteron}}} \qquad \nu_{o} = \frac{2^{3/2} \pi n e^{4} \ln \Lambda}{\left(M_{p}\right)^{1/2} E_{o}^{3/2}}$$

ORBIT CODE MODIFICATIONS

- Plasma divided into 5.104 bins in toroidal flux (magnetic surface label)
- For each time step, momentum transfer from particles to plasma through pitch angle scattering and drag recorded in each bin.
- Final particle momentum and density recorded in each bin
- Integrals over toroidal flux (bins) needed for angular rotation performed
- Angular momentum check accurate to 1 part in 5000.

NONDIMENSIONAL CENTRAL ROTATION RATE-1

1. Let Φ_0 denote the toroidal flux value where energetic particles of energy E are introduced at a rate \dot{N} .

- 2. Let $v = (2E/M)^{1/2} (\Omega_a R_a)^{-1}$ denote a nondimensional particle speed.
- 3. ORBIT computes $F(\Phi)$ = fraction of particles ending up inside flux surface Φ and the integral T_1 of the $j_r \times B_{\theta}$ torque.

$$T_{1}(\Phi) = \frac{1}{V} \int_{0}^{\Phi} \frac{d\Phi'}{q} G(\Phi') \qquad G(\Phi) = \begin{cases} F(\Phi) & \Phi \leq \Phi_{o} \\ F(\Phi) - 1 & \Phi \geq \Phi_{o} \end{cases}$$

- 4. ORBIT also calculates v $T_2(\Phi)$ = mechanical angular momentum deposited inside Φ .
- 5. Standard circular tokamak formulas, an assumed constant momentum diffusivity $\chi_M = a^2/6\tau_M$, and $\dot{N} E \tau_E = P \tau_E = W$ are employed to calculate the central rotation frequency

NONDIMENSIONAL CENTRAL ROTATION RATE-2

6. On-axis rotation rate is expressed in terms of the nondimensional rotation rate I*

$$\frac{\Omega(0)}{\dot{N}} = v^2 I^* \qquad I^* = \frac{1}{v} \int_0^{\Phi_{max}} \frac{d\Phi}{\Phi} T \qquad T = T_1 + T_2$$
$$v = \left(2E / M\right)^{1/2} \left(R_a \omega_{c,a}\right)^{-1}$$

7. Analytic considerations motivate the introduction of v so that I* is insenstive to physics parameters

ROTATION IN PHYSICAL UNITS

- **1.** Select baseline initial particle values used in computing I* to be representative of Alcator C-Mod.
 - E=48 keV, pitch = 0.25, rho = 0.165, low-field midplane, and N=2000.
 - Result:

$$I^* = 22.5$$

2. In physical variables the rotation rate is

$$\Omega(0) = \left\{ \frac{6 \text{ W}}{e B_a R_a^{3} a^2 \bar{n} (2\pi)^2} \left(\frac{\tau_M}{\tau_E} \right) \right\} I^*$$

For the shot on sheet 4, this gives $v_{tor} = \Omega(0) R_a = 7 \cdot 10^4 \text{ m/s}$, in good accord with the reported value. Results insensitive to E, pitch, and N.

INITIAL ORBITS

1. Initial orbits are characteristic of orbits near the magnetic axis



ROTATION AND TORQUE PROFILES - RUN 1

-0.8

0.2

0.4

0.6

• Rotation Profile is peaked.

Nondimensional Rotation

rho

25 **Integrated Torque Profiles - Run 1** Nondimensional Rotation Integral I^{*} 0.8 20 0.6 15 0.4 jxB torque total torque Mechanical torque Integrated Torque 10 0.2 0 5 -0.2 0 -0.4 -5 -0.6 0.2 0.8 0 0.4 0.6 1 1.2

0.8

1

1.2

SENSIVITY STUDIES

• How much does the central rotation change as E, rho, pitch, q-profile, and initial surface (ICRF resonance surface) vary ? Results expressed as I*.

Run	Objective (N=500)	I *	rho	E (keV)	pitch	q _{max}	resonance
1	Baseline (N=2000)	22.5	0.165	48	0.25	4.0	LFS
1.1	Baseline (N=200)	24.9	0.165	48	0.25	4.0	LFS
2	Pitch variation	28.3	0.165	48	0.35	4.0	LFS
3	Energy dependence	24.6	0.165	24	0.34	4.0	LFS
4	HFS vs LFS (run 1)	-18.6	0.165	48	0.25	4.0	HFS
5	q max	17.5	0.165	48	0.25	8.0	LFS
6	initial rho	13.5	0.34	48	0.5	4.0	LFS
7	initial rho	19.3	0.69	48	0.64	4.0	LFS
8	Banana vs	11.4	0.34	48	0.32	4.0	LFS
	Circulating (run 6)						
9	HFS vs LFS (run 7)	-22	0.69	48	0.64	4.0	HFS
10	On axis	7.3	0.0	48	0.35	4.0	On-axis
11	On axis - pitch	-1.9	0.0	48	0.25	4.0	On-axis

ROTATION CURVES vs INITIAL RHO



• As resonance layer is moved outward, rotation profile broadens and is lower in magnitude than baseline case.

SUMMARY - 1

- **1.** Separated regions of positive and negative torque density can generate central rotation
 - General property of a diffusion equation
- 2. ICRF generates two types of torque density on bulk plasma which are comparable in magnitude and integrate to zero net torque
 - $j_r \times B_{\theta}$ and mechanical angular momentum transfer by collisions
- 3. ORBIT code follows individual particles and computes the torque densities
 - ICRF model (initial condition for ORBIT) replaces a cold particle by an energetic particle in the mid-plane.
 - Equal numbers of co- and counter energetic particles assure not net momentum injection. Angular momentum check to 2.10-4 level.

SUMMARY - 2

- 4. Central rotation arises
 - Co-current sense, magnitude, and scaling in accord with Alcator C-Mod
 - $v_{exp}(0) = 10.0 \cdot 104$ m/s $v_{model}(0) = 7 \cdot 104$ m/s
 - Insensitve to particle energy, pitch, q_{max} , N, and initial ρ .
 - High-field-side initial ρ gives counter-current rotation.
- 5. Summary formula

$$\mathbf{v}_{tor}(0) = \left\{ \frac{6 \text{ W}}{e B_a R_a^2 a^2 n (2\pi)^2} \left(\frac{\tau_M}{\tau_E} \right) \right\} \mathbf{I}^*$$
$$\mathbf{I}^* = 10\text{-}20$$

CONCLUSIONS

- A mechanism to create central rotation in tokamaks with ICRF heating has been indentified.
- Toroidal velocity scales diamagnetically
 - Magnitude and sense in accord with C-Mod data.
- Precise treatment of angular momentum needed and provided by ORBIT code





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See Poster RP1.72 Thursday Afternoon (Also, poster RP1.71 - Y. Omelchenkov)