Periodic and Quasi-Periodic Compensation Strategies of Extreme Outages Caused by Polarization Mode Dispersion and Amplifier Noise¶

V. Chernyak[1], M. Chertkov[2], I. Kolokolov[2, 3], and V. Lebedev[2, 3, *]

[1]Corning, NY 14831, USA
[2]Theoretical Division, LANL, Los Alamos, NM 87545, USA
[3]Landau Institute for Theoretical Physics, Russian Academy of Sciences, Moscow, 117334 Russia
*e-mail: lebede@landau.ac.ru
Received July 16, 2003

The effect of birefringent disorder on the bit error rate (BER) in an optical fiber telecommunication system subject to amplifier noise may lead to extreme outages, related to anomalously large BER values. We analyze the probability distribution function of BER for various strategies of polarization mode dispersion compensation. A compensation method is proposed that is capable of more efficient extreme suppression of outages, which leads to substantial improvement of the fiber system performance. © 2003 MAIK "Nauka/Interperiodica".

PACS numbers: 42.79.Sz; 42.81.Gs; 05.40.Ca

Polarization mode dispersion (PMD) is an essential impairment for modern optical fiber systems [1–3]. Therefore, dynamical PMD compensation has become an important subject in modern communication technology [4–7]. Optical noise generated in optical amplifiers represents another impairment that may not be reduced or compensated and, therefore, should also be considered in any evaluation of fiber system performance [8]. The bit error rate (BER) calculated for a given realization of birefringent disorder by means of averaging over the amplifier noise statistics constitutes an appropriate object for characterizing the joint effect of the two impairments. In two preceding papers [9, 10], we have demonstrated that the probability of extreme outages (BER values much higher than typical) is substantially larger than one could expect from naive Gaussian estimates singling out effects of either of the two impairments. The natural object of interest is the probability distribution function (PDF) of BER and, specifically, the PDF tail corresponding to anomalously large BER. In [9] we have developed a consistent theoretical approach to calculating this tail. The case when no compensation is applied and also the effect of the simplest "setting of the clock" compensation on the PDF tail suppression have been discussed in [9]. Our investigation was then extended to study effects of the standard first- and higher-order compensations on extreme outages [10]. In the present letter we propose a compensation scheme that appears to be more efficient in reducing the extreme outages compared to the tradi-

tional high-order compensation scheme with the same number of compensating degrees of freedom.

We consider the return-to-zero (RZ) modulation format, when optical pulses are well separated in time t and thus can be analyzed as individual objects. We represent the pulse intensity measured at the system output as

$$I = \int dt\, G(t) |\mathcal{K}\boldsymbol{\Psi}(t)|^2, \qquad (1)$$

where $G(t)$ is a convolution of the electrical (current) filter function with the sampling window function. The two-component complex field $\boldsymbol{\Psi}(t)$ describes the output optical signal (the components correspond to two polarization states of the signal). The linear operator $\mathcal{K}$ in Eq. (1) represents optical filtering; it may also account for a compensating device. The compensating part of the linear operator, $\mathcal{K}_c$, is applied first, i.e., before filtering described by $\mathcal{K}_f$, resulting in $\mathcal{K} = \mathcal{K}_f \mathcal{K}_c$. Ideally, I takes two distinct values depending on whether the information slot is vacant or filled. However, the impairments enforce deviations of I from those fixed values. If the output signal intensity exceeds the decision level I_d, then "1" is associated with the slot, otherwise the slot is labeled "0." Sometimes the information is lost; i.e., the initial "1" is detected as "0" at the output or vice versa. The BER is the probability of such events that naturally depends on a specific realization of birefringent disorder in the fiber. BER must be extremely small to guarantee successful system performance. It has been demonstrated in [9] that anomalously high BER originates solely from the "1 ⟶ 0"

¶ This article was submitted by the authors in English.

events. We denote the probability of such events by B and study its sensitivity with respect to disorder.

In this letter we restrict ourselves to the linear propagation regime, when the output signal $\boldsymbol{\Psi}(t)$ can be represented as a sum of two contributions: $\boldsymbol{\varphi}$, related to the noiseless evolution of the initial pulse, and the noise-induced part $\boldsymbol{\phi}$. We consider the cases of distributed or, alternatively, lumped amplification with the fiber length Z substantially exceeding the interamplifier separation (span length) within the same scope. $\boldsymbol{\phi}$ becomes a zero-mean Gaussian variable, completely characterized by its two-point correlation function

$$\langle \phi_\alpha(t_1)\phi_\beta^*(t_2) \rangle = D_\xi Z \delta_{\alpha\beta}\delta(t_1 - t_2), \qquad (2)$$

insensitive to particular realizations of birefringent disorder and chromatic dispersion in the fiber. The product $D_\xi Z$ represents the amplified spontaneous emission (ASE) spectral density accumulated along the fiber. The coefficient D_ξ is introduced into Eq. (2) to reveal the linear growth of the ASE factor with Z [8]. The noise-independent part of the signal is governed by
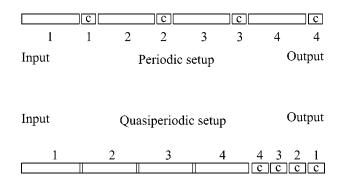
$$\partial_z\boldsymbol{\varphi} - \hat{m}(z)\partial_t\boldsymbol{\varphi} - id(z)\partial_t^2\boldsymbol{\varphi} = 0, \qquad (3)$$

z and d being the coordinate along the fiber and chromatic dispersion. The birefringence matrix can be represented as $\hat{m} = h_j\hat{\sigma}_j$, where h_j is a real three-component field and $\hat{\sigma}_j$ are the Pauli matrices. Averaging over many states of the birefringent disorder that any fiber goes through over time, or (alternatively) over the states of birefringence of different fibers, one finds that $h_j(z)$ is a zero-mean Gaussian field described by

$$\langle h_i(z_1)h_j(z_2) \rangle = D_m\delta_{ij}\delta(z_1 - z_2). \qquad (4)$$

If birefringent disorder is weak, the integral $\mathbf{H} = \int_0^Z dz\mathbf{h}(z)$ coincides with the PMD vector. Thus, $D_m = k^2/12$, where k is the so-called PMD coefficient.

In an operable communication system, typical damage caused by disorder and noise must be small, i.e., typically both impairments can cause only a small distortion to a pulse; thus, the optical signal-to-noise ratio (OSNR) and the ratio of the squared pulse width to the mean squared value of the PMD vector are both large. OSNR can be estimated as $I_0/D_\xi Z$, where $I_0 = \int dt|\Psi_0(t)|^2$ is the initial pulse intensity, the integration being performed over a single slot populated by an ideal pulse, encoding "1." Typically, B fluctuates around B_0, the zero-disorder ($h_j = 0$) value of B. A convenient auxiliary dimensionless object, $\Gamma = (D_\xi Z)\ln(B/B_0)/I_0$, depends on the birefringent disorder and the initial signal shape, as well as the details of the compensation and detection procedures; it is, however, insensitive to noise. Since the OSNR is large, even weak disorder can generate a strong increase in the value of B. This is why a perturbative (with respect to

Scheme of fiber-line-element installation corresponding to the periodic and quasi-periodic compensation strategies.

$\mathbf{h}$) calculation of Γ gives the most essential part of the PDF $\mathcal{S}(B)$ of B. If no compensation is applied, one gets $\Gamma \sim H_3/b$, b being the pulse width, and the initial signal is assumed to be linearly polarized. In the simplest case of "setting of the clock" compensation, one arrives at $\Gamma \sim (H_1^2 + H_2^2)/b^2$. This yields the powerlike tail of the PDF of B [9]. Higher order compensation leads to $\Gamma \sim (H/b)^p$, where p is an integer exceeding by one the degree of compensation, provided no additional cancellations occur, and one gets the following asymptotic expression (tail) for the PDF $\mathcal{S}(B)$ of B [10]:

$$\ln\mathcal{S} = -\mu_p b^2[D_\xi Z\ln(B/B_0)/I_0]^{2/p}/[D_m Z], \qquad (5)$$

where μ_p is a dimensional coefficient. Therefore, as anticipated, compensation suppresses the PDF tail. However, applying high-order compensation is not very efficient, since the decrease in $\mathcal{S}$ is mild as p increases.

The main purpose of this letter is to introduce more efficient compensation strategies with the same number of compensating degrees of freedom. As a first example consider the following "periodic" scheme. One divides the optical line into N segments, each having a length $l = Z/N$, and apply the first-order compensation at the end of each segment (as schematically shown in the upper panel of figure, with "c" denoting the compensating elements). The noise-independent part of the compensated signal for the "periodic compensation" strategy is determined by

$$\mathcal{K}_c\boldsymbol{\varphi} = \exp(i\eta\partial_t^2)\mathcal{K}_{1N}\hat{U}_N...\mathcal{K}_{11}\hat{U}_1\boldsymbol{\Psi}_0(t), \qquad (6)$$

$$\hat{U}_n = T\exp\left[\int_{(n-1)l}^{nl} dzh_j(z)\hat{\sigma}_j\partial_t\right], \qquad (7)$$

$$\mathcal{K}_{1n} = \exp\left[-\int_{(n-1)l}^{nl} dzh_j(z)\hat{\sigma}_j\partial_t\right], \qquad (8)$$

where $\boldsymbol{\Psi}_0(t)$ is the input signal profile, $\eta = \int_0^Z dz d(z)$ is the integral chromatic dispersion, and the ordered product on the right-hand side of Eq. (6) is taken over all the N segments (T exp is the standard notation for the so-called ordered exponential). The exponential factor $\mathcal{K}_{1n}$ represents the first-order compensation at the end of the nth segment.

This periodic compensation is not very convenient since it requires installation of compensating elements at multiple places along the fiber. However, one can naturally modify this scheme and have the same compensating elements inserted subsequently, but all at once, at the fiber output, as shown in the lower panel of the figure. If the disorder profile h_j is known (technically such end-point measurements are possible via the anti-Stokes refraction technique [11]) one can have an end-point, but multiple, compensation as $\mathcal{K}_c = \prod \mathcal{K}_{1n}$, leading to the following "quasi-periodic" modification of Eq. (6):

$$\mathcal{K}_c \boldsymbol{\varphi} = \exp(i\eta\partial_t^2)\mathcal{K}_{11}\ldots\mathcal{K}_{1N}\hat{U}_N\ldots\hat{U}_1\boldsymbol{\Psi}_0(t). \quad (9)$$

The natural idea behind this quasi-periodic compensation is obvious: to construct (in the compensating part) the best possible approximation (with a given number of compensating degrees of freedom) for the inverse of the ordered exponential $\hat{U}_N\ldots\hat{U}_1$.

Note that the (quasi-)periodic compensation does not influence the noise-dependent part of the signal; i.e., $\mathcal{K}_c\boldsymbol{\varphi}$ has the same correlation function (2) as $\boldsymbol{\varphi}$. Therefore, one arrives at the same expression $\ln(B/B_0) = \Gamma I_0/(D_\xi Z)$, with a new **h**-dependent factor Γ. Furthermore, in the main region of interest Γ can be analyzed perturbatively, just as in [9, 10]. Expanding the factors in Eq. (6) up to the second order and making use of Eqs. (1), (2) one derives

$$\Gamma \approx \frac{\mu_2'}{b^2} \sum_{n=1}^{N} \int_{a_n}^{nl} dz \int_{a_n}^{z} dz'[h_1(z)h_1(z') - h_2(z)h_1(z')], \quad (10)$$

where $a_n = (n-1)l$. Here, the dimensionless coefficient μ_2' is related to the output signal chirp produced by an initial chirp and the nonzero integral chromatic dispersion η. As follows from Eq. (9), the same expression (10) is obtained in the second order for the quasi-periodic case. Substituting Eq. (10) into the expression for B and evaluating the PDF of B, with the Gaussian statistics of **h** described by Eq. (4), leads to the following expression for the tail of the PDF of B:

$$\mathcal{S}(B)dB \sim \frac{B_0^\alpha dB}{B^{1+\alpha}}, \quad \alpha = \frac{N\pi D_\xi b^2}{2|\mu_2'|D_m I_0}. \quad (11)$$

Equation (11) holds for $\ln(B/B_0) \gg \mu_2' D_m I_0/[D_\xi b^2]$. The exponent α in Eq. (11) contains an additional factor N compared to the expression for the first order end-point

compensation, i.e., the (quasi-)periodic compensation makes the tail of $\mathcal{S}(B)$ steeper. It is instructive to compare the outage probability for the periodic case with the case of higher order end-point compensation described by Eq. (5). One finds that for higher order compensation, i.e., when $N \sim \mu_2' \ln(B_*/B_0)I_0/D_\xi Z$, the (quasi-)periodic scheme becomes more efficient compared to the straight Nth order compensation scheme. If the output signal is not chirped, $\mu_2' = 0$ and the leading term in the expansion of Γ in **h**/b is of the third order. Additional filtering efforts can be made to enforce the output pulse symmetry under the $t \longrightarrow -t$ transformation, thus removing the third-order term. Then the leading term in Γ will be of the fourth order in **h**/b. Finally, even better compensation can be achieved if the standard high-order compensation approach and the (quasi-)periodic ones are combined, i.e., if in the (quasi-)periodic setting considered above, one uses higher order compensation instead of the first order one. Formally, this hybrid case means that the first-order compensation operators $\mathcal{K}_{1n}$ in Eqs. (6), (9) should be substituted by higher order compensation operators $\mathcal{K}_{cn}$. In the hybrid periodic case, Γ can be written as the sum of $\Gamma_n \sim (\int_{a_n}^{ln} dz \mathbf{h}/b)^p$, and, since **h** is short-correlated, Γ_n related to different segments are statistically independent. This leads to the following expression for the PDF tail:

$$\ln \mathcal{S}(B) \sim -\mu_p N^{2(p-1)/p} \frac{[D_\xi Z \ln(B/B_0)b^p/I_0]^{2/p}}{D_m Z}, \quad (12)$$

which is valid at $D_\xi Z/I_0 \ln(B/B_0) \gg N^{1-p/2}(D_m Z/[\mu_p b^2])^{p/2}$.

Note that an important computational step, leading to our major results in Eqs. (11), (12), was the evaluation of Γ perturbatively in **h**. In addition, in the periodic case, Γ is a direct sum of each segment contribution Γ_n, and the perturbative treatment applies separately to each Γ_n, requiring the weakness of the PMD effect at each segment only, i.e., $D_m Z/N \ll b^2$. Therefore, one concludes that even an optical line with practically inoperable (without compensation) characteristics ($D_m Z$, which is on the order of or larger than b^2) can still be used for transmission if N is sufficiently large. Moreover, this observation on the applicability of Eqs. (11), (12) also extends to the quasi-periodic case, in the sense that Eqs. (11), (12) provide an upper bound for the PDF of BER. This is due to an additional, oscillatory with **h**, suppression of Γ_n in the quasi-periodic vs. periodic case. This suppression is especially important for segments strongly separated from their compensating segments.

For illustration purposes, let us briefly discuss the example of a fiber line with typical bit error probability, $B_0 = 10^{-12}$, and $\mu_2' = 0.14$. Assume also that the PMD coefficient, $k = \sqrt{12D_m}$, is 1.5 ps/$\sqrt{km}$; the pulse

width, $b = 25$ ps; and fiber length, $Z = 2.500$ km. Then the dimensionless parameter $D_m Z/b^2$ measuring the relevant strength of the PMD effect is $O(1)$; i.e., without compensation, the PMD effect is large, pulses are destroyed, and no successful transmission is possible: $S(B) = 0(1)$ for any $B > B_0$. If, however, the (quasi-)periodic compensation with $N = 10$ compensation units is utilized, the relevant strength of the PMD effect is substantially reduced, so that $D_m Z/[b^2 N] \approx 0.1$, and $S(B)$ begins to decay with B at $B > B_0$. The system performance can be evaluated in terms of the outage probability $\mathbb{O}$, the probability of B being larger than B_*: $\mathbb{O} = \int_{B_*}^{1} dB \mathscr{S}(B)$. One derives from Eq. (11) that in the (quasi-)periodic case, $\mathbb{O} \approx 0.06$ for $B_* = 10^{-8}$ and $\mathbb{O} \approx 0.02$ for $B_* = 10^{-6}$; i.e., the system performance is substantially improved (to become not yet perfect, but already satisfactory).

To conclude, in this letter we have proposed a (quasi-)periodic compensation scheme which appears to be a strong alternative to the standard higher order compensations. The efficiency of the scheme has been demonstrated. Even though technical implementation of this procedure needs expensive equipment, we anticipate that if this compensation technique is implemented, the reduction in the probability of extreme outages will guarantee an essential overall benefit.

REFERENCES

1. C. D. Poole and J. A. Nagel, in *Optical Fiber Telecommunications*, Ed. by I. P. Kaminow and T. L. Koch (Academic, San Diego, 1997), Vol. 3A, p. 114.

2. F. Heismann, in *ECOC'98 Digest* (1998), Vol. 2, p. 51.

3. J. P. Gordon and H. Kogelnik, Proc. Natl. Acad. Sci. USA **97**, 4541 (2000).

4. T. Ono, S. Yamazaki, H. Shimizu, and H. Emura, J. Lightwave Technol. **12**, 891 (1994).

5. F. Heismann, D. Fishman, and D. Wilson, in *Proceedings of ECOC98, Madrid, Spain* (1998), p. 529.

6. L. Moller and H. Kogelnik, in *Proceedings of ECOC'99* (1999), p. 64.

7. H. Bülow, F. Buchali, W. Baumert, *et al.*, Electron. Lett. **36**, 163 (2000).

8. E. Desurvire, *Erbium-Doped Fiber Amplifiers: Principles and Applications* (Wiley, New York, 1994).

9. V. Chernyak, M. Chertkov, I. Kolokolov, and V. Lebedev, Opt. Lett. (in press); http://arXiv.org/abs/physics/0303013.

10. V. Chernyak, M. Chertkov, I. Kolokolov, and V. Lebedev, Opt. Express **11**, 1607 (2003).

11. B. Huttner, B. Gisin, and N. Gisin, J. Lightwave Technol. **17**, 1843 (1999).